

# DEFORMATION MAPS FOR BOLTED T-STUBS

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## ABSTRACT

Deformation maps on the plastic and ultimate failure of T-stubs with a single bolt-row in tension are developed. The maps condense a large body of information within a two dimensional parameters space onto which different modes of deformation, including the régime boundaries, are plotted for any practical combination of geometric and material properties encountered in a T-stub. Its fidelity is demonstrated with experimental data from literature, and through a detailed parametric investigation by high-fidelity finite element analysis. The predictive capability of two existing analytical models are also assessed against predictions by the finite element model. Their shortcomings, and applicability, are critically assessed and discussed.

**Keywords:** T-stubs, Eurocode 3, Deformation maps, Failure modes, Joints.

## 1 INTRODUCTION

2 A bolted end plate connection is a class of moment resisting connections that  
3 is widely encountered in steel-framed structures. Its overall resistance is offered  
4 through a combination of tensile forces that act in the bolts adjacent to one flange,

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5 and compressive forces experienced by the bearing at the other. Unless significant  
6 catenary actions develop in its adjoining beam, these tensile and compressive forces  
7 are typically assumed to be equal and opposite. In general, the rotational capacity  
8 of any bolted joint is limited by the deformation within its tension zone which  
9 comprises of the column flange, end-plate and bolts in tension. In Eurocode 3  
10 (EN 1993-1-8 2005) - denoted hereinafter as EC3 - the tension zone is modelled  
11 by a T-stub and the ability to predict, *a priori*, its mode of deformation for a  
12 broad range of material (T-stub and bolt) and geometric combinations is integral  
13 to characterising the deformation capacity of a structural joint.

14 For T-stubs comprising of a single bolt-row in tension, three modes of failure,  
15 or régimes of deformation, can develop: mode 1 - complete yielding of the flange;  
16 mode 2 - bolt failure with yielding of the flange; and, mode 3 - bolt failure. Within  
17 the constitutive framework of limit analysis, Piluso et al. (2001) showed that each  
18 mode corresponds to a unique range of non-dimensional parameter  $\beta$  (Eq. 1). To  
19 extend the aforesaid to include material strain-hardening and ‘ultimate’ failure  
20 prediction, Piluso et al. (2001) adopted a piece-wise linear approximation of the  
21 true stress-strain curve of the flange material. A disadvantage is that  $\beta$  must be re-  
22 evaluated – including the critical values corresponding to the transition between  
23 régimes  $1 \rightarrow 2$  ( $\beta_{cr}^{1 \rightarrow 2}$ ) and  $2 \rightarrow 3$  ( $\beta_{cr}^{2 \rightarrow 3}$ ) – each time a different combination  
24 of material and geometry (dimensions) is encountered. At present, there is no  
25 straightforward means to visualise, and represent, the different modes and their  
26 régime boundaries for a practical range of material and geometric parameters in a  
27 compact design space that would be useful to designers.

28 Failure (both plastic and ultimate) prediction using  $\beta$  is a resistance-based ap-  
29 proach that does not shed light on the deformation of the T-stub. To this end,

30 analytical models were developed by Piluso et al. (2001) and Francavilla et al.  
 31 (2016) – it must be emphasized that both neglect three-dimensional (3D) effects,  
 32 geometric non-linearity, moment-shear interaction, to name a few – where their  
 33 predictions were shown to agree well with the force-displacement curves of T-stubs  
 34 failing in mode 1. However, the model by Piluso et al. (2001) overestimates the  
 35 ultimate displacement of T-stubs that fail in mode 2, largely as a consequence of  
 36 neglecting displacement compatibility between the bolt and flange. This simplifi-  
 37 cation was relaxed by Francavilla et al. (2016) which reduces the overestimation  
 38 of  $\Delta$  in mode 2. No experimental data were available in mode 3; consequently,  
 39 comparison to analytical predictions was not performed.

40 In this paper, we exploit the dimensionless parameter  $\beta$  (Piluso et al. 2001) to  
 41 construct deformation maps on the plastic and ultimate failure of T-stubs with a  
 42 single bolt-row in tension. The fidelity of the deformation maps will be demon-  
 43 strated through experimental data from Girão Coelho (2004), Bursi and Jaspart  
 44 (1997) and Piluso et al. (2001). Further validation is provided through a detailed  
 45 parametric investigation by three-dimensional finite element (FE) modelling, which  
 46 considers material damage in both the flange and bolt, as well as geometric non-  
 47 linearity. Predictions by the FE model are used to assess the accuracy of the two  
 48 aforementioned analytical models.

## 49 **DEFORMATION MAPS**

The mode of deformation that develops in a T-stub is related to a non-dimensional  
 parameter given by (Piluso et al. 2001)

$$\beta = \frac{2 M}{m B} \quad (1)$$

where  $m$  is the distance from the axis of the bolt-hole to the plastic hinge at the

flange-to-web connection;  $M$  is the bending moment at the plastic hinges; and,  $B$  is the tensile force in each bolt shown in Fig 1. Since  $M = M(l_{eff}, t_p, f_y)$  and  $B = B(A_s, f_{ub})$ , the parametric dependence for  $\beta$  can be expressed – by making use of dimensional analysis – as follows:

$$\beta = f\left(\frac{l_{eff}}{m}, \frac{f_y}{f_{ub}}, \frac{t_p^2}{A_s}\right), \quad (2)$$

which can be re-arranged to give

$$\frac{t_p^2}{A_s} = g\left(\beta, \frac{l_{eff}}{m}, \frac{f_y}{f_{ub}}\right) \quad (3)$$

