

A novel kinematics analysis method using quaternion interpolation – a case study in frog jumping

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1 ABSTRACT

2

3 Spherical Linear Interpolation (SLERP) has long been used in computer animation to

4 interpolate movements between two 3D orientations. We developed a forward kinematics

5 (FK) approach using quaternions and SLERP to predict how frogs modulate jump kinematics

6 between start posture and takeoff. Frog limb kinematics have been studied during various

7 activities, yet the causal link between differences in joint kinematics and locomotor variation

8 remains unknown. We varied 1) takeoff angle from 8 to 60 degrees; 2) turn angle from 0 to

9 18 degrees; and 3) initial body pitch from 0 to 70 degrees. Simulations were similar to

10 experimentally observed frog kinematics. Findings suggest a fundamental mechanism

11 whereby limb elevation is modulated by thigh and shank adduction. Forward thrust is

12 produced by thigh and proximal foot retraction with little contribution from the shank except

13 to induce asymmetries for turning. Kinematic shifts causing turns were subtle, marked only

14 by slight counter-rotation of the left versus right shank as well as a 10% timing offset in

15 proximal foot adduction. Additionally, inclining initial body tilt influenced the centre of

16 mass trajectory to determine direction of travel at takeoff. Most importantly, our theory

17 suggests firstly that the convergence of leg segment rotation axes toward a common

18 orientation is crucial both for limb extension and for coordinating jump direction; and,

19 secondly, the challenge of simulating 3D kinematics is simplified using SLERP because frog

20 limbs approximately follow linear paths in unit quaternion space. Our methodology can be

21 applied more broadly to study living and fossil frog taxa as well as to inspire new control

22 algorithms for robotic limbs.

23

24

25 1.0 INTRODUCTION

26

27 A frog's ability to perform varied locomotor behaviours (e.g. jumping, swimming, walking)  
28 is a hallmark among Anurans. Their multi-functionality has been explored from the  
29 perspective of motor recruitment (d'Avella and Bizzi, 2005; Emerson, 1979; Gillis and  
30 Biewener, 2000; Kamel et al., 1996) as well as foot-substrate interactions (Nauwelaerts et al.,  
31 2005; Nauwelaerts and Aerts, 2003) as a model for how muscular forces interact with the  
32 external environment to determine behaviour (Kargo and Rome, 2002; Kargo et al. 2002;  
33 Aerts and Nauwelaerts, 2009; Clemente and Richards, 2013; Gillis, 2000; Richards, 2011;  
34 Richards and Clemente, 2013, 2012). However, no study has provided a 3D kinematics  
35 analysis to explain how individual limb segments must move differently to achieve diverse  
36 behaviours. Consequently, we lack direct mechanical evidence to assess how the most basic  
37 anatomical features (e.g. absolute limb lengths, limb segment proportions and limb posture)  
38 might influence locomotor multi-functionality in frogs.

39

40 To understand locomotor versatility, one could first use kinematics analysis borrowed from  
41 robotics (e.g. Murray et al., 1994) to map the relationship between joint extensions and body  
42 movements. Secondly, one could record animals performing multiple behaviours, then apply  
43 the kinematics map to resolve how specific joint rotations individually contribute to motion  
44 of the body (Richards et al., 2017). In practice, this approach is restricted by behaviours  
45 animals choose within their natural ability, the limits of which are challenging to elicit in the  
46 laboratory (Astley et al., 2013). Moreover, determining the effect of a single parameter on  
47 performance is difficult because animals often modulate several parameters simultaneously.  
48 For example, frogs change their initial posture, their forelimb extension and their leg  
49 kinematics (Richards et al., 2017; Wang et al., 2014) to increase takeoff angle. To solve the  
50 above problems we use a theoretical kinematics approach where we can dictate the range of  
51 performance and test certain parameters in isolation of others.

52

53 We developed a quaternion-based theoretical forward kinematics approach based on a  
54 computer animation technique called “spherical linear interpolation” (SLERP; Shoemake,  
55 1985; see below for further details). Briefly, SLERP is a powerful technique whereby in-  
56 between motion between landmark time points (i.e. keyframes) can be smoothly interpolated  
57 to fill in gaps. For example, an animated character with an initial posture can be smoothly  
58 moved to a final posture by SLERPing between initial-final keyframes. Despite some  
59 disadvantages such as sharp accelerations with multiple (>2) keyframes (Dam et al., 1998),  
60 SLERP has a key advantage of mathematical and algorithmic simplicity, making its  
61 implementation compact and straightforward. Given the great interest in frogs as a model for  
62 understanding muscle function and muscle-tendon dynamics (e.g. Roberts & Marsh, 2003;  
63 Azizi & Roberts, 2010; Astley, 2016) we use SLERP to explore three behaviours  
64 representing a subset of a frog’s entire locomotor repertoire: 1) straight jumping to different  
65 heights; 2) turning jumps at a fixed height; and 3) straight jumping to a fixed height with  
66 variable initial body angles. Recent findings suggest that final takeoff angle is predicted by  
67 the pre-launch initial posture (Wang et al., 2014) as well as by the launch phase kinematics of  
68 the limbs (Richards et al., 2017). Our theoretical approach enabled us to independently  
69 manipulate pre-jump and takeoff posture to isolate the influence of kinematics from the  
70 effects of start posture. We tested the following hypotheses: H1) The final limb  
71 configuration can be extrapolated (in quaternion space) from the start posture with  
72 knowledge of the target body axis orientation at takeoff. H2) Increased downward rotation  
73 (adduction) is necessary and sufficient to increase jump steepness. In particular, greater  
74 inclination of the pre-jump body posture (Wang et al., 2014) may contribute to increased  
75 jump steepness, but only when followed by increased adduction of the limb segments

76 throughout the jump compared to shallower jumps. H3) A turn to one side is caused by  
77 reduced motion on that side and greater joint extension on the opposite side.

78

79 Adding to the current tools available in biomechanics, the current study introduces a  
80 powerful method to analyse biological motion with only a simple set of mathematical rules.  
81 Such tools are a crucial complement to experiments because they not only permit exploration  
82 of behaviours that are not necessarily observed *in vivo*, but also they allow isolated  
83 manipulation of certain parameters whilst holding all else consistent. Our analysis of how  
84 limb kinematics influences overall body behaviour in frogs is a step towards future  
85 examination of torque modulation (and ultimately muscle forces and neural control) to  
86 coordinate behaviour. For example, our technique can be used to make reasonable  
87 predictions of limb motion in absence of experimental data (e.g. for extinct, rare or  
88 endangered species). Beyond biology, the novel application of SLERP could be applied as a  
89 high-level motion program for bio-robotic limbs (i.e. to produce limb segment trajectories to  
90 be enforced by low-level torque-position controllers). In addition to the applications of our  
91 method, we also discuss biological findings demonstrating that diverse jumping behaviours  
92 can be generated from a single “jump kinematics template”.

93

## 94 2.0 THEORETICAL BASIS

### 95 2.1 Quaternion-based forward kinematics approach

96 To simulate different locomotor behaviours we defined a “joint coordination solver” to: 1)  
97 approximate locomotor kinematics; and 2) modulate the kinematics to achieve different  
98 behaviours. The solver is based on experimental observations demonstrating that jumping  
99 frogs re-orient the axes of rotation of their leg joints; the axes move towards a common  
100 alignment which determines takeoff angle (Richards et al. 2017; Fig. 1B-C, SI Movie 1).  
101 Given these findings, frog jumping kinematics could be simulated by rotating the joint axes

102 towards a common alignment, thus driving leg extension toward a target direction. We  
103 established that our solver obey the coordination rules: 1) rotate the body segments towards  
104 the target by the shortest path; 2) rotate the limb segments towards a common target  
105 orientation; and 3) alter the target jump direction at takeoff to change the jump behaviour. If  
106 our rules approximate realistic jump kinematics (compared to those observed during *in vivo*  
107 experiments) then we have gained insight into the actual coordination mechanisms that real  
108 frogs may employ.

109

110 To simulate frog jumping, we avoid standard forward kinematics approaches using Euler  
111 angles. Briefly, a sequence of 3 cumulative rotations (Euler angles) about Cartesian X-Y-Z  
112 axes (in any chosen order) parameterize 3D rotation. For the present study, Euler rotations  
113 are cumbersome because there exist multiple combinations of angle values leading to the  
114 same rotation. Moreover, they suffer from singularities which can lead to numerical  
115 instability in simulations (Dam et al., 1998) and potentially unnatural motion.

116

117 To avoid the problems of Euler angles, we instead used quaternion SLERP. A quaternion is a  
118 vector of 4 numbers encoding the angle of rotation about a 3D rotation axis; it contains the  
119 same information as a set of XYZ Cartesian axes defining a reference frame. Analogous to a  
120 rotation matrix, a quaternion can perform 3D rotation. There are two important quaternion  
121 properties not shared by rotation matrices. Firstly, a quaternion represents pure rotation, as  
122 opposed to a composition of three Euler rotations. Secondly quaternions can be normalised  
123 to 4D unit vectors (unit quaternions). Thus all quaternions, and therefore all 3D rotations,  
124 reside on the surface of a 4D sphere; moving between any two rotations is achieved simply  
125 by traversing the locally shortest arc along the hypersphere surface. This SLERP technique  
126 (Shoemake, 1985) revolutionised computer animation due to its simplicity and robustness.

127

128 SLERP is also useful for kinematics analysis because it analytically solves the locally  
129 minimum rotation between two orientations (Fig. 2), accomplishing coordination rule 1.  
130 Because of linearity on the hypersphere, we can extrapolate easily. We can then simulate  
131 movement by using a linear extrapolation in unit quaternion space, accomplishing  
132 coordination rule 2. Finally, the direction of extrapolation is determined by the common  
133 orientation to which the segments must converge, accomplishing coordination rule 3.  
134 Applying the rules of the solver, one can analytically determine all kinematics leading to  
135 takeoff knowing only the initial posture *a priori*. More precisely, the mathematical topology  
136 of quaternions allows us to extrapolate the final posture from the initial posture. We will  
137 refer to our approach as *limbSLERP*.

138

## 139 2.2 Animal model system

140 Our model was based on the morphology and jumping kinematics of *Kassina maculata*  
141 Duméril 1853 (the African red-legged running frog). As described in previous publications  
142 (Richards et al. 2017), skin markers placed on the joints (hip, knee, ankle and tarsometatarsal  
143 [TMT]) were assumed to represent locations of the joint centres of rotation; as confirmed by  
144 numerous dissections, the overlying skin is tightly bound to the bones and soft tissues of the  
145 knee, ankle and TMT joints and movement of the skin marker relative to the joint is  
146 negligible. 3D limb kinematics during jumping were recorded using high-speed video  
147 cameras and digitized in MATLAB (Mathworks, Natick, USA) using open source scripts  
148 (Hedrick, 2008). Based on  $\mu$ CT scanning, we used the position of the hip joint as a proxy for  
149 the centre of mass (COM; Porro et al., 2017).

