

Electron-He₂⁺ scattering calculation using the R-matrix method: resonant and bound states of He₂

M D Epée Epée¹, O Motapon^{1,2}, D Darby-Lewis³ and J Tennyson³

¹Laboratoire de Physique Fondamentale, UFD Mathématiques, Informatique Appliquée et Physique Fondamentale, University of Douala, P. O. Box 24157, Douala, Cameroon

²University of Maroua, Faculty of Science, P. O. Box 814 Maroua, Cameroon

³Department of Physics and Astronomy, University College London, Gower Street London, WC1E 6BT, UK

E-mail: j.tennyson@ucl.ac.uk

Abstract. The UK molecular R-matrix codes are used to study electron collisions with the He₂⁺ molecular ion. Full configuration interaction calculations are performed to obtain the potential energy curves of the ground X ²Σ_u⁺ and the first excited A ²Σ_g⁺ electronic states of He₂⁺. Resonances, effective quantum numbers and resonance widths as a function of the internuclear separation are determined for the lowest singlet ¹Σ_g⁺, ¹Σ_u⁺, ¹Π_g and ¹Π_u and triplet ³Σ_g⁺, ³Σ_u⁺, ³Π_g, ³Π_u and ³Δ_u states which are relevant for the study of reactive collision of He₂⁺ with low energy electron. In addition bound states are also calculated for each symmetry of He₂ at several geometries.

1. Introduction

The helium molecular ion, He₂⁺, was the first molecule to form in the Universe (Lepp et al. 2002); albeit its appearance is thought to be short-lived. Since then it is thought to have been formed in the remnant of supernovae (Lepp et al. 1990) and in the atmospheres of cool white dwarfs (Stancil 1994). On earth He₂⁺ can form in cool helium-containing plasmas, including fusion plasmas, and its curves can play a role in Penning ionisation of metastable helium atoms (Garrison et al. 1973).

In most environments where He₂⁺ forms the major destruction mechanism is dissociative recombination (DR):



This process relies on the formation of doubly-excited, metastable states of He₂ which provide the route to dissociation. DR rates for He₂⁺ have been measured in storage rings (Urbain et al. 2005). Storage rings have also been used to study inelastic collisions between electrons and He₂⁺ ions (Buhr et al. 2008).

There has been a number of theoretical studies of electron collisions with He₂⁺. These include the dissociative recombination studies of Carata et al. (1999), Royal and Orel (2005) and Royal and Orel (2007) who used He₂ curves and lifetimes computed using the Kohn variational method. These calculations are discussed further below. R-matrix calculations have also been used to study the He₂^{*} system. McLaughlin et al. (1993) computed bound and continuum curves, but only of ³Σ_u⁺ total symmetry. Recently, Celiberto et al. (2016) used an R-matrix calculation to compute cross sections for electron-impact dissociation of He₂⁺.

He₂⁺ and He₂ are, respectively, three and four electron systems; they should therefore be amenable to highly accurate electronic structure calculations. In this work we aim to compute a comprehensive and accurate set of bound and continuum curves for excited states of the helium dimer. These curves will form the input for future studies of key processes involving electron collisions of He₂⁺, such as dissociative recombination.

2. Calculations

2.1. Method

In this work we use the R-matrix method (Tennyson 2010) as implemented in the UKRMol codes of (Carr et al. 2012). This method is based on dividing the configuration space into two distinct regions (Burke 2011) by a sphere, here of radius 12a₀, centred at the centre-of-mass of the molecule. This encloses the wave function of the 3-electron target He₂⁺ ion. In the inner region, the wave functions for the target + scattering electron system (He₂⁺ + electron) is given by:

$$\Psi_k^{N+1}(x_1, \dots, x_{N+1}) = \mathcal{A} \sum_{ij} a_{ijk} \phi_i^N(x_1, \dots, x_N) u_{ij}(x_{N+1})$$

$$+ \sum_i b_{ik} \chi_i^{N+1}(x_1, \dots, x_{N+1}), \quad (2)$$

where \mathcal{A} is the anti-symmetrization operator, u_{ij} are known as continuum orbitals, x_i is the spatial and spin coordinates of electron i , ϕ_i^N is the wave functions of the i^{th} target state and χ_i are two-center L^2 functions constructed as products of target occupied and virtual molecular orbitals. The variational coefficients a_{ijk} and b_{ik} are determined by diagonalizing the hamiltonian matrix.

2.2. Target calculations

It is known that the basis sets play an important role in the quality of the calculation. We use for the present work, the cc-pVTZ Gaussian basis set for He₂⁺, which include polarization functions. An initial set of molecular orbitals was obtained by performing self-consistent field (SCF) calculations for the X ²Σ_u⁺ state of He₂⁺, although in practice the choice of orbitals is not important in a full configuration interaction (FCI) calculation. The two lowest He₂⁺ states, X ²Σ_u⁺ and A ²Σ_g⁺, were included in the close coupling expansion of the trial wave function of the scattering system; the other target states lie too high in energy to contribute significantly at the collision energies considered here. Each target state was represented by an FCI wave function. Our FCI calculations performed at selected bond lengths around the equilibrium position for the ground state X ²Σ_u⁺ and the first excited state A ²Σ_g⁺ of the He₂⁺ molecular ion (see Table 1) were in very close agreement with the high accuracy calculations of Cencek and Rychlewski (1995), Tung et al. (2012), Gadea and Paidarova (1996) and McLaughlin et al. (1993).

2.3. Scattering calculations

Scattering calculations used a two-term close-coupling expansion based on the FCI representation of the He₂⁺ X ²Σ_u⁺ and A ²Σ_g⁺ target states. Continuum functions considered partial waves up to $l = 4$ (g functions) and were taken from Faure et al. (2002). The FCI L^2 functions were generated by allowing all 4 electrons to occupy any target orbital subject only to the constraints of total symmetry.