50 where  $A_s$  is the tensile stress area of the bolt;  $f_y$  and  $f_{ub}$  are the yield strength  
51 of the flange material and the ultimate strength of the bolt, respectively; and,  
52  $t_p$  is the flange thickness. The effective length  $l_{eff}$  is the notional width defined  
53 such that, at plastic collapse, the resistance of the flange – modelled as beams  
54 – is equivalent to a T-stub whose kinematic mechanism is determined by the  
55 yield line pattern that develops. Figure 2 shows the different possible yield line  
56 patterns according to EN 1993-1-8 (2005). A circular pattern arises due to localised  
57 action of the bolts; hence, they only develop in mode 1 where prying forces  $Q_1$   
58 are small compared to that ( $Q_2$ ) in mode 2. By contrast, non-circular patterns  
59 can develop in both modes 1 and 2, and are characterised by significant prying  
60 forces that can cause premature bolt failure. In EN 1993-1-8 (2005), prying forces  
61 are assumed to act along the edges of a T-stub at a distance  $n$  from the bolt  
62 axis. According to McGuire and Winter (1978), a non-circular pattern will develop  
63 whenever the ratio  $\lambda = n/m < 1.25$ . It is worth emphasising that both circular and  
64 non-circular patterns are three-dimensional (3D) yielding mechanisms since their  
65 hinge-line profile changes along the width of the T-stub, see Fig 2. By contrast, a  
66 beam yield-line pattern is not and it develops whenever the width  $L$  is small, i.e.

67  $L < 4m + 1.25n$ . Here, only non-circular and beam patterns that induce significant  
68 prying forces are considered. Note that  $l_{eff}/m$  in Eq. 3 depends on the yield-line  
69 pattern that develops in the flange. For a non-circular pattern,  $l_{eff}/m = 4 + 1.25\lambda$   
70 (EN 1993-1-8 2005). In a beam pattern, it is assumed, according to standard  
71 specifications for HE beams (UNI 5397:1978 1978), that  $L = 2.5 \cdot m$ .

To construct the maps, consider the plastic and ultimate failure of a T-stub separately. For a fully-plastic limit state, the corresponding  $\beta$  is given by (Piluso et al. 2001)

$$\beta = \beta_{Rd} = \frac{2 M_{pl,Rd}}{m B_{Rd}}, \quad (4)$$

where  $B_{Rd}$  is the design tensile resistance of the bolt given by (EN 1993-1-8 2005)

$$B_{Rd} = 0.9 A_s f_{ub} \quad (5)$$

and  $M_{pl,Rd}$  is the design flexural resistance of the flange given by (EN 1993-1-8 2005)

$$M_{pl,Rd} = 0.25 l_{eff} t_p^2 f_y. \quad (6)$$

Substituting for  $B_{Rd}$  and  $M_{pl,Rd}$  in Eq. 4, and re-arranging according to the dimensionless groups in Eq. 3, one obtains:

$$\frac{t_p^2}{A_s} = \frac{1.8}{4 + 1.25\lambda} \beta_{Rd} \left(\frac{f_y}{f_{ub}}\right)^{-1} \quad \text{Non-circular pattern} \quad (7a)$$

$$\frac{t_p^2}{A_s} = 0.72 \beta_{Rd} \left(\frac{f_y}{f_{ub}}\right)^{-1} \quad \text{Beam pattern} \quad (7b)$$

Note that the relevant expression for  $l_{eff}/m$  corresponding to its respective yield-line pattern, given earlier, was substituted into Eq. 7. For an ultimate limit state, (Piluso et al. 2001)

$$\beta = \beta_u = \frac{2 M_f}{m B_u}. \quad (8)$$

Here,  $B_u$  is the ultimate tensile resistance of the bolt given by (Piluso et al. 2001)

$$B_u = A_s f_{ub} \quad (9)$$

72 and  $M_f$  is the bending moment at material fracture of the flange. Adopting an  
 73 idealised true stress-strain relation of the flange material in Fig. 3 – note that  
 74  $\sigma_y (= f_y)$  and  $\sigma_u (= f_u)$  is the yield and ultimate tensile strength, respectively –  
 75 the bending moment at fracture can be expressed as (Piluso et al. 2001)

$$\frac{M_f}{M_y} = \kappa = \frac{1}{2} \left[ 3 - \left( \frac{\varepsilon_y}{\varepsilon_f} \right)^2 \right] + \frac{1}{2} \frac{E_h}{E} \left( \frac{\varepsilon_f - \varepsilon_h}{\varepsilon_y} \right) \left( 1 - \frac{\varepsilon_h}{\varepsilon_f} \right) \left( 2 + \frac{\varepsilon_h}{\varepsilon_f} \right) - \frac{1}{2} \frac{E_h - E_f}{E} \frac{\varepsilon_f - \varepsilon_u}{\varepsilon_y} \left( 1 - \frac{\varepsilon_u}{\varepsilon_f} \right) \left( 2 + \frac{\varepsilon_u}{\varepsilon_f} \right), \quad (10)$$

where the corresponding moment at yield is

$$M_y = \frac{l_{eff} t_p^2}{6} f_y. \quad (11)$$

The constant  $\kappa$  in Eq. 10 is a function of the flange material so that  $\kappa = 1$  at fully plastic condition. Following the same procedure as before, one obtains

$$\frac{t_p^2}{A_s} = \frac{3}{4 + 1.25\lambda} \beta_u \frac{1}{\kappa} \left( \frac{f_y}{f_{ub}} \right)^{-1} \quad \text{Non-circular pattern} \quad (12a)$$

$$\frac{t_p^2}{A_s} = 1.2 \beta_u \frac{1}{\kappa} \left( \frac{f_y}{f_{ub}} \right)^{-1} \quad \text{Beam pattern} \quad (12b)$$

76 where the ratio  $\lambda \triangleq n/m$ . Régime boundaries are obtained by substituting  $\beta_{Rd}$   
 77 and  $\beta_u$  in Eqs. 7 and 12 with the corresponding expression for  $\beta$  in Table 1. The  
 78 final expression for the régime boundaries has the following general form:

$$\frac{t_p^2}{A_s} = h(\lambda) \frac{1}{\kappa} \left( \frac{f_y}{f_{ub}} \right)^{-1} \quad (13)$$

79 where  $h(\lambda)$  is tabulated in Table 2. Deformation maps for plastic and ultimate  
 80 failure of a T-stub – generated by plotting the régime boundaries (Eq. 13) in a