## 150 2.3 Assumptions, definitions and conventions

151 To avoid confusion, we use the following definitions for kinematics. We treated each body  
152 segment as a line, excluding the fore limbs. A single segment, torso, was used to represent  
153 the main body of the frog (head + thorax + abdomen + pelvis). Local reference frames were  
154 defined for each body segment (torso, thigh, shank, proximal foot [tarsus], distal foot), each  
155 defining a local Z axis (Fig. 1A) aligned with the segment long axis. Each reference frame  
156 originates at the proximal endpoint of its segment (origins at the snout, hip, knee, ankle, TMT  
157 for the segments torso, thigh, shank and proximal foot segments, respectively). *Pose* is a  
158 segment's orientation + XYZ position (e.g. the thigh oriented at a given angle originating at  
159 the distal end of the torso). *Configuration* is a list of poses defining the posture of all of the  
160 body segments at a single point in time. Importantly for the current study, we distinguish  
161 between *path* and *trajectory*. Here, a *path* is a particular continuum of poses (or  
162 configurations) traced by an individual segment (or the whole limb) between an initial and  
163 final pose (or configuration), without regard for time. For an example in 2D, the end point of  
164 a pair of segments connected by a hinge joint has the path of an arc traced between the flexed  
165 and extended positions. For the present study, a *trajectory* is a path traced through time. In  
166 the 2D case above, an example trajectory (among infinite possibilities) could be a constant  
167 angular velocity increase along the arc path. An alternative trajectory along the same path  
168 could involve a sinusoidal change in angle such that the trajectory is an oscillating forwards-  
169 backwards motion along the path of the arc. *Scope* is the entire range of motion of the limb,  
170 i.e., all possible paths of all segments between any chosen start or end configuration. For the  
171 present study, scope is constrained only by requiring fixed Euclidean distances between  
172 adjacent body segments. For example, frogs of the same species would share a similar scope  
173 of motion due to their shared segment length proportions, whereas a morphologically distinct  
174 species would have a different scopes of motion. For simplicity, all joints are assumed to be  
175 "ball joints" which rotate freely but do not translate (see Discussion).



176

177 For the present study we avoid analysis of joint angles in the traditional sense. We do not  
178 decompose rotation into flexion-extension, abduction-adduction, internal-external rotation  
179 corresponding to Euler angles about local Cartesian axes (e.g. Kargo & Rome, 2002).  
180 Rather, we work with pure 3D rotations in quaternion form and refer to “extension” and  
181 “flexion” as the opening or closing of a joint, regardless of the orientation of the segments.  
182 More formally, we define “flexion-extension” as a scalar angle within a plane defined by two  
183 connected limb segments, regardless of their orientation in space. For example, vectors  
184 representing the thigh and shank form an invisible plane that can tilt as the femur rotates  
185 about its axis. We say that the knee is “extending” if the shank is moving away from the  
186 thigh, regardless of the plane’s orientation. In practice, extension can be calculated by the  
187 angle between two 3D vectors (Eq. A1) representing two adjacent body segments  
188 (disregarding long-axis rotation). Alternatively, extension can be calculated as a 4D angle  
189 between two quaternions (by the same equation) representing the local reference frames of  
190 two adjacent segments. We later quantify a segment orientation with respect to the global  
191 vertical and horizontal planes (see Section 3.5) rather than with respect to adjacent segments.

192

193

#### 194 *2.4 Unit quaternions and SLERP*

195 Unless otherwise noted, all quaternions in the current work will be unit quaternions and all  
196 angles will be in radians. Unit quaternions have the form

$$\mathbf{q} = \left[ \cos\left(\frac{\theta}{2}\right), \hat{a}_i \sin\left(\frac{\theta}{2}\right), \hat{a}_j \sin\left(\frac{\theta}{2}\right), \hat{a}_k \sin\left(\frac{\theta}{2}\right) \right] \quad (1)$$

197 where  $\hat{\mathbf{a}}$  is a unit vector for the axis of rotation and  $i, j, k$  are its  $x, y, z$  components.  $\theta$  is the  
198 rotation angle about the axis of rotation (in 3D space). Spherical Linear Interpolation

199 (SLERP; Shoemake, 1985) is a method to interpolate intermediate positions between two unit

200 quaternions,  $\mathbf{q1}$  and  $\mathbf{q2}$ . It was developed for efficiently and smoothly computing paths of  
201 moving objects for computer animation. An interpolated quaternion ( $\mathbf{qI}$ ) is calculated as  
202 follows.

$$\mathbf{qI}_{(q1,q2,\tau)} = \frac{\mathbf{q1} \sin[(1 - \tau) \theta] + \mathbf{q2} \sin(\tau \theta)}{\sin(\theta)} \quad (2)$$

203 Where  $\tau$  is relative time between 0 and 1, and  $\theta$  is the 4D angle between  $\mathbf{q1}$  and  $\mathbf{q2}$   
204 (Appendix A). At the endpoints  $\tau=0$  or  $\tau=1$ , the above equation reduces to  $\mathbf{qI} = \mathbf{q1}$  or  $\mathbf{qI} =$   
205  $\mathbf{q2}$ , respectively. For intermediate values of  $\tau$ ,  $\mathbf{qI}$  is the weighted average of  $\mathbf{q1}$  and  $\mathbf{q2}$  in  
206 unit quaternion space.

207

### 208 3.0 FORWARD KINEMATICS ALGORITHM: *limbSLERP*

#### 209 3.1 *Workflow outline*

210 Our workflow is summarized in the following steps: 1) “Quaternionization”: The initial  
211 posture of the left leg and body is “quaternionized” to express each  $i^{th}$  body segment as a  
212 quaternion,  $\mathbf{q}_i$ . 2) Path planning: A target (endpoint) COM location and body orientation is  
213 chosen and expressed as a quaternion,  $\mathbf{q}_{target}$ . Using the quaternionized limb as a starting  
214 point, a path of 3D kinematics is solved analytically using SLERP to derive the segment  
215 paths required to move the body towards the target. 3) Kinematic extrapolation: Using  
216 continuously varying time,  $\tau$ , the quaternionized limb is “SLERPed” towards the target until  
217 the target body pose is reached.

218

#### 219 3.2 *Quaternionization*

220 The orientation of each leg segment (thigh, shank, proximal foot and distal foot) was  
221 expressed as a quaternion relative to the adjacent proximal segment. For the torso, we used  
222 the z-axis [0, 0, 1] as a fixed global reference vector  $\mathbf{v}_{ref}$ . Thus, to describe the local  
223 orientations for each of the 5 body segments, we gathered all quaternions into a vector,  $\mathbf{Q}$ ,

224 containing the 5 quaternions calculated above. Using this convention, a “null” rotation of 0°  
 225 ( $\mathbf{q} = [1, 0, 0, 0]$ ) would result in two adjacent segments aligned end-to-end along their long  
 226 axes. For example, at the null position, (each element of  $\mathbf{Q} = [1, 0, 0, 0]$ ), all segments would  
 227 be aligned end-to-end along the global z-axis (Fig. 1A inset). As another example, for a 90°  
 228 protraction of the left hip (but no rotation at other joints), the femur orientation would point  
 229 to the left side of the body. Since the orientation of the shank is defined with respect to the  
 230 thigh, no relative rotation would be required for the shank ( $\mathbf{q}_{\text{thigh,shank}} = [1, 0, 0, 0]$ ), and  
 231 similarly for the remaining segments (SI Movie 2).

232

233 The starting pose for jumping was quaternionized from the first video frame from an example  
 234 trial collected from a previous data set (Richards et al., 2017). At each time sample we have  
 235  $\mathbf{P}$ , a matrix containing XYZ coordinates for  $ns$  number of segments (i.e.  $ns$  rows X 3  
 236 columns) ordered from proximal to distal. The point of ground contact (i.e. distal foot) is the  
 237 Cartesian origin, XYZ = [0, 0, 0]. Moving from proximal to distal,  $\mathbf{P}$  is converted to  $\mathbf{V}$ , a  
 238 matrix of local segment vectors ( $\mathbf{P} \rightarrow \mathbf{V}$ ):

$$\mathbf{V} = [\mathbf{v}_{i,i+1} \dots, \mathbf{v}_{ns-1,ns}] = [\mathbf{p}_{i+1} - \mathbf{p}_i \dots, \mathbf{p}_{ns-1} - \mathbf{p}_{ns}] \quad (3)$$

239 which become

$$\mathbf{V} = [\mathbf{v}_{\text{torso}}, \mathbf{v}_{\text{thigh}}, \mathbf{v}_{\text{shank}}, \mathbf{v}_{\text{prox.foot}}, \mathbf{v}_{\text{dist.foot}}] \quad (4)$$

240 then each vector is converted to quaternions using the transformation,

$$\mathbf{v1} \xrightarrow{q_{v1,v2}} \mathbf{v2} \quad (5)$$

241 creating local reference frames such that each  $\mathbf{q}_i$  represents the orientation of  $\mathbf{q}_i$  relative to  
 242  $\mathbf{q}_{i-1}$ .

243

244 Finally, the entire left limb (expressed as a set of vectors,  $\mathbf{V}$ ) is quaternionized to  $\mathbf{Q}$ , a vector  
 245 of quaternions representing the postural configuration of the limb at a given time ( $\mathbf{V} \rightarrow \mathbf{Q}$ ).

$$\mathbf{Q} = [\mathbf{q}_{i,i+1} \dots, \mathbf{q}_{ns-1,ns}] = [\mathbf{q}_{\text{torso}} \dots, \mathbf{q}_{\text{prox.foot}}] \quad (6)$$

246 The whole procedure of quaternionization ( $\mathbf{P} \rightarrow \mathbf{Q}$ ) is detailed in Appendix B.

247

### 248 *3.3 Path planning*

249 Firstly, the start pose of the left leg,  $\mathbf{Q}_0$ , was calculated from the first video frame of an  
 250 example jump experimental trial. Secondly, a target end pose (e.g. takeoff)  $\mathbf{Q}_1$  must be given  
 251 to allow SLERP to compute the intermediate kinematics between  $\mathbf{Q}_0$  and  $\mathbf{Q}_1$ . If both start  
 252 and takeoff postures are known, the kinematics of a jump can be SLERPed between the two  
 253 poses to approximately reconstruct experimentally collected data. For the present study we  
 254 wish to produce a hypothetical range of takeoff poses beyond those observed experimentally.  
 255 Thus, the final limb configuration ( $\mathbf{Q}_1$ ) is not known, but rather extrapolated from  $\mathbf{Q}_0$ . To  
 256 extrapolate, we assume that the final configuration ( $\mathbf{Q}_1$ ) lies somewhere between the initial  
 257 ( $\mathbf{Q}_0$ ) and a fully straightened leg. Since extrapolation is essentially a guess based on this  
 258 assumption,  $\mathbf{Q}_1$  must be modified later as explained below.  $\mathbf{Q}_1$  was determined in three steps.  
 259 1) The body heading and orientation at takeoff were specified by an elevation angle,  $\psi$ , in the  
 260 vertical plane and a turn angle,  $\alpha$ , in the horizontal plane. These angles were used to make a  
 261 first guess regarding the composition of  $\mathbf{Q}_1$ . Specifically, varying takeoff targets were chosen  
 262 with respect to a nominal takeoff configuration (takeoff pitch = 33°; yaw = 0) representing an  
 263 exemplar trial from experimental recordings. 2) The relative duration of the interpolated  
 264 kinematics was adjusted to prevent the leg from over-extending to a fully straight posture  
 265 (see below). 3) The kinematics of the opposite (right) leg were solved by mirroring  $\mathbf{Q}_1$  to  
 266 yield  $\mathbf{Q}_{1R}$ . Because turns are asymmetrical, small additional adjustments were made using  
 267 inverse kinematics (Appendix B).