Calculations were performed for singlet and triple spin symmetries and using C_{2v} point group symmetry. Below results have been recast using standard linear molecule symmetry notation. Calculations were repeated at 25 geometries in the range $R = 1.8$ to $6.0 a_0$.

2.4. Resonance detection and fitting

For the resonant states, the R-matrix is propagated (Morgan 1984) to the distance for results to stabilise, and then matched with a Gailitis expansion (Noble and Nesbet 1984). Here a distance of $200.1 a_0$ was used. Resonances were detected and fitted to a Breit-Wigner profile to obtain their energy (E) and width (Γ) using program RESON (Tennyson & Noble 1984). The calculations used an initial energy grid of 0.73×10^{-3}

Ryd; upon detection of a resonance, RESON then lays down an appropriate grid based on the estimated resonance width. Complex quantum defects were obtained from the resonances using the relations

$$E_r = E_t - \frac{1}{\nu^2}, \quad \Gamma = \frac{4\beta}{\nu^3} \quad (3)$$

where E_t is the energy of the threshold to which the resonance state is associated. The effective quantum number ν is related to the real part of the quantum defect by $\nu = n - \alpha$ where n is an integer. For resonances, the complex quantum defect μ is given by $\mu = \alpha + i\beta$, where estimates of α and β which are assumed to vary smoothly and usually with n , can be obtained by performing scattering calculations above the threshold (Seaton 1983, Tennyson 1988).

The recent study of Little and Tennyson (2014) on resonance states of N₂ found it necessary to develop an enhanced method based on time-delays to characterise N₂^{*} resonances (Little et al. 2016). Here, however, the resonances are generally well spaced and the absence of extra, nearby excited target states means that it is not necessary to consider intruder states. As a result the single, isolated resonance model assumed in the Breit-Wigner fits gives good results.

2.5. Bound states

After solving the inner region problem, the solutions were used to build the R-matrix on the boundary. Outer region wave functions were then integrated to a distance of 30.1 a₀ from where an asymptotic expansion due to Gailitis (Noble and Nesbet 1984) is used. For this work, an improved Runge-Kutta-Nystrom integration procedure, as implemented by Zhang et al. (2011) was used. Bound states were then found using the searching algorithm of Rabadán and Tennyson (1996) with the improved nonlinear, quantum defect-based grid of Sarpal et al. (1991).

3. Results and Discussion

We present here our calculations of resonances and bound states of He₂, effective quantum numbers and widths of singlet and triplet symmetries. Results are shown graphically but a complete set of data is provided in a spreadsheet as supplementary material to this article. These results are compared with the quasidiabatic potentials of Royal and Orel (2005).

The resonance curves of singlet total symmetries ¹Σ_g⁺, ¹Σ_u⁺, ¹Π_g and ¹Π_u are shown in Figure 1. Those corresponding to triplet total symmetries ³Σ_g⁺, ³Σ_u⁺, ³Π_g, ³Π_u and ³Δ_g are represented in Figure 2 where they are compared with those of Royal and Orel (2005). These resonance curves are Rydberg states of the first excited state A ²Σ_g⁺ of He₂⁺ which correspond in the diabatic picture to neutral dissociative states of He₂. Below the crossing point between the resonances and the ion ground state potential

electronic curve, they are extended using the program BOUND. These curves are shown in Figures 1 and 2 with the same colour convention as the corresponding resonances. It is to be noted that there are infinite number of bound Rydberg states which converge on the He₂⁺ X ²Σ_u⁺ ground state ion curve. This means that dissociating resonance undergo a series of interactions close to the crossing zone. These can be seen in the adiabatic potential energy curves of He₂ calculated by Cohen (1976). The states calculated in this zone are highly perturbed. This accounts for the lack of smoothness observed in some of our curves. In any case, the smoothness of potential curves after the crossing point has little influence on the DR cross section.

Figure 2 shows that there is some similarity but not complete agreement between our curves and those of Royal and Orel (2005). However, the complex Kohn variational calculations of Royal and Orel were only performed at three internuclear distances: $R = 2.0, 2.5$ and 2.7 a₀; full curves were obtained by use of smooth functions which matched with the appropriate asymptote. This means that direct comparison of the results of these two calculations is difficult. There are slight differences that can be accounted for by the difference in approaches and choices of model. Figure 2 also shows that the crossing points of our resonances and those of Royal and Orel (2005) with the X ²Σ_u⁺ ground state of the He₂⁺ molecular ion are not at the same position. Our lowest ³Σ_g⁺ and ³Π_g resonance curves in Figure 2 cross the ion curve at a lower energy than the correspondings one of Royal and Orel (2005). Since it is well known that the magnitude of the DR cross sections are very sensitive to the crossing point between the dissociative state and the ion curve, for both the direct or the indirect mechanism, the shift between our curves and those of Royal and Orel (2005) should result in differences in the DR cross section in terms of magnitude and structure.