81 plot of  $t_p^2/A_s$  versus  $\kappa f_y/f_{ub}$  – are shown in Figs. 4 and 5, respectively. Notice  
82 that each boundary is an isoline corresponding to a constant  $\lambda$  value. Any pair  
83 of geometric ( $t_p^2/A_s$ ) and material ( $\kappa f_y/f_{ub}$ ) parameters now uniquely locates a  
84 point on the 2D map. From the map, one is able to determine the deformation  
85 mode for a given value of  $\lambda$  (or by interpolation between any two  $\lambda$  values plotted  
86 in Figs. 4 and 5, if required). Alternatively, the map allows a designer to select  
87 the combination of geometric and material parameters for a T-stub to deform in  
88 a desired mode. Note that only a single boundary demarcates the transition from  
89 mode 2→3 for the beam pattern since it is independent of  $\lambda$ . The fidelity of the  
90 maps will be validated later against experimental data from existing literature and  
91 numerical predictions by FE model to be developed next. It is worth emphasising  
92 that the deformation maps are only as accurate as the calculations/assumptions  
93 undertaken in their generation.

## 94 **FINITE ELEMENT MODELLING**

95 Three-dimensional finite element (FE) models are developed for T-stubs sub-  
96 jected to quasi-static tension with ABAQUS/Standard V6.13. Predictions will be  
97 validated against experimental data from three separate independent sources: viz.  
98 WT1 (Girão Coelho 2004), T1 (Bursi and Jaspart 1997) and T15 (Ribeiro et al.  
99 2015). The acronyms correspond to that used in their respective original source.  
100 The aforesaid were performed for identical flange (S355 steel) and bolt (Grade 8.8)  
101 materials, but each have a different  $\lambda$  (WT1: 0.9, T1: 1.0 and T15: 0.7). All three  
102 failed in either mode 1 or 2 depending on the non-dimensional parameter  $t_p^2/A_s$   
103 which is different for each data-set. It is worth highlighting that both WT1 and  
104 T1 contained two bolt-rows. Since they both develop a beam yield line pattern,

105 3D effects can be neglected and interactions between the two bolt-rows need not  
 106 be considered (EN 1993-1-8 2005); consequently, they can be treated as equivalent  
 107 T-stubs with a single bolt row.

### 108 **Constitutive model**

109 All the flange and bolts in the experiments were constructed from S355 steel  
 110 and Grade 8.8 bolt (of Young's modulus  $E = 210$  GPa and Poisson's ratio  $\nu =$   
 111  $0.33$ ), respectively. Their subsequent plastic responses are modelled using the  
 112 conventional  $J_2$  plasticity flow theory to allow implementation of the progressive  
 113 degradation of material stiffness. Figure 6 shows the nominal, and corresponding  
 114 true stress-strain curves of the S355 flange material. Characteristic points on the  
 115 nominal and true stress-strain curves are indicated as follows: yield ( $y$ ), necking  
 116 ( $n$ ), rupture ( $r$  - the last point on the softening branch just before the stress drops  
 117 to 0) and fracture ( $f$ ). An index  $i$  denotes a generic data point on the stress-strain  
 118 curve (nominal or true) connecting points  $y$ ,  $n$ ,  $r$  and  $f$ . The  $i$ -th data point  
 119 of the plastic response is extracted from the nominal stress-strain curve through  
 120 (ABAQUS 2009)

$$\sigma_i = \begin{cases} \sigma_i^{\text{nom}}(1 + \varepsilon_i^{\text{nom}}), & y \leq i < n \\ \sigma_n^{\text{nom}}(1 + \varepsilon_i^{\text{nom}}), & n \leq i \leq f \end{cases} \quad (14)$$

121 and

$$\varepsilon_i^{\text{pl}} = \ln(1 + \varepsilon_i^{\text{nom}}) - \varepsilon_y, \quad (15)$$

122 where  $\sigma_n^{\text{nom}}$  is the nominal stress at necking and  $\varepsilon_y$  is the strain at yield. To account  
 123 for the effects of post necking strain localisation, we follow Pavlović et al. (2013)

124 by defining the nominal strain  $\varepsilon_i^{\text{nom}}$  as

$$\varepsilon_i^{\text{nom}} = \begin{cases} \Delta l_i / l_i, & y \leq i < n \\ \varepsilon_{i-1}^{\text{nom}} + (\Delta l_i - \Delta l_{i-1}) / l_i, & n \leq i \leq f, \end{cases} \quad (16)$$

125 where  $l_i$  – given in Eq. 17 – represents the gauge length at the  $i$ -th data point  
 126 and  $\Delta l_i = \varepsilon_i^{\text{nom}} l_0$ . For a cylindrical tensile coupon of diameter  $d$ , the initial  
 127 gauge length  $l_0$  ( $= 50$  mm) starts decreasing as the material softens. At the  
 128 point of rupture, the gauge length becomes  $l_{\text{loc}} = 0.5d$  (Panontin and Sheppard  
 129 1999): here,  $l_{\text{loc}} = 6$  mm. Following Pavlović et al. (2013), the reduction of the  
 130 gauge length is assumed to obey a power law through the localisation rate factor  
 131 ( $\alpha_L = 0.5$ ) given by

$$l_i = \begin{cases} l_0, & y \leq i < n \\ l_0 + (l_{\text{loc}} - l_0)[(\Delta l_i - \Delta l_n) / (\Delta l_r - \Delta l_n)]^{\alpha_L} & n \leq i \leq f, \end{cases} \quad (17)$$

132 where  $\Delta l_n = \varepsilon_n^{\text{nom}} l_0$  and  $\Delta l_r = \varepsilon_r^{\text{nom}} l_0$  are the elongations of the gauge length at  
 133 necking and rupture, respectively. The effects of strain localisation were not taken  
 134 into account for the bolt since experiments by Girão Coelho (2004) have shown  
 135 that the area reduction following the onset of necking affects the entire length of  
 136 its shaft. Instead, the true stress-plastic strain curve obtained by Girão Coelho  
 137 (2004) – this is plotted in Fig. 6 – is used for the bolt.