268

269 The calculations are as follows.

270 Path planning step 1: A desired takeoff elevation angle,  $\psi$ , was expressed as a quaternion  
 271 about the body pitch axis ([1, 0, 0]; see Fig. 3A):

$$\mathbf{q}_{pitch} = \left[ \cos\left(\frac{\psi - \pi/2}{2}\right), \sin\left(\frac{\psi - \pi/2}{2}\right), 0, 0 \right] \quad (7)$$

272 Note that an offset angle of  $-\pi/2$  was added in order to define  $\psi$  with respect to horizontal as  
 273 done previously (i.e.  $\psi = 0$  is a horizontal jump; Richards et al., 2017). Importantly,  
 274  $\psi$  determines the orientation of the torso axis at the instant of takeoff, but does not dictate the  
 275 orientation of the centre of mass velocity vector. Regardless, steeper pitch angles will result  
 276 in steeper jumps (i.e. greater jump height at the moment of takeoff). For the present study,  
 277 we use “steep”, “high jump”, “high pitch”, “high elevation” synonymously to refer to a large  
 278  $\psi$  value resulting in greater vertical displacement at takeoff.

279

280 Similarly, a desired turn angle,  $\alpha$  was chosen about the yaw axis ([0, 0, 1]; see Fig. 3B):

$$\mathbf{q}_{yaw} = \left[ \cos\left(\frac{\alpha}{2}\right), 0, 0, \sin\left(\frac{\alpha}{2}\right) \right] \quad (8)$$

281 The two quaternions were then multiplied to give a target rotation composed of a pitch  
 282 followed by a yaw rotation:

$$\mathbf{q}_{target} = \mathbf{q}_{yaw} \otimes \mathbf{q}_{pitch} \quad (9)$$

283 Where  $\otimes$  denotes quaternion multiplication. As a nominal simulation we chose a  
 284 representative jump reaching  $33^\circ$  elevation at takeoff (Richards et al., 2017) with no turning  
 285 ( $\psi = 0.576$ ;  $\alpha = 0$ ).  $\mathbf{Q}_1$  was then defined:

$$\mathbf{Q}_1 = [\mathbf{q}_{torso}, \mathbf{q}_{thigh}, \mathbf{q}_{shank}, \mathbf{q}_{prox.foot}, \mathbf{q}_{dist.foot}] = [\mathbf{q}_{target}, \mathbf{q}_0, \mathbf{q}_0, \mathbf{q}_0, \mathbf{q}_0] \quad (10)$$

286 Where  $\mathbf{q}_0$  is the null rotation resulting in a 0 angle between adjacent segments. Since all  
 287 quaternions describe relative rotations between segments,  $\mathbf{q}_0$  simply means straight  
 288 orientation with respect to the proximal segment. Importantly,  $\mathbf{Q}_1$  is a crude guess which sets  
 289 the direction of motion, but not necessarily the destination configuration. Our theory

290 proposes that as limb segment orientations move towards the null orientation (i.e. straight),  
291 the limb will extend via joint angle trajectories that minimize overall motion (via SLERP).  
292 However, frogs do not fully extend their legs prior to takeoff (Richards et al., 2017).  
293 Accordingly, we allow simulations to move towards  $\mathbf{Q}_1$ , *but are never allowed to reach*  $\mathbf{Q}_1$   
294 by adjusting  $\tau$  (see below). Otherwise, the posture at takeoff would be a fully extended limb  
295 (Fig. 1A, inset).

296

297 Path planning step 2: SLERP was used to interpolate the COM displacement throughout the  
298 jump.

$$\mathbf{QI}_\tau = [\mathbf{qI}_{\text{torso}}, \dots, \mathbf{qI}_{\text{qdist.foot}}]_\tau \quad (11)$$

299 As explained above, we never allowed the interpolation to reach  $\tau = 1$ . Instead, the  
300 interpolation was stopped at time  $\tau'$  when the angle between the torso and thigh segment  
301 reached  $\sim 130^\circ$  to mimic the configuration of the limb just prior to takeoff (Fig. 3B from  
302 Richards et al., 2017).  $\tau'$  was then used for path planning step 3. We define  $\mathbf{Q}_1'$  as the final  
303 configuration at  $\tau = \tau'$ .

304

305 To summarize steps 1 & 2,  $\mathbf{Q}_0$  is first sampled from recorded data - it is the only parameter  
306 known from experimental observation. Then, extrapolation (in unit quaternion space) is used  
307 to guess  $\mathbf{Q}_1$  which is later refined to  $\mathbf{Q}_1'$ . Importantly, neither  $\mathbf{Q}_0$  nor  $\mathbf{Q}_1$  alone contain  
308 information regarding leg kinematics - they only specify configurations at two separate  
309 moments in time to bracket the jump. However, as soon as both  $\mathbf{Q}_0$  and  $\mathbf{Q}_1$  are defined, the  
310 full kinematics of the jump are known (i.e. all trajectories for all body segments) simply by  
311 substituting a time value ( $\tau$ ) into Eq. 11. Along this interpolated path exists a limb  
312 configuration ( $\mathbf{Q}_1'$  at  $\tau = \tau'$ ) that brings the torso midline axis close to the target. Therefore,  
313 the final target pitch and yaw are specified, but the configuration of the leg segments is

314 unknown *a priori*. Thus, our procedure requires extrapolation to guess the final configuration  
315 at a single time point (takeoff), but uses interpolation to derive the motions in between the  
316 two time points. If both  $\mathbf{Q}_0$  and  $\mathbf{Q}_1$  are known, extrapolation would not be needed - one could  
317 skip directly to the interpolation in step 2.

318

### 319 Path planning step 3:

320 Firstly, each  $i^{\text{th}}$  point along the left side of the body was mirrored about the midline body axis  
321 using a reflection matrix

$$\begin{bmatrix} x_r \\ y_r \\ z_r \\ 1 \end{bmatrix}_i = \mathbf{R}_{reflect} \cdot \begin{bmatrix} x \\ y \\ z \\ 1 \end{bmatrix}_i \quad (12)$$

322 where  $\mathbf{R}_{reflect}$  is a 4x4 matrix (Kovács, 2012; Appendix A) and  $x_r, y_r, z_r$  are the reflected XYZ  
323 coordinates. Secondly, for turning simulations, the left and right leg kinematics are  
324 necessarily asymmetric. To solve the asymmetric leg kinematics for turns, the body segment  
325 orientations of the mirrored limb were adjusted using an iterative inverse kinematics (IK;  
326 Appendix C) algorithm. Briefly, IK calculates the minimum changes in joint angles required  
327 to move the limb endpoint to a target. In the current study, the target was the left hip and the  
328 right limb's endpoint was the right hip. At each time point, the right limb (mirrored) was  
329 incrementally moved towards the left leg until the two halves join at the hip. Although not  
330 always necessary (i.e., for symmetric jumps, see Discussion), this IK adjustment was applied  
331 to all simulations.

332

### 333 *3.4 Kinematic interpolation*

334 The previous steps yield a nominal final pose,  $\mathbf{Q}_1'$ , from which we can vary the takeoff  
335 direction of the torso/limbs to simulate jumps of varying steepness and degree of turning.  
336 Using Eqn's 7-10 we modify the nominal simulation by choosing takeoff pitch and yaw

337 angles relative to the nominal condition. Specifically, we alter  $\mathbf{Q}_1$  by modifying its first row  
338 which is  $\mathbf{q}_{\text{target}}$ .

339

340 Because we are using linear interpolation in unit quaternion space, linearly advancing the  
341 time from  $\tau=0$  to  $\tau=1$  produces linear changes in the orientation of each body segment with  
342 time. Thus, a linear increase in  $\tau$  results in constant rotational velocity (Shoemaker, 1985)  
343 which is unrealistic for jumps which require acceleration throughout (Marsh, 1994). This  
344 problem is solved simply by using a nonlinear function for  $\tau$ . Without altering the kinematics  
345 paths derived above, interpolation time can be defined using a function,  $T$ , such that the slope  
346 constantly increases (hence constant increase of velocity). At least a second order curve (e.g.  
347  $T = \tau^2$ ) is needed to guarantee acceleration of the COM throughout. Given that *in vivo*  
348 acceleration patterns are not constant (Roberts & Marsh, 2003), we used a higher order  
349 function based on the computed displacement from our representative nominal jump using  
350 the location of the hip as a proxy for the COM (Richards, 2017), normalized the data by  
351 maximum displacement then fit it to a 4<sup>th</sup> order polynomial to create a function for  $T$ .

$$T(\tau) = a + b * \tau^2 + c * \tau^3 + d * \tau^4 \quad (13)$$

352

353 where  $\tau$  is the adjusted relative time ( $0 < \tau \leq \tau'$ ) with fit coefficients  $a=0.019$ ,  $b= 0.145$ ,  
354  $c=1.383$  and  $d= -0.549$ . To perform kinematic extrapolation,  $T$  is substituted for  $\tau$  in Eq. 11.

355

### 356 *3.5 Analysis of simulated kinematics*

357 All analysis for the present study was performed in Mathematica 10 (Wolfram, Hanborough,  
358 UK). Quaternion interpolation was implemented using two programming loops. The “outer  
359 loop” repeats for  $nt$  time samples and selects a relative time value ( $0 < \tau \leq \tau'$ ) at an arbitrary  
360 sampling interval,  $dt$ , (e.g.  $dt = 0.01$  to give  $nt = 100$  time samples). For each incremental



361 value of  $\tau$ , the “inner loop” repeats for  $ns$  iterations for each body segment. Within this inner  
362 loop,  $\mathbf{qI}$  for each segment (Eq. 11) is calculated and gathered into a vector  $\mathbf{QI}$  at the given  $\tau$ ,  
363 yielding a  $ns \times 4$  matrix of quaternion values. The lowest matrix dimension is always 4  
364 representing the 4 numbers in each quaternion. The completion of both loops yields an  $nt \times$   
365  $ns \times 4$  matrix storing one  $\mathbf{QI}$  for each time point. Finally, quaternionization is performed in  
366 reverse ( $\mathbf{Q} \rightarrow \mathbf{P}$ ) to yield a set of 3D body segment vectors that are assembled end-to-end to  
367 construct the schematic body configuration at time  $\tau$  (Appendix B) analogous to an  
368 experimentally collected 3D kinematics data set.

369

370 Following kinematic interpolation, simulated limb segment kinematics were analysed using a  
371 similar approach as on experimental data (see Richards et al. 2017). Two metrics were used  
372 to quantify kinematic differences between varied behaviours: 4D angles and limb segment  
373 orientations in polar coordinates. 4D angles were used as a scalar measure of “rotational  
374 distance” between two orientations. By analogy, the distance travelled along the surface of a  
375 sphere (in 3D) can be measured as a scalar angle between any two positions on the surface.  
376 One can similarly measure a scalar distance between any two 4D unit vectors whose tips lie  
377 on the surface of a 4D sphere. Since all 3D rotations, when expressed as unit quaternions,  
378 reside on the surface of a 4D sphere, any two orientations are represented by two points on  
379 the 4D sphere. Similar to the 3D analogy, scalar “distance” between two orientations can be  
380 represented by the 4D angle between them (Eq. A1). Whereas 4D angles indicate rotational  
381 displacement, limb segment polar angles describe orientation by treating each segment as the  
382 radius about an imaginary sphere centred at its joint of origin. For example, one can imagine  
383 the hip joint as the centre of a sphere whose radius is the femur. The polar orientation of the  
384 femur is then described by two angles: 1) A protraction-retraction angle in the horizontal  
385 plane and 2) an abduction-adduction angle relative to the vertical ( $z$ ) axis (Fig. 3):

$$retraction\ angle(\theta) = \tan^{-1}\left(\frac{y}{x}\right) - \frac{\pi}{2} \quad (14)$$

386

$$adduction\ angle(\phi) = \frac{\pi}{2} - \cos^{-1}\left(\frac{z}{segment\ length}\right) \quad (15)$$

387

388 Where  $x$ ,  $y$  and  $z$  are Cartesian coordinates of a given segment vector (e.g.  $\mathbf{V}_{thigh}$  from Eq. 4  
 389 for calculating thigh retraction-protraction and abduction-adduction). Note that the  $\pi/2$   
 390 offsets are for convenience such that the retraction angle sweeps rearward from the torso  
 391 midline and adduction sweeps downward from the horizontal plane.