Effective quantum numbers and resonance widths corresponding to the resonances given in Fig. 1 are presented in Figs. 3 and 4 respectively. In the same way effective quantum numbers and resonance widths that correspond to the resonances given in Fig. 2 are shown in Fig. 5 and 6 where they are compared with the results of Royal and Orel (2005). The approximate s, p, d, f character of the quantum defect shown in Figs. 3 and 5 are labeled according to the separated atom limit of the resonance state to which they belong. Our widths presented in Figs 4 and 6 show significant structure. This structure is a result of the avoided crossings between the curves of the same symmetry and the interaction of many molecular states in the crossing zone as mentioned above. The significant amount of structure in the widths could be an artifact of the Breit-Wigner fitting process used. However, our fitting program, RESON (Tennyson & Noble 1984), computes a goodness factor which is defined as the sum of the absolute residues for the 25 points used in the eigenphase fit. This allows a ready assessment of the quality of the fit. For the lowest two resonances of each symmetry the goodness factor for our fits is excellent, typically 10⁻⁷ radians when the eigenphases change by π radians over the range of the fit. For some of the higher resonances this goodness factor increases; in the worst case it is 0.05 radians in which case the widths may not be well-determined This means, however, that the structure detected in the widths of the lower-lying resonances

is a genuine feature of our calculations which can be attributed to the interaction and crossing of resonances state as the Rydberg series converging on each threshold. These interactions are clearly visible in plots of the effective quantum numbers, Figs. 3 and 5. Such interactions can cause the widths to change dramatically as a function of internuclear separation. This behaviour was also found in detailed R-matrix studies performed by Little and Tennyson (2013) and Chakrabarti and Tennyson (2015) on electron collision with N_2^+ and BeH^+ , respectively.

The comparison of our triplet resonance widths as a function of internuclear separation, R , with those of Royal and Orel (2005) given in Figure 6 is interesting. Royal and Orel, who do not present detailed results for their singlet calculations, assumed that the widths have a Gaussian dependence on internuclear separation. Only in one case, the second resonance of $^3\Pi_u$ symmetry is this behaviour even approximately followed by our calculated results. A previous study on electron impact vibration excitation and dissociation of N_2 (Laporta et al. 2014) has already showed that use of idealised, Gaussian widths can lead to significant differences compared to calculations based on use of the true $\Gamma(R)$. The difference in the crossing points of the ion ground state and the resonance widths will probably lead to differences between the dissociative recombination and excitation cross sections obtained from our data, and those of Royal and Orel (2005).

Given the relatively simple electronic structure of the four electron He_2 system it is worth considering the residual sources of uncertainty in the present calculation. Guidelines for uncertainty quantification in collision problems have recently been given by Chung et al. (2016). Starting first with the target electronic structure: this has been treated using a full CI so, unusually, there are no issues with the convergence of the CI model beyond those of the choice of the original 1-electron basis set. Analysis of the results of Table 1 show very good agreement between our model and calculations with more extended basis sets. In particular, our excitation energies differ by only about 0.01 eV from the most accurate predictions which are due to Gadea and Paidarova (1996). This is the crucial parameter for the scattering calculation since the resonances are associated with Rydberg series converging on the excited $A\ ^2\Sigma_u^+$ state. Considering the scattering calculation, again the use of a full CI model both removes issues of balance between the target and scattering calculation (Tennyson 1996) and is likely to lead to good treatment of target polarisation effects. The unusual nature of the electronic structure of He_2^+ means that there is a very large gap between the first and second electronic excited states meaning that for low-energy calculations the truncation of the calculation at two states is unlikely to introduce any significant approximation. This is not generally true and inclusion of higher states is usually important (Jones and Tennyson 2010, Brigg et al. 2014). The resonance parameters presented in this work involve fitting the eigenphases to a Breit-Wigner form which in itself can be a cause of uncertainty, particularly in the case of the resonance widths. However, as discussed above, our fits for the lowest resonances proved to be very accurate and stable. Finally our curves show some structure as a function of bondlength. Such structures can be

Table 1. Comparison of energies (in Hartree) for the ground state X ²Σ_u⁺ and the first excited state A ²Σ_g⁺ of the He₂⁺ molecular ion at selected bond length.

State	$R(a.u)$	This work	Cencek et al. ^a	Tung et al. ^b	Gadea & Paidarova ^c	McLaughlin et al. ^d	
						Truncated CI	Full CI
X ² Σ _u ⁺	2.0	-4.988 475 2	-4.994 402 35		-4.990 317	-4.982 802 5	-4.990 725 8
	2.042	-4.988 673	-4.994 644 2	-4.994 643	-4.990 579		
	2.5	-4.974 214 6	-4.980 371 96		-4.976 980	-4.968 760 5	-4.976 836 1
	3.0	-4.949 241 1	-4.955 776 65		-4.952 694	-4.943 661 9	-4.951 917 3
A ² Σ _g ⁺	2.0	-4.603 143 0			-4.605 518	-4.594 455 0	-4.604 711 1
	2.042	-4.622 843 4					
	2.5	-4.765 283 4			-4.767 643	-4.757 254 5	-4.766 982 9
	3.0	-4.835 105 2			-4.837 597	-4.827 267 8	-4.836 745 1

^a Extrapolated value of Cencek and Rychlewski (1995).

^b Explicitly correlated Gaussian calculation of Tung et al. (2012).

^c Full CI calculation of Gadea and Paidarova (1996).

^d R-matrix calculation of McLaughlin et al. (1993) which used a (4s, 2p, 2d) Slater type orbital (STO) basis.

resolved by performing calculations on a very fine grid (Little and Tennyson 2014), although in practice DR calculations have thus far chosen to ignore this structure (Little et al. 2014). All the curves presented in this paper are adiabatic and ignore any effects due to the breakdown of the Born-Oppenheimer approximation. In particular, no attempt has been made to produce so-called energy-dependent resonance widths which result from non-adiabatic effects (Nestmann 1998). Given the reliability of other stages of the calculation, it is likely that effect of these non-adiabatic couplings is the most significant approximation in our calculations.