### 138 **Damage modelling**

139 The damage initiation criterion by ABAQUS (2009) is used to predict the onset  
 140 of damage over a wide range of stress states, given by

$$\omega_d = \int \frac{d\bar{\varepsilon}^{\text{pl}}}{\bar{\varepsilon}_0^{\text{pl}}(\theta)} = 1 \quad (18)$$

141 where  $\omega_d$  is a state variable that increases monotonically with  $\bar{\varepsilon}^{\text{pl}}$  (expressed as a  
 142 function of stress triaxiality  $\theta$ ). Here, the equivalent plastic strain at the onset of  
 143 damage  $\bar{\varepsilon}_0^{\text{pl}}$  is defined as (Pavlović et al. 2013)

$$\bar{\varepsilon}_0^{\text{pl}}(\theta) = \varepsilon_n^{\text{pl}} \cdot \exp[-1.5(\theta - 1/3)] \quad (19)$$

144 where  $\varepsilon_n^{\text{pl}}$  is the true plastic strain at necking. In the post necking regime ( $n \leq$   
 145  $i \leq f$ ), the damage process is controlled by the evolution of the damage variable  
 146  $D_i$  which is expressed as a function of the equivalent plastic displacement  $\bar{u}_i^{\text{pl}}$ .  
 147 Following Pavlović et al. (2013),  $D_i$  is given by

$$D_i = \begin{cases} (1 - \bar{\sigma}_i/\sigma_i) \cdot 1.5, & n \leq i \leq r \\ 1, & i = f \end{cases} \quad (20)$$

148 where  $\sigma_i$  is the true stress (Eq. 14) and  $\bar{\sigma}_i = \sigma_i^{\text{nom}}(1 + \varepsilon_i^{\text{nom}})$  for  $y \leq i \leq f$ . Notice  
 149 that no data points were considered between points  $r$  (rupture) and  $f$  (fracture)  
 150 in Fig. 6 since rupture is defined as the last point on the softening branch of the  
 151 nominal stress-strain curve, i.e.  $D_i = 1$  beyond point  $r$ .  $\bar{u}_i^{\text{pl}}$  is given by (Pavlović  
 152 et al. 2013)

$$\bar{u}_i^{\text{pl}} = \bar{u}_f^{\text{pl}}(\varepsilon_i^{\text{pl}} - \varepsilon_n^{\text{pl}})/(\varepsilon_f^{\text{pl}} - \varepsilon_n^{\text{pl}}), \quad n \leq i \leq f \quad (21)$$

153 where the equivalent plastic displacement at fracture  $\bar{u}_f^{\text{pl}}$  (in Eq. 21) is

$$\bar{u}_f^{\text{pl}} = \lambda_S \lambda_E L_E (\varepsilon_f^{\text{pl}} - \varepsilon_n^{\text{pl}}) \quad (22)$$

154 and  $L_E = \sqrt[3]{V}$  ( $V$  is the volume of the element) is the characteristic length of the  
 155 element (ABAQUS 2009, Sui et al. 2017). The factor  $\lambda_E$  must be calibrated to  
 156 the element type used in the FE model. The dependence of  $\bar{u}_f^{\text{pl}}$  on the mesh size is  
 157 removed by introducing an element size factor  $\lambda_S$ , which is obtained by reproducing  
 158 the tensile test in FE with different element size of  $L_{E,a}$ , where  $a \in [0, t]$  and  $t$  is

159 the number of element sizes considered. The reference element size is  $L_{E,a} = L_{E,0}$   
 160 for  $a = 0$ . The final expression for  $\lambda_S$  is (Pavlović et al. 2013)

$$\lambda_S = \sqrt[3]{\frac{L_{E,0}}{L_{E,a}}}, \quad a \in [1, t]. \quad (23)$$

161 It is worth emphasising that  $\lambda_E$  is used to model damage evolution in the S355  
 162 steel. For the Grade 8.8 bolt, a linear damage evolution law is used for simplicity  
 163 since it does not require the introduction of  $\lambda_E$ .

#### 164 **Calibration of $\lambda_E$**

165 Tension loading of dog-bone specimens for the S355 steel were simulated to  
 166 calibrate  $\lambda_E$ . Four specimens with cylindrical cross-sectional area, and dimen-  
 167 sions stipulated by the ASTM Standard E8/E8M-15a (2015), are discretised with  
 168 C3D8R (8 nodes linear continuum elements with reduced integration) elements of  
 169 a uniform size throughout its gauge section. The element sizes considered were:  
 170  $L_{E,0} = 0.89\text{mm}$  (186 elements across diameter),  $L_{E,1} = 1.00\text{mm}$  (140 elements),  
 171  $L_{E,2} = 1.25\text{mm}$  (96 elements) and  $L_{E,3} = 1.59\text{mm}$  (60 elements). Each element  
 172 size was chosen so that their aspect ratio is  $\approx 1$ . A displacement boundary condi-  
 173 tion of  $0.05\text{mm/s}$  (corresponding to a nominal strain rate of  $0.001\text{s}^{-1}$ ) is applied  
 174 to one end of the specimen, with the other end fully clamped, to simulate tensile  
 175 loading. The true stress-strain curve in Fig. 6 is used and damage initiation follows  
 176 Eq. 19.

177 A flow chart summarising the procedure to calibrate  $\lambda_E$  is shown in Fig. 7.  
 178 An initial FE model discretised uniformly with C3D8R elements of size  $L_{E,0}$  is  
 179 first considered with a trial  $\lambda_E$  value.  $\lambda_S$  is, then, evaluated - note that  $\lambda_S = 1$   
 180 (Eq. 23) for  $L_{E,a} = L_{E,0}$  - and Eqs. 20, 21 and 22 are used to obtain  $D_i$ ,  $\bar{u}_i^{\text{pl}}$   
 181 and  $\bar{u}_f^{\text{pl}}$ , respectively. The predicted stress-strain curve is then compared to its

182 experimental counterpart. If their differences  $Er(\varepsilon_r^{\text{nom}}) \leq 2\%$ , then the trial  $\lambda_E$   
183 value is satisfactory and the same procedure repeated for  $L_{E,1}$ ,  $L_{E,2}$  and  $L_{E,3}$ .  
184 If, however,  $Er(\varepsilon_r^{\text{nom}}) > 2\%$  for any  $L_{E,a}$ , the initial value of  $\lambda_E$  is updated and  
185 the simulations repeated for all  $L_{E,a}$ . Pavlović et al. (2013) have shown that  
186  $\lambda_E \in [2.5, 3.2]$  for C3D8R elements and the final value to use depends on the  
187 ductility of the material. Here, the calibrated value of  $\lambda_E = 2.5$  lies within the  
188 range stipulated by Pavlović et al. (2013).