392

## 393 4.0 RESULTS

### 394 4.1 Simulated versus observed jump kinematics

395 Compared to data collected experimentally (Richards et al., 2017), jump kinematics predicted  
 396 by *limbSLERP* produced similar patterns when expressed as 4D angles (Fig. 4). Using the  
 397 current convention of “quaternionization”, a 4D angle of 0 indicates that a segment is at the  
 398 “zero” position; i.e. the segment’s long axis is aligned with that of the proximal segment.  
 399 Note this zero position is a theoretical extreme which is not anatomically possible (see  
 400 below). In both experimental and simulated jumps, 4D angles decreased through time  
 401 indicating that the entire limb extends (i.e. “straightens”) as each segment’s orientation  
 402 converges towards a common orientation. For all segments except the proximal foot,  
 403 SLERPed simulations followed trajectories within the variation of trajectories observed *in*  
 404 *vivo*. The proximal foot showed the same downward trend, but at consistently lower values  
 405 than natural frog kinematics indicating that the proximal foot remained slightly “straighter”  
 406 with respect to the femur compared to experimental observations (Fig. 4C). Discrepancies  
 407 between simulation and real kinematics perhaps result from using external skin markers  
 408 (Richards et al., 2017). Despite negligible relative motion between skin and underlying joints

409 (see above), a pair of external markers per segment is not sufficient to reconstruct as a proxy  
410 for internal bone motion which our theory simulates. Additionally, our model does not  
411 account for small translations possibly occurring at the knee joint (Kargo et al., 2002) which  
412 might cause the observed angular offset (see Discussion).

413

#### 414 *4.2 Simulated modulation of jump performance*

415 From an identical starting configuration based on experimentally collected data (Richards et  
416 al., 2017),  $\mathbf{q}_{\text{target}}$  was varied to modulate both the vertical steepness of jump angle (pitch) and  
417 the turning angle (yaw). Three contrasting takeoff targets were chosen: Nominal takeoff  
418 configuration (takeoff pitch = 33°; yaw = 0), a steep jump (pitch = 60°; yaw = 0) and a left  
419 turn (pitch = 15°; yaw = -18°). Note that we used a shallow pitch to allow turns to occur  
420 mainly in the horizontal plane. This allowed us to reduce confounding effects of “banking  
421 turns” (i.e. rolling about the torso axis) to better isolate the kinematic mechanism of turning.  
422 Animations from simulations show smooth motion of the body segments extending to carry  
423 the torso towards the target orientation (Fig. 5; SI Movie 3). For the three example cases,  
424 kinematics were qualitatively similar; extension of the main driving joints can be seen clearly  
425 in top view for the hip and knee (Fig 5, second row) and in side view for the ankle (Fig 5,  
426 third row). For the nominal simulation, the hip, knee, ankle and TMT joints underwent a net  
427 extension (max angle – min angle) of 80, 93, 80 and 37°, respectively. Notably, magnitudes  
428 of joint extension remained consistent across different jumps; total excursion in terms of 3D  
429 and 4D angles (see Section 3) varied within only a few degrees difference from the nominal  
430 extension values (nominal values  $\pm 1$  to 5° versus 1 to 14° for varying pitch versus yaw). The  
431 consistent amount of extension across simulated behaviours suggests that limb segment  
432 orientation governs jump trajectory (rather than differential magnitudes of joint extension).

433 Thus, the present analysis focuses on how varying degrees of upward, forward and medio-  
434 lateral motion of limb segments can modulate the jump path.

435

436 To observe the relative contribution of upward/downward versus forward/rearward versus  
437 medial/lateral segment motions for the example cases above, each segment was decomposed  
438 into polar coordinates to quantify orientations in terms of retraction angle (caudal-directed  
439 rotation to push the body forward) or adduction angle (downward rotation to push the body  
440 upward; Fig. 6). During straight jumps and turns, all leg segments (thigh, shank, proximal  
441 and distal foot) retracted and adducted throughout jumps (except the shank which adducted,  
442 but remained at a fixed retraction angle; Fig. 6A&B).

443

#### 444 *4.3 Theoretical kinematics comparing a nominal jump, a steep jump and a left turn*

445 For straight jumps, protraction-retraction and abduction-adduction kinematics were nearly  
446 symmetrical between left and right legs. However, SLERP introduced slight lateral drift of  
447 the body midline axis (Fig. 5B&F) which caused a small asymmetry in segment kinematics  
448 (see Discussion). Regardless for steep jumps, both left and right legs showed a reduction in  
449 the extent of retraction, particularly for the thigh and proximal foot which were reduced by  
450  $\sim 50^\circ$  and  $\sim 90^\circ$ , respectively, with the shank retraction nearly constant. In contrast, adduction  
451 for the thigh and shank segments increased by  $\sim 14^\circ$  and  $\sim 26^\circ$ , respectively, but decreased by  
452  $\sim 15^\circ$  for the proximal foot.

453

454 During simulated turns, left and right legs extended with subtle asymmetries (Fig. 6A vs. B;  
455 Fig. 7). Unexpectedly, there was no dramatic right leg bias in either retraction or adduction  
456 angles. There were three notable subtle asymmetries. Firstly, the thigh and proximal foot  
457 segments retracted to a greater extent in the left leg. Secondly, the left shank protracted

458 slightly (“pulling” the limb towards the inside of the turn) whereas the right shank retracted  
459 slightly (“pushing” the limb leftwards; Fig. 7A). Thirdly, the right proximal foot adducted  
460 earlier and to a greater extent than the left (Fig. 7B).

461

#### 462 *4.4 Theoretical kinematics of modulating jump height and turn angle*

463 To better understand how kinematics are modulated, we simulated a hypothetical range of  
464 takeoff pitch angles (with no turning) and a range of turning angles (at fixed takeoff pitch).

465 Over the range of pitch angles, retraction of the thigh and proximal foot segments decreased  
466 dramatically  $\sim 50^\circ$  and  $130^\circ$ , respectively, from the shallowest to the steepest jumps (Fig. 8A).

467 In contrast, adduction excursion for the thigh and shank segments increased  $\sim 50^\circ$  over the  
468 range of pitch angles revealing a forward kinematic mechanism of exchanging retraction for  
469 adduction to increase takeoff pitch.

470

471 For turning we swept the left leg turn angle from positive (right turn) to negative (left turn) to  
472 cause a functional shift from the outside leg (“pushing” the leg toward the opposite side) to  
473 the inside leg (“pulling” the leg into the turn). Surprisingly, thigh and proximal foot  
474 retraction increased as turn angle decreased (i.e. “pushed” more on the inside of the turn).

475 The shank kinematics, however, shifted  $\sim 30^\circ$  from retraction (left leg turning right) to  
476 protraction (left leg turning left), with negligible protraction or retraction at  $0^\circ$  turn angle  
477 (Fig. 8B).

478

#### 479 *4.5 The influence of jump preparation angle (initial pitch angle)*

480 To test the influence of the ‘jump preparation angle’ (Wang et al., 2014), we chose a fixed  
481 target takeoff pitch of  $60^\circ$  (yaw = 0) whilst varying the initial pitch angle of the body segment  
482 (at  $\tau=0$ ). Across a range of initial pitch angles of  $0^\circ$  (horizontal) to  $\sim 75^\circ$  (nearly vertical), the

483 paths of the segment rotations changed, causing shifts in the final direction of centre of mass  
484 (COM) velocity at takeoff (Fig. 9). At low preparation angles ( $< 45^\circ$ ), the simulation would  
485 leave the ground moving upwards and backwards, despite a body axis heading of  $60^\circ$   
486 pointing upwards/forwards. In contrast, preparation angles above  $45^\circ$  caused forward/upward  
487 motion at takeoff. Over the range of initial angles, the COM takeoff velocity angle decreased  
488 sharply, reaching an optimum of  $60^\circ$  at a preparation angle of  $\sim 47.5^\circ$  where the body would  
489 continue in the target direction.

490

## 491 5.0 DISCUSSION

492

### 493 *5.1 limbSLERP predicts jump kinematics*

494 The goal of the present study was threefold: Firstly, we aimed to create a  
495 mathematical/computational method to predict limb motion from simple geometric  
496 information such as limb segment proportions and their connections. Our use of SLERP is  
497 novel because, to our knowledge, it provides one of the simplest sets of mathematical rules  
498 that predict realistic limb motion in the absence of detailed physical, anatomical and  
499 physiological constraints. It is intended as a first step to provide predictions and insights to  
500 assist the development of more rigorous dynamics analyses (forward dynamics; inverse  
501 dynamics; musculoskeletal simulation) to follow. Secondly, we intended to establish a  
502 theoretical framework to simulate frog hind limb kinematics over a range of jumping  
503 performance. Our aim was not to faithfully reproduce or fit experimental data. Rather, we  
504 created a template model with minimal mathematical constraints capturing the essential traits  
505 of a frog jump (Fig. 4) to probe for insights into the coordination of movement. Thirdly, we  
506 sought to tease apart whether pre-jump posture versus dynamic modulation of leg kinematics  
507 are most crucial in steering the COM forwards, upwards or laterally. During jumps, frogs  
508 naturally vary their takeoff pitch angle (Kargo & Rome, 2002; Wang et al., 2014; Richards et

509 al., 2017; Porro et al., 2017) as well as their turn angle (C.T. Richards & L. B. Porro,  
510 unpublished observations). Whilst controlling for initial pre-jump posture, we tested whether  
511 a set of simple coordination rules could reproduce frogs' natural pitching and turning  
512 behaviours.

513 The success of *limbSLERP* for simulating realistic kinematics, despite its simple assumptions  
514 and neglect of dynamics (i.e. forces), stems from two principles we hypothesise to be crucial  
515 for frog jumping. Firstly, based on experimentally observed behaviour (Richards et al., 2017;  
516 SI Movie 1), the joint axes of rotation converge prior to takeoff. Secondly, the limb  
517 kinematics between the start and end configurations result from the minimisation of segment  
518 rotation. Theoretically, there exist infinite paths along which the limb segments could travel  
519 between start and end postures. However, from a kinematics perspective (without knowing  
520 the dynamics) our theory proposes that the most sensible path is that which minimises  
521 motion. Supporting our first hypothesis (H1) the above coordination principles approximate  
522 natural kinematics (Fig. 4), providing evidence that frogs may coordinate their limb  
523 movements by converging rotation axes and economising motion.