4. Conclusion

Using the UK R-matrix molecular codes, we have studied electron collisions with the He₂⁺ molecular ion. Electronic energy curves of resonances, widths and effective quantum numbers were generated as a function of the geometry for the singlet ¹Σ_g⁺, ¹Σ_u⁺, ¹Π_g and ¹Π_u and triplet ³Σ_g⁺, ³Σ_u⁺, ³Π_g, ³Π_u and ³Δ_g symmetries of the He₂ molecule. Below the ion ground state, bound states were also calculated. Some differences are found between our resonance curves and widths and the former computations as well as their crossing positions. For example, at some internuclear distances, our resonance widths are larger than those of Royal and Orel (2005). Moreover, resonance curves and widths are calculated for many more bond lengths, making them suitable for appropriate use in DR computations, as the fit of the autoionization widths in this case would lead to more realistic electronic couplings. As a consequence, use of the present data set is likely to result in differences in the cross sections and rate coefficients predicted for some DR and

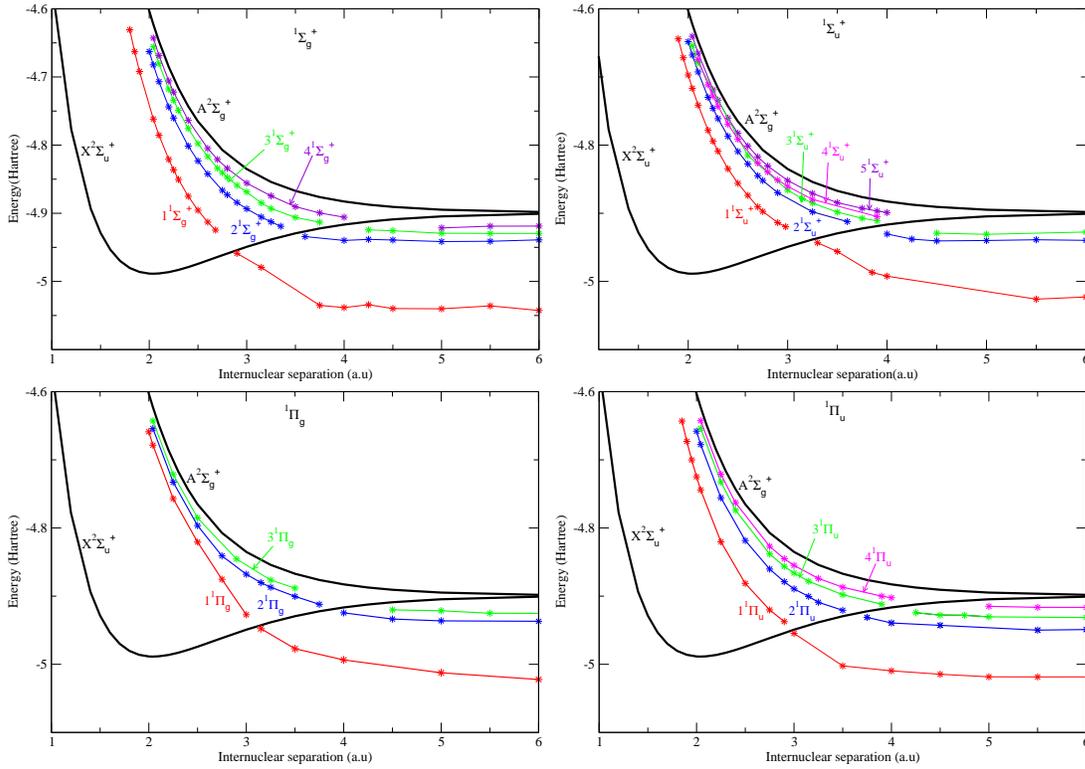


Figure 1. Electronic energy curves of He₂ resonance states of singlet symmetry, with calculated points indicated by stars. The symmetry of each set of the resonances is indicated in the panel, colour is used to match the molecular state corresponding to each curve with the resonance widths given in figure 4. The two black thick full curves are the ground X $2\Sigma_u^+$ and the first excited A $2\Sigma_g^+$ states of He₂⁺.

related processes compared to the ones already available. The present data will be used to produce new dissociative recombination and dissociative excitation cross sections in the near future.

5. Acknowledgements

The authors are grateful to the International Atomic Energy Agency (IAEA, Vienna) for scientific and financial support through the Contract No. 16712 with the University of Douala (CAMEROON) and via the Coordinated Research Projects "Atomic and Molecular Data for State-Resolved Modeling of Hydrogen and Helium and their Isotopes in Fusion Plasma". They also thank Professor Ioan F. Schneider of the University of Le Havre for many helpful discussions. MDEE thanks the Department of Physics and Astronomy of the University College of London for hospitality and a Royal Society Wolfson Research Merit Award to JT which helped support his visit.