189 The difference in the predicted nominal strain at rupture  $\varepsilon_r^{\text{nom}}$  with experiment  
190 is less than 2% for all  $L_{E,a}$ . A linear damage evolution law is assumed for the bolt,  
191 with an equivalent plastic displacement at fracture given by  $\bar{u}_f^{\text{pl}} = \varepsilon_f L_E$  (ABAQUS  
192 2009).  $L_E = 2\text{mm}$  is used here so that there is a minimum of 12 to 16 nodes across  
193 the bolt diameter (Viridi 1999); and,  $\varepsilon_f = 0.13$  follows from Girão Coelho (2004).  
194 Figure 8 shows the damage initiation and evolution curves that were implemented  
195 into the FE models.

## 196 Simulations of T-stub in tension and validation

197 Figure 9 shows a schematic of the tensile test simulated by FE. Recall that  
198 WT1 and T1 comprises of two bolt-rows; hence, only a quarter of the T-stub was  
199 simulated and  $u_z = 0$  must be specified for the  $x - y$  plane (Fig. 9a), unlike in  
200 T15. The bottom web is fully clamped and a displacement boundary condition  
201 of  $\dot{u}_y = 0.01\text{mm/s}$  (Girão Coelho 2004) was applied to the top web. Table 3  
202 lists the value of each geometric parameter shown in Figure 9b. Since T15 was  
203 obtained from an IPE300 beam profile (Bursi and Jaspart 1997), there is no weld  
204 to be modelled. By contrast, both WT1 (Girão Coelho 2004) and T1 (Ribeiro  
205 et al. 2015) were constructed by welding two plates together. In our FE model,

206 both welds and flange are assumed to be made of the same material (Girão Coelho  
 207 2004; Girão Coelho et al. 2004; Ribeiro et al. 2013; Latour and Rizzano 2012). The  
 208 same is also assumed for T15 given that no information was provided by Ribeiro  
 209 et al. (2015) regarding the weld material. The entire FE model is discretised using  
 210 C3D8R elements (Latour and Rizzano 2012; Latour et al. 2014). The bolts are  
 211 modelled as a solid cylinder with an equivalent cross-sectional area of  $A_s$  (Bursi  
 212 and Jaspert 1997). Figure 10a shows the results of three element sizes (4, 2 and  
 213 1.5mm) – the flange is meshed as in B4C5D4 of Fig. 10d (this will be justified  
 214 later) – that were used to mesh the bolt. Since a 1.5mm element predicted a  
 215 similar ultimate displacement  $\Delta_u$  – black circular dot in Fig. 10 – to its 2mm  
 216 counterpart, the latter was used to reduce computational cost. The mesh for the  
 217 bolt, used in subsequent parametric study, is shown in Fig. 9d. To determine the  
 218 mesh size for the flange, we partition the flange into four zones (A, B, C and D)  
 219 as shown in Fig. 10d. Zone A is discretised using C3D8R elements of different  
 220 sizes, while the rest of the flange was meshed as shown in B4C5D4 (Fig. 10d),  
 221 and their corresponding results are shown in Fig. 10b. After this, three mesh  
 222 densities were tested for zones B, C and D – see Fig. 10d – where the number of  
 223 elements along the x-axis is given after each letter. Figure 10a shows a negligible  
 224 difference between the predicted  $\Delta_u$  by B4C5D4 and B6C10D4. Hence, Region B  
 225 is discretised with four elements across the circle; while region C with four elements  
 226 between regions A & B and between regions B & D. Region D is discretised with 4  
 227 elements. The mesh for the flange, used in subsequent parametric study, is shown  
 228 in Fig. 9c. Surface-to-surface contact formulation with small sliding is prescribed  
 229 for all contact pairs – top and bottom flange, flanges and head/nut of the bolt,  
 230 bolt shank and hole – with a coefficient of Coulomb friction  $\mu = 0.25$  (Bursi and

231 Jaspert 1997).

232 Figure 11 compares the predicted deformation history of the T-stubs to its  
233 experimental counterpart. The resistance ( $F_u$ ) of a T-stub is defined as the max-  
234 imum resultant reaction force acting at the bottom web and its corresponding  
235 displacement is the ultimate displacement ( $\Delta_u$ ) - they are shown in Fig. 11 by red  
236 (experiment) and black (FEM) dots. The predicted force-displacement ( $F - \Delta$ )  
237 curve closely matches that of the experiments where the percentage differences for  
238  $F_u$  (WT1: 3.4%, T1: 3.1%, T15: 5.3%) and  $\Delta_u$  (WT1: 0.7%, T1: 5.1%, T15:  
239 2.6%) are small. In addition, comparison can also be made for the bolt elongation  
240  $\Delta_b$  in WT1 (note that this was not measured in T1 and T15) as shown in Fig. 11d.  
241 The discrepancy in the final bolt elongation between experiment and FE arises be-  
242 cause the former was halted at  $\Delta_b = 0.4\text{mm}$  to prevent equipment damage (Girão  
243 Coelho 2004). Notwithstanding, a good general agreement - differences of less  
244 than 5.1% - is observed until  $\Delta_b = 0.4\text{mm}$ . The mode of failure is also successfully  
245 predicted by the FE models; this is evident from the distributions of equivalent  
246 plastic strain ( $\varepsilon^{\text{pl}}$ ) and damage variable ( $D_i$ ) shown in Fig. 11. Note that the con-  
247 tour plots for the flange and bolt correspond to the ultimate displacement. Both  
248 WT1 and T15 failed in mode 2. A plastic hinge also develops at the weld-toe in the  
249 FE model which is highlighted in Figs. 11a and c: here,  $\varepsilon^{\text{pl}}$  exceeds the threshold of  
250 0.05 through the thickness, as suggested by Ribeiro et al. (2015). In Figs. 11a and  
251 c, the damage variable  $D_i = 1$  is reached in the bolt; consequently, elements are  
252 removed from the mesh. Specimen T1 developed two plastic hinges; at both the  
253 weld-toe and bolt hole. However, Fig. 11b also shows that  $D_i$  is close to unity in  
254 the bolt. This suggests that the bolt had undergone significant plastic deformation  
255 which agrees with observations by Bursi and Jaspert (1997) who suggested that

256 the actual failure mechanism is between mode 1 and 2. The location of fracture in  
257 the experiment was not indicated by Bursi and Jaspart (1997); however, the value  
258 of  $D_i$  in our FE model suggests that bolt fracture is imminent.