524

## 525 *5.2 Theoretical evidence for how frogs modulate jump height and turn angle*

526 Current findings suggest that changes in leg segment adduction had the strongest direct effect  
527 on takeoff height supporting hypothesis H2. As expected, greater downward rotation of  
528 segments, particularly the thigh and shank, “pushed” the body upwards to cause steeper  
529 jumps for a given functional leg length (Fig. 6C & 8A) similar to experimentally observed  
530 jumps (Richards et al., 2017). Although simulations predicted little adduction contribution  
531 from the proximal foot, this does not imply that the ankle joint is inactive. Inverse dynamics  
532 analysis indicates that increased torque from the ankle and hip drive steeper jumps (Porro et

533 al., 2017). Thus, increased thigh and shank adduction are likely powered by torques at the  
534 hip and ankle, respectively. In exchange for greater adduction, retraction decreased (Fig. 6A  
535 & 8A) to translate the body upwards rather than forwards (SI Movie 3). Notably, altering the  
536 pre-jump body pitch did not influence the ability to reach the target jump orientation. Rather,  
537 adjustment of initial body inclination enabled the simulation to travel in the direction that the  
538 body was pointed (i.e. aligning the body orientation with heading; Fig. 9). We speculate that  
539 inclining the pre-jump body angle not only is an indicator of fore limb push off (Wang et al.,  
540 2014), but also is a mechanism to aid neuro-muscular control of takeoff velocity. This is  
541 because appropriate inclination of the initial body posture allows the limb to travel in the  
542 direction of its body simply by straightening the limb. Using an analogy to reinforce this  
543 point, we imagine a toy robot which must be programmed with joint angle trajectories. If the  
544 hobbyist wishes to program a steeper jump, doing so with a higher initial body pitch will  
545 simplify the programming of the hindlimb kinematics. For the most extreme takeoff  
546 steepnesses, the cumulative rotations of all segments caused the COM to accelerate  
547 backwards (Fig. 9). Although such strong backwards motion is unlikely *in vivo*, real frogs do  
548 generate short periods of rearwards force as they shift their weight and pitch their body  
549 rearwards during the steepest jumps (Porro et al., 2017) which could theoretically be  
550 corrected by inclining the body prior to launch.

551

552 The joint kinematics for turns did not behave as expected. Instead of greater retraction on the  
553 outside (right) limb segments, thigh and proximal foot retraction increased more on the left  
554 (Fig. 6A), becoming greater with sharper turns (Fig. 8B). Although this increased retraction  
555 is counterintuitive, further inspection reveals that it likely has little impact on overall limb  
556 kinematics due to shifts in adduction. The greatest change in retraction occurs at the  
557 proximal foot; however, this segment also undergoes greater adduction approaching  $\sim 90^\circ$



558 (vertically downwards) which nearly cancels any impact of retraction. Using a globe analogy,  
559 adduction/abduction is analogous to moving north/south on a globe whereas protraction-  
560 retraction refers to east/west. When a segment is adducted to  $90^\circ$  (i.e. at the South Pole by  
561 our definition) protraction-retraction has no effect on segment orientation in the same sense  
562 that one cannot move east or west whilst on a pole. Given that the most visible changes have  
563 little kinematic effect, what is the key asymmetry that causes turns? For the same reason that  
564 the proximal foot becomes ineffective at retraction, the shank has the strongest effect due to  
565 its horizontal orientation which causes the greatest XY displacement for any given  
566 protraction/retraction. Furthermore the shank of our model frog, *Kassina maculata*, is the  
567 longest leg segment (Fig. 1A) and hence exerts greater displacement as it rotates. In support  
568 of hypothesis H3, the orientation of the shank (in the XY plane) remained nearly stationary  
569 across all jump conditions (Fig. 6A&B; SI Movie 3) except for turns in which left shank  
570 retraction switched to protraction causing differential rotation “pulling” the leg backward on  
571 the left whilst “pushing” forward on the right. For the reasons above, our model predicts that  
572 such a subtle shift (Fig. 7A) is sufficient to steer the frog.

573

574 Our above results, though intuitive to understand, do not reflect the only possible mechanism  
575 for how frogs kinematically modulate jump direction. Because of the high number of  
576 degrees-of-freedom of frog hindlimb (Kargo & Rome, 2002) there hypothetically exist  
577 multiple possible solutions to how a frog might differentially rotate its segments to increase  
578 jump steepness and turn. Our current method based on minimal rotation offers a sensible  
579 starting prediction and outperforms Euler angle-based inverse kinematics (IK; e.g. Bus 2004).  
580 When attempting to derive a jump trajectory for the left leg using IK, the simulated  
581 kinematics diverged towards an alternative, but unnatural extension of the limb. Specifically,

582 the hip hyperextended while the ankle compensated by migrating medially to arrive on the  
583 opposite (right) side of the frog (Supplementary information, Fig. S1).

584

### 585 *5.3 Theoretical evidence for a "jump kinematics template"*

586 Despite variation in performance, do all frog jumps share similar underlying biomechanical  
587 traits, regardless of the jump direction? Although the final answer can only be found using  
588 dynamics analysis, the present theoretical kinematics approach gives some insight. In  
589 particular, if the underlying limb segment dynamics (torques and accelerations) follow a  
590 qualitatively similar pattern across jump performance, we expect kinematics should also be  
591 similar (and vice versa). We propose that these similarities constitute a theoretical  
592 "kinematics template" which can be morphed to vary jump performance within *Kassina* (and  
593 potentially other morphologically similar frog species) whilst maintaining the fundamental  
594 characteristics of a jump. A common pattern can be distilled from all observed jumps both  
595 simulated (present study) and *in vivo* (Richards et al. 2017) in three main elements: 1) The  
596 thigh and proximal foot segments rotate rearwards (retract) with negligible shank retraction.  
597 This is most clearly seen in top view where the shank orientation appears to remain fixed  
598 whilst the neighbouring segments retract (SI Movie 3). 2) The thigh, shank and proximal  
599 foot segments adduct throughout the jumps. 3) The joint axes of rotation converge  
600 throughout the jump (SI Movie 1).

### 601 *5.4 Limitations of the present kinematics approach*

602 The foremost limitation of any kinematics analysis, including the present study, is the neglect  
603 of dynamics. Simulations do not account for limb masses or moments of inertia nor do they  
604 consider muscle force and power properties which are known to limit frog jumping ability  
605 (Galantis and Woledge, 2003; Lutz and Rome, 1994; Peplowski and Marsh, 1997; Roberts

606 and Marsh, 2003). For example, the present analysis cannot address speed effects which  
607 could alter kinematics due to higher joint torques and greater ground reaction forces to drive  
608 farther jumps. Regardless, as we explain below, our kinematics approach gives important  
609 insights that could be overlooked with more detailed dynamics modelling. A second  
610 limitation is our lack of mathematical constraints other than the requirement that all limb  
611 segments connect end-to-end and that joints only rotate. We avoided additional algebraic  
612 approaches to imposing additional motion constraints (e.g. obstacle avoidance; Murray et al.,  
613 1994) because they are algebraically messy and, more importantly, to avoid *a priori* biases  
614 from prior knowledge of frog behaviour. Due to our lack of constraints, the torso segment  
615 drifted medially to cross the body midline which does not occur naturally (Fig. 5B&F; SI  
616 Movie 3). In practice, one can easily perform minor *post hoc* corrections using inverse  
617 kinematics to correct left-right drift of the hip joint. For the current study, we allowed the left  
618 leg to drift slightly then used inverse kinematics for the right leg to join the right hip with the  
619 left (see Appendix C). This drift correction caused the small left-right asymmetry during  
620 straight jumps (Fig. 6 A&B). Despite this small issue, lack of additional constraints  
621 strengthens confidence in our model which approximates natural behaviour without  
622 “knowing” rotational limits of joints or that left segments must remain on the left side.  
623 Finally, the current implementation of quaternion interpolation assumes that frog joints are  
624 “ball joints” (i.e. no translation) which is not representative for all joints (e.g. the knee in  
625 ranid frogs; Kargo et al., 2002). This oversimplification is possibly the cause for the slight  
626 downward offset of the proximal foot segment compared to experimental results (Fig. 4C).  
627 However, given that the above discrepancy was small and that the remaining leg segment  
628 patterns matched those predicted, *limbSLERP* is a simple and powerful starting point for  
629 exploring 3D limb kinematics which could be further developed in the future by adding  
630 translations using dual quaternions (Kavan et al., 2008).

631

632 *5.5 Summary, interpretation and broader context*

633 The present study had three main findings. Firstly, frogs straighten their legs by moving their  
634 leg segments along paths of minimal rotation (i.e. paths on the unit quaternion sphere)  
635 between the pre-jump posture and a fully straightened limb. Moreover, the target orientation  
636 of the fully straightened limb determines final jump direction. Secondly, limb segment  
637 adduction is the key determinant of jump steepness, particularly the thigh and shank. In light  
638 of prior work showing the importance of forelimb push-off (Wang et al. 2014) versus  
639 hindlimb adduction (Richards et al., 2017), we sought evidence to determine which  
640 alternative mechanism is the most important. Although forelimb push-off does indeed  
641 contribute vertical force (Wang et al., 2014), our findings suggest additionally that hindlimb  
642 adduction is necessary and sufficient to produce steeper jumps. Instead of playing a direct  
643 mechanical role, pre-jump “preparation angle” is perhaps important for control by  
644 influencing whether the body travels in the direction that it is initially oriented. Thirdly,  
645 turning is caused by a subtle switch from shank retraction to protraction on the inside leg of  
646 the turn. Overall, our theoretical model predicts that jumps of different direction share the  
647 same fundamental kinematic mechanism whereby the thigh and shank adduct to drive limb  
648 elevation whereas the thigh and proximal foot retract to thrust the body forward. Among all  
649 of the kinematic shifts observed in simulations, the orientation of the shank segment acted as  
650 the principle steering mechanism (due to its length and straight orientation) to modulate jump  
651 height as well as turn angle.

652

653 Beyond qualitative description, we aim that our kinematics theory provides precise  
654 hypotheses for further testing. Until future dynamics analyses are performed, we cannot fully  
655 claim that our theory has direct bearing on biomechanics beyond those explained by simple

656 geometry. However, we argue our approach is a critical first step in understanding the  
657 various limits of the locomotor system. We propose a multi-step approach focusing first on  
658 kinematics, then rigid body dynamics followed by musculoskeletal dynamics. Kinematics  
659 analysis explores the outermost limit to behaviour which is bounded by geometry (i.e.  
660 segment lengths and their anatomical relationships). Within the scope of geometrically  
661 permissible motions, there is a subset which is physically possible, as determined by rigid  
662 body dynamics analysis. For example, how far could a frog jump given knowledge of its  
663 limb inertial properties, ground contact properties and above kinematic limitations?  
664 Furthermore, within the scope of physically possible, there are the physiological limits of  
665 muscle force, velocity and power (Josephson, 1999) as well as limits to bone stress  
666 (Biewener, 1989) and soft tissue structure of joints to influence passive forces and range of  
667 motion (Kargo et al., 2002). Finally, there is the smallest subset encompassing what animals  
668 are willing to do behaviourally (particularly within experimental setups; Astley et al., 2013).  
669 Because researchers cannot dictate behaviour, we reiterate the value of a modelling approach  
670 where key aspects (such as initial posture) can be held constant to better highlight causal  
671 relationships and underlying mechanisms not detected with traditional experimentation. We  
672 argue that one cannot fully understand experimentally observed behaviour until we are able  
673 to explain the “lower level” limitations of the system components. Moreover, we propose  
674 that experimental approaches containing individual variation and measurement noise may not  
675 be sufficiently sensitive to discern subtle behavioural shifts such as those presently observed  
676 during turns.

#### 677 *5.6 Applications and future work*

678 In addition to the biomechanical implications of our approach, we hypothesise that our  
679 theoretical kinematics template is a basic coordination strategy for frog jumping. Despite the  
680 mathematical abstraction of quaternions and 4D hyperspheres, the theory has a physical basis

681 which we propose can be exploited physiologically. A nervous system need not perform  
682 calculations on quaternions; perhaps all that is needed is to generate torques which straighten  
683 the limb whilst minimising segment rotation (perhaps by minimising muscle shortening).  
684 Future forward dynamics modelling could be used to test whether *limbSLERP* could be used  
685 as a high level controller to generate physiologically realistic torques and ground reaction  
686 forces either for robotics approaches or further studies in musculoskeletal dynamics.  
687 Furthermore *limbSLERP* is a simple analytical approach that can compute entire trajectories  
688 for all limb segments extremely fast, making it potentially useful for control of 3D robotic  
689 limbs.