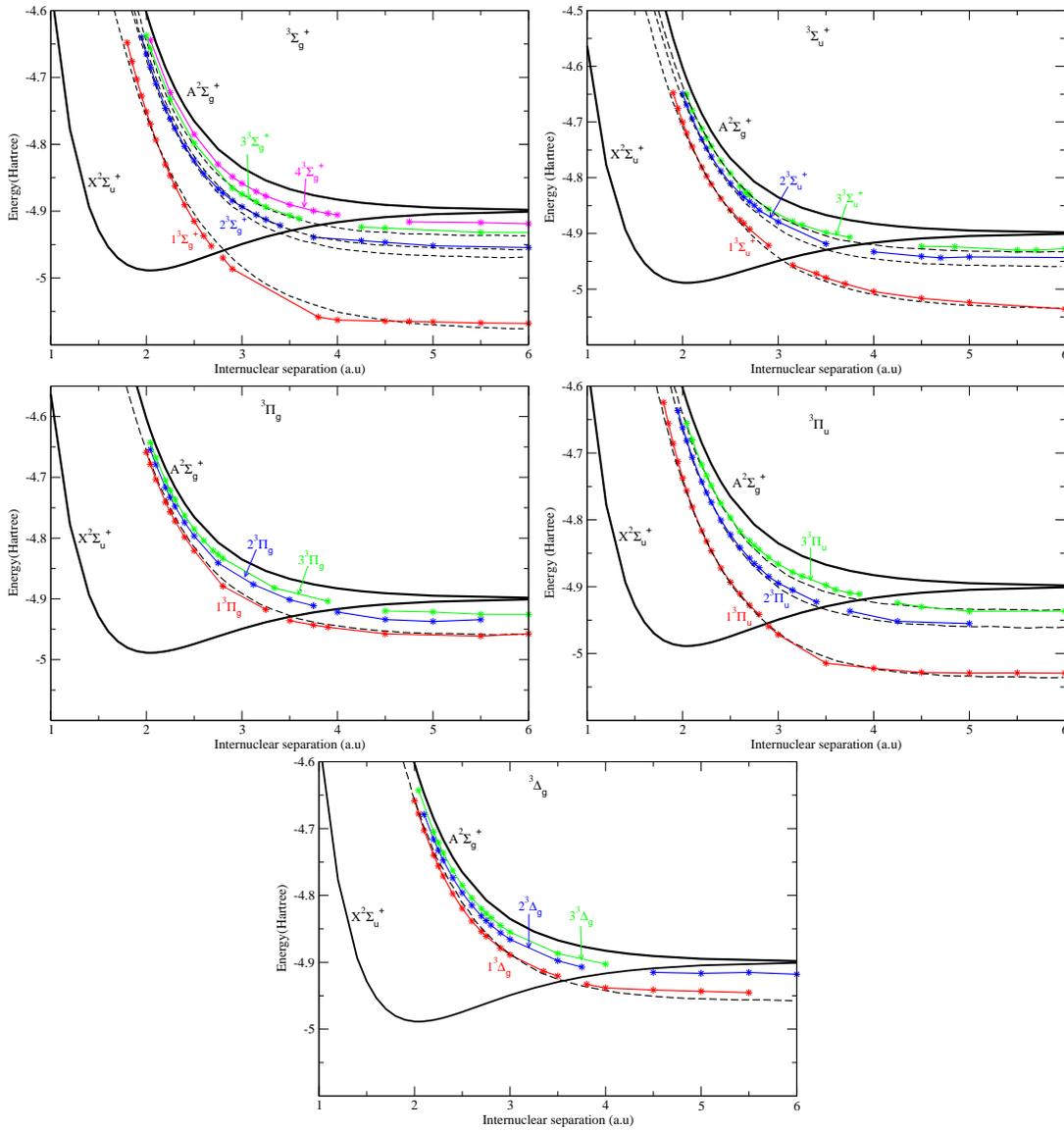


Figure 2. Electronic energy curves of He₂ resonance states of triplet symmetry with calculated points indicated by stars. The symmetry of each set of resonances is indicated in the panel; colour is used to match the molecular state corresponding to each curve with the resonance widths given in figure 6. The two black thick full curves are the ground X $2\Sigma_u^+$ and the first excited A $2\Sigma_g^+$ states of He₂⁺. Dashed curves are the quasideiabatic potentials of Royal and Orel (2005)

References

- Brigg W J, Tennyson J & Plummer M 2010 *J. Phys. B: At.Mol. Opt. Phys.* **47**, 185203.
 Buhr H, Pedersen H B, Altevogt S, Andrianarijaona V M, Kreckel H, Lammich L, Novotny S, Strasser D, Hoffmann J, Lange M, Lestinsky M, Mendes M B, Motsch M, Novotny O, Schwalm D, Urbain X, Zajfman D & Wolf A 2008 *Phys. Rev. A* **77**, 032719.
 Burke P G 2011 *R-Matrix Theory of Atomic Collisions: Application to Atomic, Molecular and Optical Processes* (Berlin: Springer-Verlag)