## 259 **RESULTS**

260 The FE model is now employed in a parametric study to investigate how the  
261 deformation capacity of a T-stub is affected by its geometric and material pa-  
262 rameters; and, the results will also be used to critically assess the accuracy, and  
263 limitations, of two existing analytical models developed by Piluso et al. (2001) and  
264 Francavilla et al. (2016).

### 265 **Parametric study**

#### 266 *Geometry and material properties*

267 Figure 9b shows a schematic of the T-stub and the geometric dimensions  
268 considered. Note that the tensile response of a T-stub that develops beam yield  
269 line pattern had been extensively studied, both experimentally and numerically, by  
270 others; see, for example, Piluso et al. (2001), Girão Coelho (2004) and Ribeiro et al.  
271 (2015). In addition, analytical models also exist that could accurately predict their  
272  $F - \Delta$  relationship (Piluso et al. 2001; Francavilla et al. 2016). By contrast, there  
273 are relatively fewer studies on non-circular yield line patterns. For this reason,  
274 all the T-stubs here were sized to develop this; hence, their width  $L$  must satisfy  
275  $L > 4m + 1.25n$ . Since  $B$  is fixed at 200mm in all the models,  $m$  and  $n$  are varied  
276 to obtain three different  $\lambda$  ( $\triangleq n/m$ ) values as follows: (1) A maximum value of  
277  $\lambda_{\max} = 1.25$  determined by considering the minimum standard bolt spacing of  
278  $p_{\min} = 98\text{mm}$  for a HEA200 beam (UNI 5397:1978 1978). This is taken as a  
279 reference beam profile since it is characterised by a width which is identical to  $B$

280 (200mm). Note that if  $\lambda > \lambda_{\max}(= 1.25)$ , non-circular patterns will not develop  
281 according to McGuire and Winter (1978); (2) A minimum value of  $\lambda_{\min} = 0.9$   
282 determined by considering the maximum allowable bolt spacing of  $p_{\max} = 14t_p$  as  
283 stipulated in EN 1993-1-8 (2005); (3) An intermediate value of  $\lambda_{\text{inter}} = 1.1$  - this  
284 was selected to lie between  $\lambda_{\max}$  and  $\lambda_{\min}$ . Three flange thicknesses  $t_p$  (8, 9 and  
285 10mm) and four bolt diameters  $d_b$  (10, 12, 20 and 27mm) are considered. Table 4  
286 tabulates the combinations of dimensions that were modelled.

287 Four different grades of structural steel (S235, S275, S355 and S450) are mod-  
288 elled for the flange material. Their respective true stress-strain curve is represented  
289 using a piece-wise approximation similar to Piluso et al. (2001) in Fig. 3. Key val-  
290 ues of stresses ( $\sigma_y$  and  $\sigma_u$ ), strains ( $\varepsilon_y$ ,  $\varepsilon_h$ ,  $\varepsilon_u$  and  $\varepsilon_f$ ) and moduli ( $E$ ,  $E_h$  and  $E_f$ )  
291 are tabulated in Table 5. For the S450 steel,  $\sigma_y$  and  $\sigma_u$  were obtained from EN  
292 1993-1-8 (2005),  $E_f = \sigma_u$  (Piluso et al. 2001) and  $\varepsilon_f = 0.17$  (EN 10025-2 2004);  
293 and,  $E$ ,  $E_h$ ,  $\varepsilon_y$ ,  $\varepsilon_h$  and  $\varepsilon_u$  are assumed to be identical to those of S355. The dam-  
294 age initiation criterion for each steel is given by Eq. 19 and their corresponding  
295 values tabulated in Table 6. Damage evolution is modelled as shown in Fig. 8 for  
296 all the four grades of steel because the actual nominal stress-strain curves were not  
297 provided by Girão Coelho (2004) for S235, S275 and S450. This is an acceptable  
298 assumption given that all four steels are characterised by similar  $\varepsilon_u$  and  $\varepsilon_f$  (Table  
299 5). The same Grade 8.8 bolt is used throughout and is modelled as previously  
300 described.

### 301 *Results and Discussions*

302 Results of the parametric study are plotted in Fig. 12. Analytical predictions  
303 by Piluso et al. (2001) and Francavilla et al. (2016) are then compared to the

304 FE predictions of the ultimate displacement  $\Delta_u$  and failure mode for  $\lambda_{\max} = 1.25$   
 305 in Table 4. Since the analytical models are, hitherto, mostly applied only to  
 306  $\lambda \approx 1$ , comparison will be made here for  $\lambda_{\max} = 1.25$ . The deformation capacity  
 307 of a T-stub is characterised by its non-dimensional ultimate displacement  $\delta \triangleq$   
 308  $\Delta_u/t_p$ , and this is plotted against  $\Gamma \triangleq \kappa f_y/f_{ub} \cdot t_p^2/A_s$ . Notice that  $\Gamma$  is a product  
 309 of two dimensionless groups that were previously used to delineate the régime  
 310 boundaries. From Eq. 13, it is clear that  $\Gamma = h(\lambda)$ ; hence, for a constant  $\lambda$ , the  
 311 régime boundaries depend only on  $\Gamma$ .