690 An additional application is to supplement data collected from X-ray Reconstruction of  
691 Moving Morphology (XROMM) experiments (Brainerd et al., 2010). In particular, the  
692 rotation of bones about their long axis (long axis rotation; LAR) can be an important feature  
693 of kinematics (Kambic et al., 2014; Rubenson et al., 2007), yet its measurement can be  
694 difficult because it requires at least three non co-linear implanted markers to be visible on a  
695 single structure. Impressively, LAR measurements have been performed on small animals  
696 such as frogs (Astley &, Roberts, 2014). However because the markers in frog bones are tiny  
697 and move at high speeds, some trials may be lost due to failure of image processing software  
698 to track certain markers. In such cases like frog jumping where marker visibility may be  
699 intermittent, *limbSLERP* could be used to supplement frog XROMM data either to fill in the  
700 gaps or, perhaps even to predict LAR in the absence of a third marker (given that there is  
701 some information about a bone's initial long axis orientation).

702 Finally, the most important application to our theoretical approach is to provide a simple tool  
703 for evolutionary morphologists. Our present study did not apply our theoretical kinematics  
704 template to other species with different limb segment length proportions. For example, if the  
705 shank were relatively shorter (as in some burrowing frog taxa; Emerson, 1976), would frogs

706 rely on other segments to be the main drivers for increased jump height and turning? In  
707 contrast, would longer relative shank lengths (as exhibited by tree frogs and terrestrial  
708 jumping taxa; Emerson, 1982) increase potential range of jumping performance? The current  
709 technique could be used to fully map the space of feasible/optimal initial postures and  
710 segment kinematics given the diversity of limb proportions among frog species. Additionally,  
711 *limbSLERP*, with its simple coordination rules, provides an objective and replicable way to  
712 simulate locomotion in extinct anuran species without relying on taxon-specific experimental  
713 kinematics data (derived from species specialized in a particular locomotor mode). More  
714 broadly, our kinematics approach combined with subsequent dynamics analyses can be used  
715 to generate and test precise hypotheses relating evolutionary changes in skeletal structure  
716 (e.g. Emerson 1982; Reilly & Jorgensen 2011) to changes in limb function.

717

718

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726

#### COMPETING INTERESTS

727 There were no competing interests in the current study.

728

729

#### AUTHOR CONTRIBUTIONS

730 C. Richards developed the theory, performed the kinematics analysis and co-wrote the  
731 manuscript. L. Porro collected, processed and analysed the experimental data used to verify  
732 the model. L. Porro also co-wrote and edited the manuscript.

733

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836

## 837 FIGURE LEGENDS

838 FIG1.

839 Schematic view of a frog jump. (A) Inset showing top view (XY plane) with body segments  
840 and joints in a natural configuration. Dashed lines extending back from the centre of mass  
841 (red circle) show leg segments in the "zero" position using null quaternions. (B) An exemplar  
842 jump in top view and (C) rear view showing the first (gray) and final (black) frames of the  
843 left leg. Dashed lines show axes of rotation for the hip (black), knee (red) and ankle (blue)  
844 highlighting how their orientations change in the direction of the arrows to align throughout

845 the jump (see also SI Movie 1) (D) Initial limb (gray) and end-jump configurations (black) on  
846 a floor (grey square) representing the global XY plane. Local reference frames are shown  
847 with the local X (red) and Z (cyan dashed) in all frames (global frame, snout frame, torso  
848 frame, hip frame) to illustrate reference frame transformation using "quaternionization".  
849 Local Z-axes, by definition, align along each body segment. Y-axes as well as ankle and  
850 TMT reference frames have been omitted for clarity. (E) End-jump, rear view, (F) side view  
851 and (G) top view.

852 FIG2.

853 Calculating a smooth path of minimal rotation between two orientations is challenging in 3D,  
854 but trivial in 4D. (A) A Cartesian reference frame (Black) is rotated to a new orientation  
855 (Red) along a path determined by Spherical Linear Interpolation (SLERP) to calculate the  
856 minimum rotation required. Paths of the X axis (open circles) and the Z axis (closed circles)  
857 illustrate the smooth curved motion between the two orientations. (B) The rotation shown in  
858 (A) is parameterized in 3D space (Euler angles) versus 4D space (quaternions, dashed). Euler  
859 angles about X (Red) followed by Y (Blue) then Z (Green) axes are unpredictable and  
860 nonlinear and thus are difficult to extrapolate meaningfully. However, quaternion  
861 displacement via SLERP is linear (i.e. a great arc on the hypersphere surface), making  
862 extrapolation trivial.

863 FIG3.

864 Segment angle definitions. Schematic of left leg in (A) side view showing the target pitch  
865 angle with respect to horizontal (B) top view showing the target yaw angle and the retraction  
866 angle of the thigh with respect to the global y axis and (C) rear view showing the adduction  
867 angle with respect to horizontal. The black circle marks the hip. Note that segment angles  
868 (pitch, yaw, retraction, adduction) are all calculated in the global reference frame (as opposed  
869 to segment quaternions which are in local reference frames; see text).

870 FIG4.

871 Experimental versus simulated frog jump kinematics. Scalar angles (in 4D) were calculated  
872 from quaternion unit vectors for body segments of (A) the thigh, (B) shank, (C) proximal  
873 foot and (D) distal foot compared to the unrotated "ground" reference frame. Grey dots are  
874 experimentally collected data points from 24 intermediate-height jumps (Richards et al.,  
875 2017). Red lines indicate kinematics simulated using SLERP. In the present convention, 4D  
876 angles of 0 would indicate that segments are at the "zero" (null) position (i.e. leg segments  
877 straightened caudally from the hip; Fig. 1 inset). Large angle values indicate a large  
878 deviation from the null configuration of the limb. In a frog jump, the limb joints extend to  
879 straighten the limb. Thus, the angle values decrease rather than increase through time as the  
880 entire limb extends.

881 FIG5.

882 Simulated kinematics animations. Jump animations are shown for (A-D) nominal takeoff  
883 kinematics, (E-H) a steep jump, (I-L) a left turn shown in angled view (row 1), top view (row  
884 2), side view (row 3) and rear view (row 4). A fixed subset of evenly-spaced animation  
885 frames are shown in each view. For the non-turning jumps (A-H) only the left leg is shown  
886 (Black), whereas both left and right (Red) legs are shown for the turning simulation (I-L).  
887 Note that initial configurations (Bold lines) are identical for each condition. The x and y axes  
888 of the global coordinate frame are shown in black and the z-axis is red.

889 FIG6.

890 Simulated limb segment protraction-retraction and abduction-adduction angles for varying  
891 jump steepness and turning. Traces are for thigh (black), shank (red) and proximal foot (blue)  
892 for (A,B) retraction in the cranio-caudal direction and (C,D) adduction in the dorso-ventral  
893 direction shown for both left and right legs during a nominal jump (solid), a steep jump

894 (dashed) and a turn (dotted-dashed) as in Fig. 5. The dashed line (A,B) represents a line  
895 drawn posterior from the hip joint from which protraction-retraction angles were referenced.  
896 Trending towards the line denotes segment caudal rotation (retraction) to push the body  
897 forward. The x-axis (C,D) represents the horizontal axis. Downward slopes indicate  
898 downward rotation (adduction) to push the body upwards. Kinematics of the distal foot are  
899 similar to the pattern for the proximal foot and therefore have been omitted for clarity.

900 FIG7.

901 Left versus right limb kinematics for a left turn. Data traces are from Fig. 6, but rearranged to  
902 highlight left-right asymmetries. Traces are for left leg (solid) and right leg (dashed) using the  
903 same colours as in Fig. 6. (A) Retraction in the cranio-caudal direction and (B) adduction in  
904 the dorso-ventral direction. Note in (A) how the right shank trends downward towards the  
905 dashed line indicating retraction to push the limb forward versus the left shank which trends  
906 upwards (protracts) to push the limb backwards on the inside of the turn. Kinematics of the  
907 distal foot are similar to the pattern for the proximal foot and therefore have been omitted for  
908 clarity.

909 FIG8.

910 Left leg angular excursion for varying jump steepness and turning. (A) Varying pitch angle  
911 relative to horizontal (yaw = 0). Total retraction excursion (max retraction angle - min  
912 retraction angle [closed circles]) and adduction excursion (max adduction angle - min  
913 adduction angle [open circles]) of the thigh (black), shank (red) and proximal foot (blue).  
914 Each data point represents a single simulation beginning from the nominal initial limb  
915 configuration and ending at the specified target angle. Note that increasing jump steepness  
916 requires increased thigh and shank adduction while retraction decreases. (B) Varying turn  
917 angle (constant takeoff pitch = ~8 deg). Turns range from left (negative values) to right

918 (positive). Negative excursion values indicate protraction. Note that unlike varying pitch,  
919 changes in the magnitude of retraction modulate turn angle.

920 FIG9.

921 The effect of initial angle on jump trajectory and centre of mass (COM) takeoff velocity.

922 Each point represents a single simulation whose entire kinematic path is influenced by initial

923 angle. The initial pitch angle of the torso segment was varied from 0 (horizontal) to nearly

924 vertical, leaving the leg segments unchanged. The final pitch of the body axis was held at 60

925 degrees for all simulations. The dashed black lines represent the optimal initial pitch angle

926 which allows the COM takeoff velocity to align with the takeoff body orientation (i.e. the

927 frog COM will travel in the appropriate direction). Stick figure animations for minimum,

928 maximum and optimum initial angles show the initial posture (bold) and subsequent

929 animation frames (gray). Red arrows indicate the direction of the takeoff velocity vector.

930 Note that as the initial pitch angle increases, takeoff velocity direction shifts from  $>90$

931 (jumping upwards and backwards) to  $<90$  (upwards and forwards).

932 FIG S1

933 Euler angle-based inverse kinematics (IK) versus quaternion-based kinematics (limbSLERP).

934 Kinematics of the left leg and body are shown in grey with the final posture at takeoff in red.

935 Experimental observations (left) are compared with limbSLERP (middle) versus IK (right).

936 Global XYZ axes are shown with Z in red. limbSLERP predicts jump behaviour reasonably

937 well, however IK is unreliable for frog jumps. Note the unnatural hyperextension of the hip

938 (black arrow) and, more problematically, the migration of the ankle joint to the opposite side

939 of the frog (red arrow). We also note that in our implementation using Mathematica on OSX,

940 IK was  $\sim 35x$  slower to compute.

941

942 SI Movie 1. Animation of experimentally observed kinematics example trial (used as the  
943 basis for the nominal simulation). The left leg rear view is shown on the left and top view on  
944 the right. On the rear view, the instantaneous joint axes of rotation (dashed) have been  
945 estimated to be the normal vector to the adjoining segments. Axes are shown for the hip  
946 (black), knee (red) and ankle (blue). The centre of mass (hip) is shown in red. The global  
947 reference frame is shown in black with the Z axis bold and the X axis dashed. Note that the  
948 distal foot segment has been mathematically fixed to the ground.

949 SI Movie 2. Screen recording of interactive forward kinematics computations. The values  
950 making the matrix  $\mathbf{Q}_0$  are shown with each row a quaternion representing a body segment.  
951 This illustrates how changes in proximal segment orientations cause all distal limbs to follow.  
952 For example, changing the torso pitch angle (while leaving the remaining quaternions as null  
953 values) transforms  $\mathbf{Q}_{0\text{-torso}}$ , moving the torso along with the remaining segments. The centre  
954 of mass (hip) is the large circle. For this demo, the joint angles are protraction-retraction  
955 only.