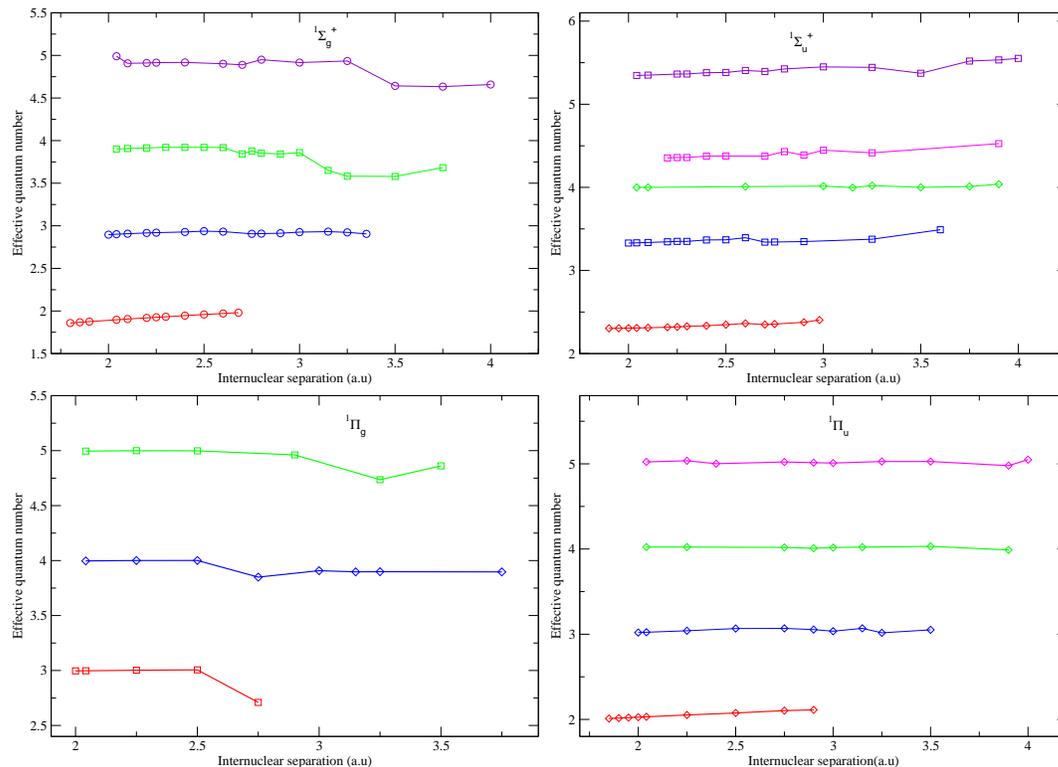


Figure 3. Effective quantum numbers corresponding to the resonances shown in Figure 1 as a function of the internuclear distance R , for the $1\Sigma_g^+$, $1\Sigma_u^+$, $1\Pi_g$ and $1\Pi_u$ symmetries. The nature of the states is indicated with the symbols o : s-state, \diamond : p-state, \square : d-state

- Carata L, Orel A E & Suzor-Weiner A 1999 *Phys. Rev. A* **59**, 2804-2812.
- Carr J M, Galiatsatos P G, Gorfinkiel J D, Harvey A G, Lysaght M A, Madden D, Masin Z, Plummer M & Tennyson J 2012 *Euro. J. Phys. D* **66**, 58.
- Celiberto R, Baluja K L, Janev R K & Laporta V 2016 *Plasma Phys. Control. Fusion* **58**, 014024.
- Cencek W & Rychlewski J 1995 *J. Chem. Phys.* **102**, 2533.
- Chakrabarti K & Tennyson J 2015 *J. Phys. B: At.Mol. Opt. Phys.* **48**, 235202.
- Chung H-K, Braams B J, Bartschat K, Császár A G, Drake G W F, Kirchner T, Kokouline V & Tennyson J 2016, *J. Phys. D: Appl. Phys.* **49**, 363002.
- Cohen J S 1976 *Phys. Rev. A* **13**, 86-98.
- Faure A, Gorfinkiel J D, Morgan L A & Tennyson J 2002 *Computer Phys. Comms.* **144**, 224-241.
- Gadea F X & Paidarova I 1996 *Chem. Phys.* **209**, 281-290.
- Garrison B J, Miller W H & Schaefer H F, 1973 *J. Chem. Phys.* **59**, 3193-3198.
- Jones M & Tennyson J 2010 *J. Phys. B: At.Mol. Opt. Phys.* **43**, 045101.
- Laporta V, Little D A, Celiberto R & Tennyson J 2002 *Plasma Sources Sci. Technol.*, **23**, 065002.
- Lepp S, Dalgarno A & McCray A 1990 *Astrophys. J.* **358**, 262-265.
- Lepp S, Stancil P C & Dalgarno A 2002 *J. Phys. B: At.Mol. Opt. Phys.* **35**, R57-R80.
- Little D A & Tennyson J 2013 *J. Phys. B: At.Mol. Opt. Phys.* **46**, 145102.
- Little D A & Tennyson J 2014 *J. Phys. B: At.Mol. Opt. Phys.* **47**, 1105204.
- Little D A, Chakrabarti K, Schneider I F & Tennyson J, 2014 *Phys. Rev. A* **90**, 052705.
- Little D A, Tennyson J, Plummer M, Noble C J & Sunderland A 2016 *Computer Phys. Comms.* (in press) doi=10.1016/j.cpc.2017.01.005.
- McLaughlin B M, Gillan C J, Burke P G & Dahler J S 1993 *Phys. Rev. A* **47**, 1967-1980.