312 Figure 12 shows a general reduction in the deformation capacity  $\delta$  with  $\Gamma$   
 313 and/or when the mode switches from 1 to 3. It is hardly surprising that the  
 314 ductility of a T-stub in mode 1 is highest due to the collapse mechanism it develops.  
 315 And, since the collapse mechanism is affected by both geometric and material  
 316 parameters (EN 1993-1-8 2005), the ductility of T-stubs is highly dependent on  
 317 the dimensionless parameters  $t_p^2/A_s$  and  $\kappa f_y/f_{ub}$ . Notice that the data points for  
 318 mode 1 are much more disperse, further confirming the sensitivity of ductility to  
 319 geometric parameters and material properties. By contrast, mode 3 deformation  
 320 is dictated by the deformation of the bolts and is, consequently, less sensitive to  
 321 geometric and material properties of the flange.

322 Table 4 tabulates the value of  $\delta$  for each T-stub. It can be seen that, for  
 323 a constant  $\kappa f_y/f_{ub}$ ,  $\delta$  increases for T-stubs constructed with a weak flange and  
 324 strong bolts. These T-stubs fail predominantly in either mode 1 or 2. By contrast,  
 325 T-stubs with strong flanges and weak bolts deform primarily in mode 3, and they  
 326 have low deformation capacity ( $\delta$ ) that is nearly constant for  $\Gamma > 1.5$ . Furthermore,  
 327 Table 4 also shows that, for the same combination of  $t_p$  and  $A_s$ ,  $\delta$  reduces with  
 328 increasing  $\kappa f_y/f_{ub}$  for all modes of failure. In addition, the effects of  $\kappa f_y/f_{ub}$  on  $\delta$  is

329 greatest for  $\lambda = \lambda_{\min}$  (mode 1:  $5.3 \leq \delta \leq 6$  and mode 2:  $1.9 \leq \delta \leq 4.5$ ) compared  
 330 to  $\lambda_{\text{inter}}$  (mode 1:  $4.9 \leq \delta \leq 5.5$  and mode 2:  $1.9 \leq \delta \leq 4.0$ ) and  $\lambda_{\max}$  (mode  
 331 1:  $4.6 \leq \delta \leq 5.1$  and mode 2:  $1.4 \leq \delta \leq 3.6$ ). This is because when  $\lambda = \lambda_{\min}$ ,  
 332  $n$  is small compared to  $m$  which implies that the bolt spacing  $p$  is large. Girão  
 333 Coelho (2004) observed that a large  $p$  is responsible for a reduction of the T-stub  
 334 stiffness because the flange is not as rigidly constrained between the weld-toe and  
 335 bolt line compared to cases where  $\lambda = \lambda_{\max}$  ( $n \gg m$ ). Increasing stiffness of a  
 336 T-stub is accompanied by a consequential increase in its ductility – T-stubs with  
 337  $\lambda = \lambda_{\min}$  are more compliant – which leads to comparatively higher ductility of the  
 338 T-stub (Girão Coelho 2004). By contrast, the effects of both  $t_p^2/A_s$  and  $\kappa f_y/f_{ub}$   
 339 are negligible on  $\delta$  for large  $\Gamma$  as the data points eventually flatten out.

340 Table 4 also compares the FE results for  $\lambda = \lambda_{\max}$  to analytical predictions  
 341 by Piluso et al. (2001) and Francavilla et al. (2016). In general, both models  
 342 under-predict  $\delta$  in mode 1 because the deformation mechanism was assumed to be  
 343 two-dimensional. This simplification is only valid for a beam yield line pattern;  
 344 instead, a non-circular yield line pattern characterised by 3D effects develops in  
 345 the flange. The model by Piluso et al. (2001) over-predicts the displacement in  
 346 mode 2 because it neglects the compatibility condition between the elongation of  
 347 the bolt and the deformation of the flange; this was subsequently addressed by  
 348 Francavilla et al. (2016). Notwithstanding, discrepancies remain between FE and  
 349 analytical predictions in mode 2 because geometric non-linearities were neglected  
 350 in both analytical models. However, the difference between FE and analytical  
 351 predictions is small in mode 3, which suggests that both analytical models are,  
 352 in general, accurate if applied to T-stubs that undergo small displacements. The  
 353 good agreement in mode 3 is also, partly, because of the insensitivity of mode 3

354 deformation to the yield line pattern that develops.

### 355 **Failure deformation maps - validation**

356 The fidelity of the deformation maps is demonstrated for non-circular yield line  
357 pattern in Fig. 13 by plotting the data from the parametric study; and for the beam  
358 yield line pattern with experimental data from Girão Coelho (2004), Piluso et al.  
359 (2001), Bursi and Jaspart (1997) and Ribeiro et al. (2015). Plastic (dotted lines)  
360 and ultimate (solid lines) régime boundaries are plotted for  $\lambda_{\min}$  (Fig. 13a),  $\lambda_{\text{inter}}$   
361 (Fig. 13b) and  $\lambda_{\max}$  (Fig. 13c). Note that the boundaries corresponding to mode  
362 2→3 for  $\lambda_{\max}$  lie below the ones for  $\lambda_{\min}$ , similarly in Figs. 4 and 5. This is because  
363 T-stubs with a smaller  $\lambda$  value tend to be more ductile; consequently, they are more  
364 likely to fail either in mode 1 or 2. Only a limited combinations of  $t_p^2/A_s$  and  $f_y/f_{ub}$   
365 causes mode 3 failure. The plastic failure mode predicted by the map is, in general,  
366 conservative since T-stubs that fail in mode 1 (or 2) were predicted to fail in mode  
367 2 (or 3). This is unsurprising given that the régime boundaries in the plastic failure  
368 map were constructed within the constitutive framework of limit analysis (for the  
369 flange material). By contrast, the ultimate régime boundaries were constructed  
370 by assuming a linear piece-wise approximation of the stress-strain curve (Fig. 3)  
371 which gives a better prediction of the failure mode. The data points indicated  
372 by blue arrows in Fig. 13 – they were identified by their row number in Table 4  
373 and the flange material – are outliers due to the approximate nature of the piece-  
374 wise idealisation of the flange material. Apart from the outliers, the failure maps  
375 predict well the mode of deformation predicted by the parametric study.