956 SI Movie 3. A demo animation of the three exemplar simulations (nominal followed by a  
957 steep jump then a turn). Limb kinematics traces are shown for the hip (black), knee (red) and  
958 ankle (blue). Note how the kinematics traces instantly update for any changes in target  
959 orientation. This is because the trajectories are known for all values of  $t$  as long as  $\mathbf{Q}_0$  and  $\mathbf{Q}_1$   
960 are known.

961

List of Abbreviations Used (bold values are vectors or matrices)

$\alpha_{\text{tum}}$	Takeoff turn angle: angle of the body midline axis with respect to the y-axis at takeoff.
$\hat{\mathbf{a}}$	3D axis of rotation (unit vector)
$\mathbf{A}$	Matrix of rotation axes (ns rows x 3 columns)
COM	Centre of mass
$d$	Gain multiplier for inverse kinematics (value between 0-1)
$dt$	Time sample interval
<b>error</b>	The 3D vector between the inverse kinematics target and current position
$e$	The error value: Norm of the <b>error</b> vector (= Euclidean distance between inverse kinematics target and the current position).
$e_{\text{min}}$	Minimum error value
$i$	Loop iterator for body segments
$[i]$	Index from a vector or matrix. E.g. $\mathbf{P}[2]$ would be the 2 <sup>nd</sup> row in the matrix $\mathbf{P}$ (i.e. XYZ for the 2 <sup>nd</sup> body point).
$[i,j]$	Row and column indices from a matrix. E.g. $\mathbf{P}[2,3]$ would be 3 <sup>rd</sup> column from the 2 <sup>nd</sup> row in the matrix $\mathbf{P}$ (i.e. the Z coordinate for the 2 <sup>nd</sup> body point).
$\mathbf{J}$	The Jacobian matrix (6 rows x ns columns)
$\mathbf{J}_T$	The translational component of the Jacobian (3 rows x ns columns)
$\mathbf{J}_R$	The rotational component of the Jacobian (3 rows x ns columns)
$nt$	Number of time samples for simulated kinematics
$\hat{\mathbf{n}}$	Normal unit vector
$ns$	Number of body segments
$\theta$	Angle of rotation
$\mathbf{p}$	X, Y, Z coordinates of a digitized point
$\mathbf{p}'$	X, Y, Z coordinates of an interpolated point
$\mathbf{P}$	Matrix of XYZ coordinates for segment endpoints along the body. Its dimensions are ns X 3
$\mathbf{P}_0$	Initial configuration: matrix of XYZ coordinates at the beginning of jump.
$\mathbf{P}_1$	Final configuration: matrix of XYZ coordinates at takeoff.
$\mathbf{q}$	A unit quaternion
$\mathbf{q}^*$	The conjugate of a quaternion
$\mathbf{q}_0$	The null rotation [1, 0, 0, 0] resulting in no rotation
$\mathbf{q}_i$	The quaternion for the i <sup>th</sup> segment (= $\mathbf{Q}[i]$ )
$\mathbf{qI}(\mathbf{q}_1, \mathbf{q}_2, \square)$	A unit quaternion interpolated between $\mathbf{q}_1$ and $\mathbf{q}_2$ at time $\square$ .
$\mathbf{Q}$	Quaternionized limb: vector of body segments expressed as quaternions
$\mathbf{Q}_0$	Quaternionized limb at the initial configuration (pre-jump)
$\mathbf{Q}_1$	Quaternionized limb at the final configuration (takeoff), initial guess
$\mathbf{Q}_{1R}$	Mirror image of $\mathbf{Q}_1$ representing the final configuration of the opposite (right) leg
$\mathbf{Q}_1'$	Quaternionized limb at the final configuration (takeoff) used for simulation
$\Theta$	ns x 1 vector of rotation angles (i.e. $\square \Theta = [\theta_1, \theta_2, \dots, \theta_{ns}]^T$ )
$\Delta \Theta$	ns x 1 vector of small changes in rotation angles for inverse kinematics
$\tau$	Relative time (from 0 to 1)



$\tau'$	Adjusted relative time ( $0 \leq \tau' < 1$ ) to prevent overshoot of COM position at takeoff
$T(\tau)$	A non-linear time function used to simulate acceleratory motion
$\mathbf{v}$	A 3D vector
$\mathbf{v}_q$	A 3D vector expressed as a (non-unit) quaternion
$\mathbf{v}_{\text{ref}}$	An arbitrary reference vector to represent the “zero” orientation, usually chosen to be the z-axis, [0, 0, 1]
$\psi$	The takeoff pitch angle: angle of the body midline axis with respect to horizontal at takeoff

962

963 APPENDIX A: *Miscellaneous calculations*

964 *Calculating angles between vectors*

965 Any two 3D vectors can form a plane. The angle between the vectors in this invisible plane is  
 966 calculated by the following:

$$angle = \cos^{-1} \frac{\mathbf{v1} \cdot \mathbf{v2}}{\|\mathbf{v1}\| \|\mathbf{v2}\|} \quad (\text{A1})$$

967

968 Where  $\cdot$  is the dot product and  $\mathbf{v1}$  and  $\mathbf{v2}$  are vectors of any dimension. For example,  $\mathbf{v1}$  and  
 969  $\mathbf{v2}$  can be 3D vectors for a 3D angle or the can be quaternions to compute the 4D angle  
 970 between them.

971 *Mirroring the leg kinematics between left and right sides*

972 To mirror the left leg we defined a plane of symmetry by calculating a normal vector to the  
 973 plane:

$$\hat{\mathbf{n}} = \frac{\mathbf{a} \times \mathbf{v}_{\text{ref}}}{\|\mathbf{a} \times \mathbf{v}_{\text{ref}}\|} \quad (\text{A2})$$

974

975 where  $\mathbf{a}$  is an axis within the plane (the body midline axis was used in the current study) and  
 976  $\mathbf{v}_{\text{ref}}$  is a reference vector in the plane ( $\mathbf{v}_{\text{ref}} = [0, 0, 1]$  for the present study). A reflection  
 977 matrix to reflect an XYZ point about an arbitrary plane is given by (Kovács, 2012):

$$R_{\text{reflect } 4 \times 4} \tag{A3}$$

$$= \begin{bmatrix} 1 - 2n_1^2 & -2n_1n_2 & -2n_1n_3 & o_1 - (1 - 2n_1^2)o_1 + 2n_1n_2o_2 + 2n_1n_3o_3 \\ -2n_1n_2 & 1 - 2n_2^2 & -2n_2n_3 & 2n_1n_2o_1 + o_2 - (1 - 2n_2^2)o_2 + 2n_2n_3o_3 \\ -2n_1n_3 & -2n_2n_3 & 1 - 2n_3^2 & 2n_1n_3o_1 + 2n_2n_3o_2 + o_3 - (1 - 2n_3^2)o_3 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

978  
 979 where  $n_1, n_2$  and  $n_3$  are the x, y and z components of  $\hat{\mathbf{n}}$  and  $o_1, o_2$  and  $o_3$  are the x, y and z  
 980 components of the local origin (the proximal end of the thigh was used for the present study).

981  
 982 *Quaternion arithmetic*

983 Quaternion arithmetic is required for performing rotations. A 3D vector can be expressed as  
 984 a quaternion:

$$\mathbf{v}_q = [\mathbf{0}, \mathbf{v}[1], \mathbf{v}[2], \mathbf{v}[3]] \tag{A4}$$

985  
 986 Where  $\mathbf{v}_q$  is used to denote a vector expressed as a quaternion (it is not necessarily a unit  
 987 quaternion, thus we avoid calling it “q”) and  $\mathbf{v}[1], \mathbf{v}[2], \mathbf{v}[3]$  are the XYZ components of the  
 988 3D vector  $\mathbf{v}$ .

989 The conjugate of a quaternion,  $\mathbf{q}^*$ :

$$\mathbf{q}^* = [\mathbf{q}[1], -\mathbf{q}[2], -\mathbf{q}[3], -\mathbf{q}[4]] \tag{A5}$$

990  
 991 Where  $\mathbf{q}[1], \mathbf{q}[2], \mathbf{q}[3], \mathbf{q}[4]$  are the 4 scalar values of the quaternion. A rotation operation is  
 992 as follows:

$$\mathbf{v}' = \mathbf{q} \otimes (\mathbf{v}_q \otimes \mathbf{q}^*) \tag{A6}$$

993

994 Where  $\otimes$  denotes quaternion multiplication. The first element of the rotated vector  $\mathbf{v}'$  should  
 995 be discarded to yield a 3D vector. In practice, Eq. A4-A6 can be combined to a single  
 996 function,  $\mathbf{v} \xrightarrow{\mathbf{q}} \mathbf{v}'$ .

997

## 998 APPENDIX B: *Converting between quaternion rotations and XYZ coordinates*

999

### 1000 *Quaternionization (P → Q)*

1001 The process of “quaternionization” converts a list of segment vectors (XYZ coordinates) to  
 1002 quaternions. It is computed in the following steps in a loop from  $i=1$  to  $i=ns$  (number of body  
 1003 segments):

1004 Step 1. Establish a reference vector,  $\mathbf{v}_{ref}$  ( $= [0, 0, 1]$  for the present study to represent the  
 1005 global reference segment).

1006 Step 2. Establish an empty matrix of limb segment coordinates,  $\mathbf{Q}$  which has dimensions ( $ns$   
 1007  $\times 3$  ( $= 5 \times 3$  for the present study)).

1008 Step 3. Begin the loop: For segment  $i \dots$

1009 Step 4. Calculate the quaternion between adjacent segment vectors  $\mathbf{v}_{ref}$  and  $\mathbf{V}[i]$  (i.e. the  $i^{th}$   
 1010 row of  $\mathbf{V}$ ). This is done by first calculating the axis:

$$\mathbf{A}_i = \frac{\mathbf{v}_{ref} \times \mathbf{V}[i]}{\|\mathbf{v}_{ref} \times \mathbf{V}[i]\|} \quad (\text{A7})$$

1011

1012 then calculating the angle between  $\mathbf{v}_{ref}$  and  $\mathbf{V}[i]$  using Eq. A1. Finally,  $\mathbf{q}_i$  is obtained by  
 1013 substitution into Eq. 1.

1014 Step 5. Overwrite  $\mathbf{v}_{ref}$ :  $\mathbf{v}_{ref} = \mathbf{V}[i]$  .

1015 Step 6. Insert  $\mathbf{q}_i$  into matrix  $\mathbf{Q}$  at the  $i^{th}$  row:  $\mathbf{Q}[i] = \mathbf{q}_i$  .

1016 Step 7. Increment  $i$ :  $i=i+1$ ; then return to step 4 until  $i=ns$ .

1017

1018 *Forward kinematics computation ( $\mathbf{Q} \rightarrow \mathbf{P}$ )*

1019

1020 Forward kinematics are performed in the following steps in a loop from  $i=1$  to  $i=ns$  (number  
1021 of body segments):

1022 Step 1. Establish a reference vector,  $\mathbf{v}_{\text{prox}}$  ( $= [0, 0, 1]$  for the present study to represent the  
1023 proximal segment) and a proximal joint to anchor each segment,  $\mathbf{v}_{\text{joint}} = [0, 0, 0]$ .