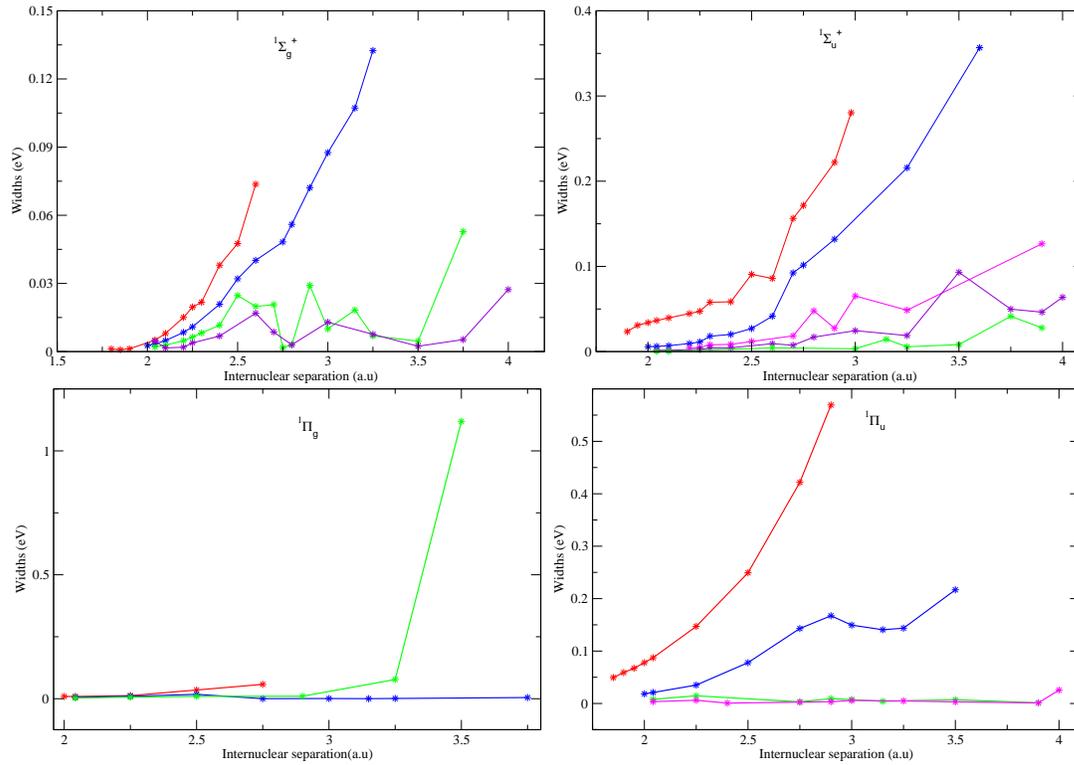


Figure 4. Resonance widths corresponding to the resonances shown in Figure 1 as a function of the internuclear distance R , for the $1\Sigma_g^+$, $1\Sigma_u^+$, $1\Pi_g$ and $1\Pi_u$ symmetries.

- Morgan L A 1984 *Computer Phys. Comms.* **31**, 419.
 Nestmann B M 1998 *J. Phys. B: At.Mol. Opt. Phys.* **31**, 3929-3948.
 Noble C J & Nesbet R K 1984 *Computer Phys. Comms.* **33**, 399-411
 Rabadán I & Tennyson J 1996 *J. Phys. B: At. Mol. Opt. Phys.* **29**, 3747
 Reagen P N, Browne J C & Matsen F A, 1963 *Phys. Rev.* **132**, 304
 Royal J & Orel A E 2005 *Phys. Rev. A* **72**, 022719.
 Royal J & Orel A E 2007 *Phys. Rev. A* **75**, 052706.
 Sarpal B K, Branchett S E, Tennyson J & Morgan L A 1991 *J. Phys. B: At.Mol. Opt. Phys.* **24**, 3685
 Seaton M J, 1983 *Rep. Prog. Phys.* **46**, 167.
 Stancil P C 1994 *Astrophys. J* **430**, 360-370.
 Tennyson J 1988 *J. Phys. B: At.Mol. Opt. Phys.* **21**, 805.
 Tennyson J 1996 *J. Phys. B: At.Mol. Opt. Phys.* **29**, 6185-6201.
 Tennyson J 2010 *Phys. Rep.* **491**, 29-76.
 Tennyson J & Noble C J 1984 *Comput. Phys. Commun.* **33**, 421-424.
 Urbain X, Djuric N, Safvan C P, Jensen M J, Pedersen H B, Sogaard L V & Andersen L H 2005 *J. Phys. B: At.Mol. Opt. Phys.* **38**, 43-50
 Tung W-C, Pavanello M & Adamowicz L 2012 *J. Chem. Phys.* **136**, 104309.
 Zhang R, Baluja K L, Franz J & Tennyson J 2011 *J. Phys. B: At. Mol. Opt. Phys.* **44**, 035203