1024 Step 2. Establish an empty matrix of limb segment coordinates,  $\mathbf{P}$  which has dimensions ( $ns$   
1025  $+1$ )  $\times 3$  ( $= 6 \times 3$  for the present study). Set the first row of  $\mathbf{P}$  to be the limb anchor ( $\mathbf{P}[1] =$   
1026  $[0, 0, 0]$ ). Each row of  $\mathbf{P}$  will become a point on the body (i.e.  $\mathbf{P} = [\mathbf{P}[1], \mathbf{P}[2], \mathbf{P}[3], \mathbf{P}[4],$   
1027  $\mathbf{P}[5], \mathbf{P}[6)]^T = [\text{snout}, \text{hip}, \text{knee}, \text{ankle}, \text{TMT}, \text{foot}]^T$ ).

1028 Step 3. Begin the loop: For segment  $i$  ...

1029 Step 4. Calculate the distal vector,  $\mathbf{v}_{\text{dist}}$ , using quaternion rotation via the  $i^{\text{th}}$  quaternion (i.e.  
1030 the  $i^{\text{th}}$  row of  $\mathbf{Q}$ ):  $\mathbf{v}_{\text{prox}} \xrightarrow{\mathbf{q}_i} \mathbf{v}_{\text{dist}}$ ; then normalize the new vector:  $\mathbf{v}_{\text{dist}} = \frac{\mathbf{v}_{\text{dist}}}{\|\mathbf{v}_{\text{dist}}\|}$

1031 Step 5. Overwrite  $\mathbf{v}_{\text{prox}}$ :  $\mathbf{v}_{\text{prox}} = \mathbf{v}_{\text{dist}}$ .

1032 Step 6. Update the joint anchor position:  $\mathbf{v}_{\text{joint}} = \mathbf{v}_{\text{joint}} + l_i(\mathbf{v}_{\text{prox}})$  where  $l_i$  is the length of the  $i^{\text{th}}$   
1033 segment.

1034 Step 7. Insert  $\mathbf{v}_{\text{joint}}$  into matrix  $\mathbf{P}$  at position  $i + 1$ :  $\mathbf{P}[i+1] = \mathbf{v}_{\text{joint}}$ .

1035 Step 8. Increment  $i$ :  $i=i+1$ ; then return to step 4 until  $i=ns$ .

1036 Step 9. Anchor the frog at  $\text{XYZ} = [0, 0, 0]$  so that the frog leg extends upward. This is done  
1037 by subtracting the final point from each  $i^{\text{th}}$  XYZ point ( $\mathbf{P}[i] = \mathbf{P}[ns+1] - \mathbf{P}[i]$ ).

1038 Finally, if needed,  $\mathbf{P}$  can be converted to local vectors,  $\mathbf{V}$ , using Eq. 3.

1039

1040 APPENDIX C

1041 *Inverse kinematics: deriving the Jacobian matrix*

1042 For each simulated time value ( $\tau$ ), the left leg was mirrored to create the right leg (Appendix  
1043 A) which does not guarantee that the left and right hips join. Inverse kinematics (IK) was  
1044 used on the right leg to apply slight a correction to allow the hips to meet. This process was  
1045 repeated for each value of  $\tau$ .

1046 IK was briefly described previously (Richards et al., 2017), although a more complete  
1047 treatment will be necessary here. A common problem for models with multiple linked  
1048 segments (e.g. animal limbs or robotic manipulators) is that joint angles must be controlled to  
1049 guide the “end effector” of the limb (e.g. hand) to a specific target in 3D space. In the  
1050 present study, the limb is the right leg, the end effector is the right hip and the target is the  
1051 left hip. Problematically, there are often multiple solutions; i.e. there can be multiple different  
1052 limb configurations that allow the end effector to reach the target. A standard approach is to  
1053 1) calculate the error (Euclidean distance) from the target 2) move incrementally in the  
1054 direction of the target 3) return to step 1 and repeat until the error,  $e$ , falls below a given  
1055 tolerance,  $e_{min}$ . The error is simply  $\mathbf{p}_{current} - \mathbf{p}_{target}$  which itself is a velocity correction vector,  
1056  $\mathbf{v}_{corr}$ ; i.e. moving in the direction of the vector will bring the end effector closer to the target.  
1057 This is achieved using a Jacobian matrix,  $\mathbf{J}$ , which converts small changes in joint angles into  
1058 end effector velocity. Specifically,

$$\mathbf{J}_T \cdot \Delta\Theta = \mathbf{v}_{corr} \quad (C1)$$

1059 Where  $\mathbf{J}_T$  is the 3 X  $ns$  translational portion of the Jacobian matrix (see below),  $\Delta\Theta$  is the  $ns$   
1060 X 1 vector of joint angle changes for  $ns$  number of segments ( $ns = 5$  in the present study  
1061 representing torso, thigh, shank, proximal foot, distal foot). In other words,  $\Delta\Theta$  is a list of  
1062 unknown small changes (corrections) in each joint angle to produce incremental motion

1063 towards the target. Importantly, joint angles here are not Euler angles. Rather, they are  
 1064 angles about instantaneous rotation axes embedded in the quaternions (see below). The  
 1065 pseudoinverse of  $\mathbf{J}_T, \mathbf{J}_T'$ , allows us to solve for the unknown  $\Delta\Theta$ .

$$\mathbf{J}_T' \cdot \mathbf{v}_{\text{corr}} = \Delta\Theta \quad (\text{C2})$$

1066 We took an unconventional approach of defining  $\mathbf{J}$  using axis-angle coordinates to avoid  
 1067 pitfalls of Euler angles and to allow direct conversion to/from quaternions, (i.e. mapping  
 1068  $\mathbf{Q} \rightarrow \mathbf{J}$  via axis angle parameters). Recalling that a unit quaternion can be composed of a 3D  
 1069 rotation axis and an angle about that axis (Eq. 1), quaternions can likewise be decomposed to  
 1070 axis-angle parameters. For each segment, an  $ns \times 1$  vector of angles,  $\Theta$ , were calculated:

$$\Theta[i] = 2 \cos^{-1} \mathbf{Q}[i, 1] \quad (\text{C3})$$

1071 Where  $\mathbf{Q}[i, 1]$  (a scalar value) is the first quaternion element taken from  $i^{\text{th}}$  quaternion of  $\mathbf{Q}$   
 1072 (specifically,  $\mathbf{Q}_{\mathbf{R}}[i, 1]$ ). Given the angles, the axes can then be computed. Each row of  $\mathbf{A}$  is  
 1073 computed in a loop iteration in proximal to distal order:

$$\mathbf{A}[i] = \frac{\mathbf{q}_{iv}}{\sin \Theta[i]} / \left\| \frac{\mathbf{q}_{iv}}{\sin \Theta[i]} \right\| \quad (\text{C4})$$

1074 Where  $\mathbf{q}_{iv}$  is the vector component of the  $i^{\text{th}}$  quaternion (i.e. the 2<sup>nd</sup>, 3<sup>rd</sup> and 4<sup>th</sup> elements of  
 1075  $\mathbf{q}_i$ ).

1076 Finally, we can assemble  $\mathbf{J}$ , a matrix with 6 rows and  $ns$  columns. Each column is computed  
 1077 in single a loop iteration from  $i = 1$  to  $i = ns$  in proximal to distal order from snout to toe  
 1078 (snout, hip, knee, ankle, TMT).

$$\mathbf{J}_T[i] = \mathbf{A}[i] \times (\mathbf{p}[ns + 1] - \mathbf{p}[i]) \quad (\text{C5})$$

$$\mathbf{J}_R[i] = \mathbf{A}[i]$$

$$\mathbf{J}[i] = \begin{bmatrix} \mathbf{J}_T[i] \\ \mathbf{J}_R[i] \end{bmatrix}$$

1079 Where  $\mathbf{J}_T[i]$  is the  $i^{th}$  column of  $\mathbf{J}_T$  (similarly for  $\mathbf{J}_R[i]$ ),  $\mathbf{p}[i]$  is the XYZ point of the proximal  
1080 end point of the  $i^{th}$  segment and  $\mathbf{p}[ns+1]$  is the most distal endpoint of the most distal  
1081 segment (i.e. the foot point of ground contact).  $\mathbf{A}[i]$  is the  $i^{th}$  row of the rotation axes matrix  
1082 (see below) which is the instantaneous axis of planar rotation between segments  $i$  and  $i+1$ .  
1083 To assemble the full Jacobian,  $\mathbf{J}_T[i]$  and  $\mathbf{J}_R[i]$  are stacked to make  $\mathbf{J}_i$ , a column of 6 rows.  $\mathbf{J}_R$   
1084 is provided here for completeness, however it was not used in the present analysis.

1085 In practice, IK is done over several iterations moving a small fraction ( $d$ ) of the calculated  
1086  $\Delta\Theta$ .

$$\Delta\Theta = d(\mathbf{J}'_T \cdot \mathbf{v}_{corr}) \quad (C6)$$

1087 Where  $d$  is a small value ( $0 < d \leq 1$ ; we used  $d=0.1$ ). Larger values move the limb faster  
1088 towards the target, but excessively large values risk overshooting the target.

$$\Delta\Theta_{corr} = \Delta\Theta + \Theta \quad (C7)$$

1089 Then the corrected  $\mathbf{Q}$  is then computed by substituting into Eq. 1:

$$\mathbf{q}_i = \left[ \cos\left(\frac{\Delta\Theta_{corr}[i]}{2}\right), \mathbf{A}[i, 1] \sin\left(\frac{\Delta\Theta_{corr}[i]}{2}\right), \right. \quad (C8)$$

$$\left. \mathbf{A}[i, 2] \sin\left(\frac{\Delta\Theta_{corr}[i]}{2}\right), \mathbf{A}[i, 3] \sin\left(\frac{\Delta\Theta_{corr}[i]}{2}\right) \right]$$

1090 Where  $\mathbf{q}_i$  is the quaternion for the  $i^{th}$  segment within  $\mathbf{QIR}$  and  $\Delta\Theta_{corr}[i]$  is a scalar angle  
1091 correction value for the  $i^{th}$  segment

1092 *Inverse kinematics: numerical algorithm*

1093 Numerical integration was used to implement the IK correction using the following algorithm  
1094 for each time sample (i.e. each division of  $0 < \tau \leq \tau'$  up until  $nt$  samples).

1095 At time  $\tau = t \dots$

1096 Step 1. Mirror the left leg segments (Eq. 11) and quaternionize (Eq. 6), giving  $\mathbf{QI}_R$  at time  $t$ .

1097 Step 2. Calculate the error vector:  $\mathbf{error} = \mathbf{p}_{\text{hip,left}} - \mathbf{p}_{\text{hip,right}}$  and its magnitude ( $e =$   
1098  $\text{Norm}[\mathbf{error}]$ ).

1099 Step 3. Loop the following steps while  $e > e_{\text{min}}$  (the present study used  $e_{\text{min}} = 0.001$ ). If  $e \leq$   
1100  $e_{\text{min}}$  then skip to step 9.

1101 Step 4. Calculate the Jacobian (Eq. B3-B5) for  $\mathbf{QI}_R$  at time  $t$  and its pseudoinverse.

1102 Step 5. Calculate the correction angles then update  $\mathbf{QI}_R$  using Eq. B6-B8.

1103 Step 6. Perform forward kinematics to convert the quaternions  $\mathbf{QI}_R$  to limb XYZ coordinates  
1104 ( $\mathbf{Q} \rightarrow \mathbf{P}$ ; Appendix B steps 1-9).

1105 Step 7. Re-calculate the error:  $\mathbf{error} = \mathbf{p}_{\text{hip,left}} - \mathbf{p}_{\text{hip,right}}$  and its magnitude ( $e = \text{Norm}[\mathbf{error}]$ ).

1106 Step 8. Return to step 3.

1107 Step 9. Advance to the next time step,  $t = t + dt$ , then return to step 1 until  $t = \tau'$ . (i.e. for all  
1108 allotted timesteps.

1109