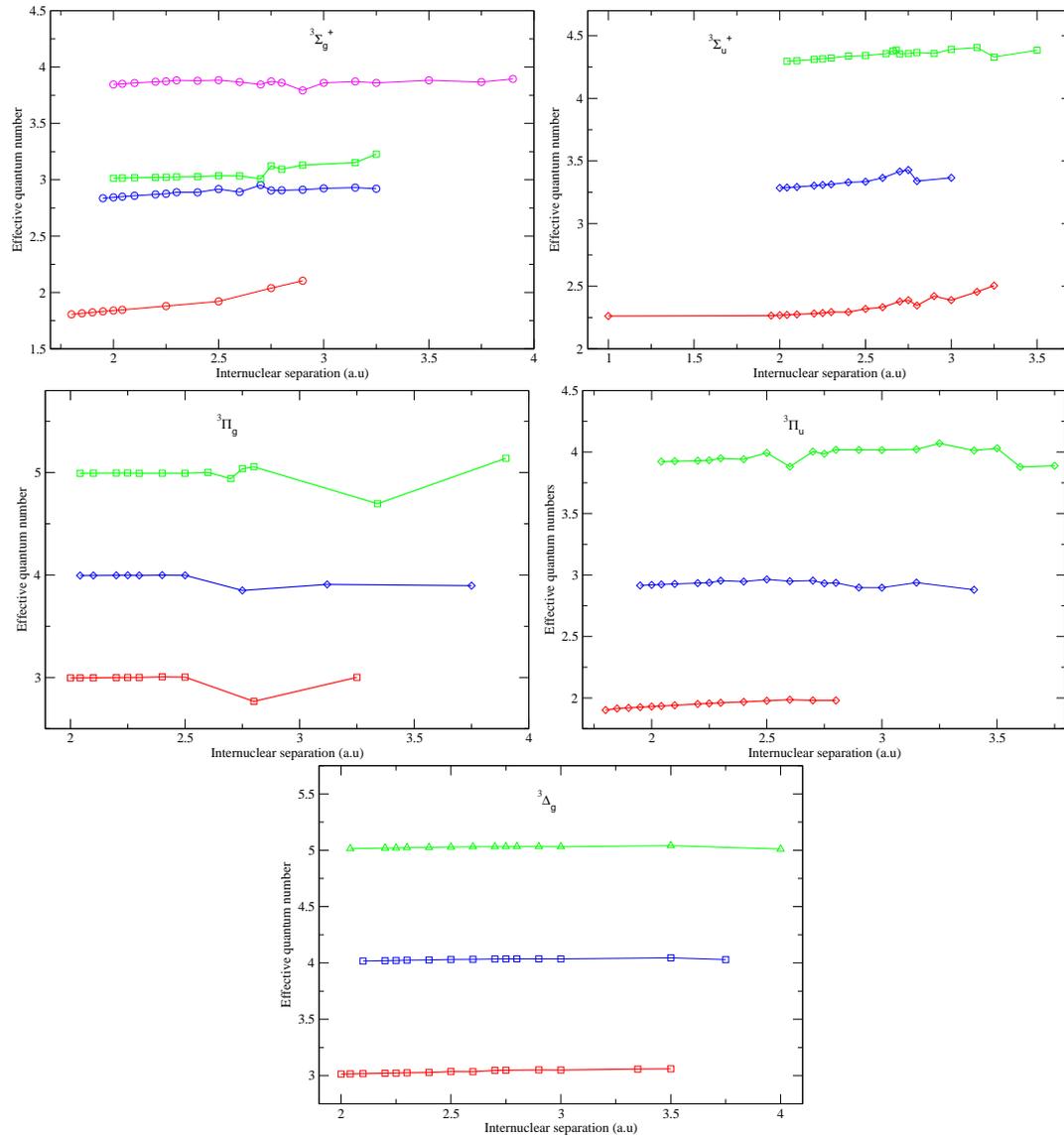


Figure 5. Effective quantum numbers corresponding to the resonances shown in Figure 2 as a function of internuclear distance R , for the ${}^3\Sigma_g^+$, ${}^3\Sigma_u^+$, ${}^3\Pi_g$, ${}^3\Pi_u$ and ${}^3\Delta_u$ symmetries. The nature of the states is indicated with the symbols \circ : s-state, \diamond : p-state, \square : d-state, \triangle : f-state

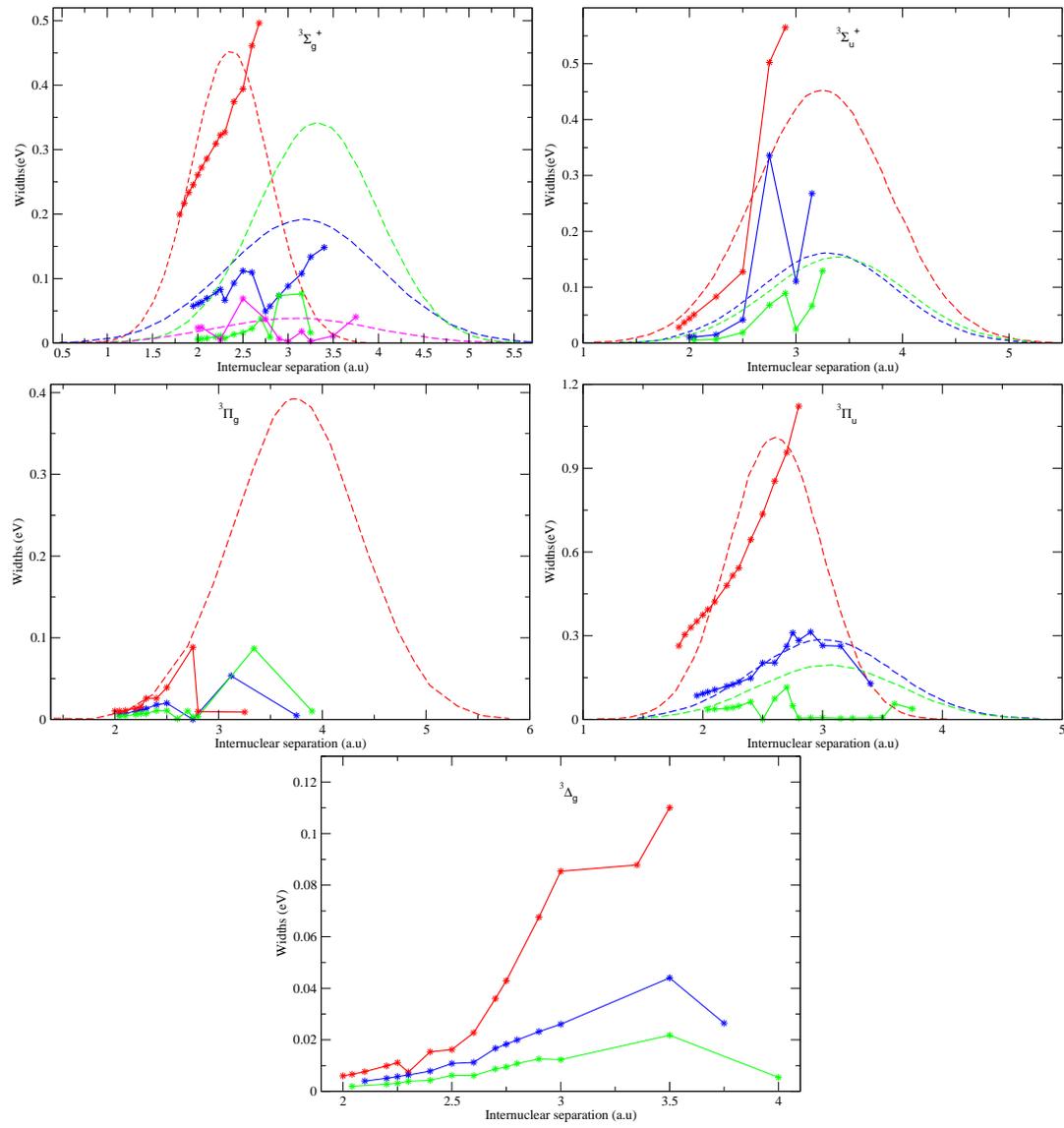


Figure 6. Resonance widths corresponding to the resonances shown in Figure 2 as a function of the internuclear distance R , for the ${}^3\Sigma_g^+$, ${}^3\Sigma_u^+$, ${}^3\Pi_g$, ${}^3\Pi_u$ and ${}^3\Delta_u$ symmetries. Line with star: present work. Dashed curves: Royal and Orel (2005)