376 The maps also demonstrate how failure mode is influenced by  $t_p^2/A_s$  and  $\kappa f_y/f_{ub}$ .  
377 Increasing  $t_p^2/A_s$  leads to a shift in the mode of failure from 1 to 3 which is evident

378 from the columns of data – each column corresponds to a constant  $\kappa f_y/f_{ub}$  – in  
 379 Fig. 13. Similarly, increasing  $\kappa f_y/f_{ub}$  at a constant  $t_p^2/A_s$  value also leads to a shift  
 380 towards a less ductile mode ( $1 \rightarrow 2$  or  $2 \rightarrow 3$ ). It is worth noting that since  $f_{ub}$   
 381 is identical in the parametric study, increasing  $\kappa f_y/f_{ub}$  corresponds to a stronger  
 382 flange. If the flange has an increased resistance to deformation, then the bolt is  
 383 more likely to fail even for lower  $t_p$  and higher  $A_s$  values - see, for example, data  
 384 points 19 (S450), 23 (S355) and 23 (S450) in Fig. 13.

385 To examine the effects of varying  $A_s$  whilst keeping  $t_p$  constant, consider the  
 386 T-stub configurations 1 to 4 listed in Table 4. Figure 13(c) shows that a T-stub  
 387 with a greater  $A_s$  tends to fail in mode 1 – as exemplified by rows of data points  
 388 labelled 3 and 4 – while a smaller  $A_s$  leads to failure in mode 2 or 3 (exemplified  
 389 by rows of data points 1 and 2). The plastic failure map – their régime boundaries  
 390 are plotted as dotted lines – suggests that one would need to increase the diameter  
 391 of the bolt in order for the T-stub to deform in mode 1. This is in contrast to what  
 392 the ultimate failure map would suggest. Hence, using the ultimate map prevents  
 393 the over-sizing of bolts which is a common, yet expensive, strategy adopted by  
 394 structural steel designers. If instead one focuses on configurations 1, 5 and 9  
 395 – they are characterised by a constant  $A_s$  – it is evident from Fig. 13(c) that  
 396 increasing  $t_p$  leads to a less ductile failure mode; see, for example, the rows of data  
 397 points labelled 5 and 9.

398 Table 7 compares the predicted deformation using the maps to existing exper-  
 399 imental test data. It is clear that the plastic régime boundaries are excessively  
 400 conservative as they tend to predict a mode 2 deformation for T-stubs failing in  
 401 mode 1. The ultimate boundaries are evidently more accurate and are capable of  
 402 subdividing the geometric ( $t_p^2/A_s$ ) and material ( $\kappa f_y/f_{ub}$ ) parameters space into

403 correct modes of failure that are consistent with the experimental data.

404 Finally, the maps are only as good as the theory (specifically, the constitutive  
405 idealisation of the flange material) used to construct them. But they are useful  
406 in spite of their inexactness for both designing and interpreting experiments, and  
407 in selecting and understanding the behaviour of T-stubs for engineering applica-  
408 tions. And, by identifying the places where data or theory are poor, they can be  
409 systematically improved.

## 410 CONCLUSIONS

411 Failure deformation maps were constructed for the plastic and ultimate failure  
412 of a T-stub with a single bolt-row in tension. The maps allow to avoid iterative  
413 pre-design calculations by condensing a large body of information within the 2D  
414 parameters space  $t_p^2/A_s - \kappa f_y/f_{ub}$ . It was found that the failure mode is sensitive  
415 to the non-dimensional parameters  $t_p^2/A_s$  and  $\kappa f_y/f_{ub}$ . The maps show that a  
416 ductile failure mode (mode 1) is induced by either decreasing  $t_p$  or increasing  $A_s$ .  
417 The fidelity of the maps is demonstrated through existing experimental data and  
418 through a FE parametric investigation. It was shown that the analytical models  
419 by Piluso et al. (2001) and Francavilla et al. (2016) under-predict the ultimate  
420 displacement in mode 1 arising from the assumption of a beam yield line pattern.  
421 Both models are also shown to be accurate if applied to T-stubs undergoing small  
422 displacements where geometric non-linearities are negligible.

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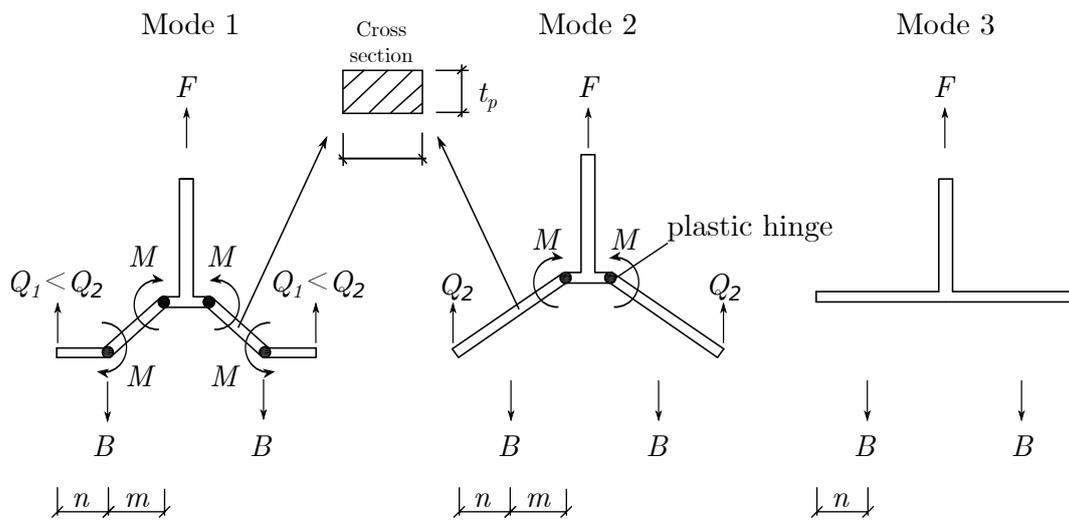
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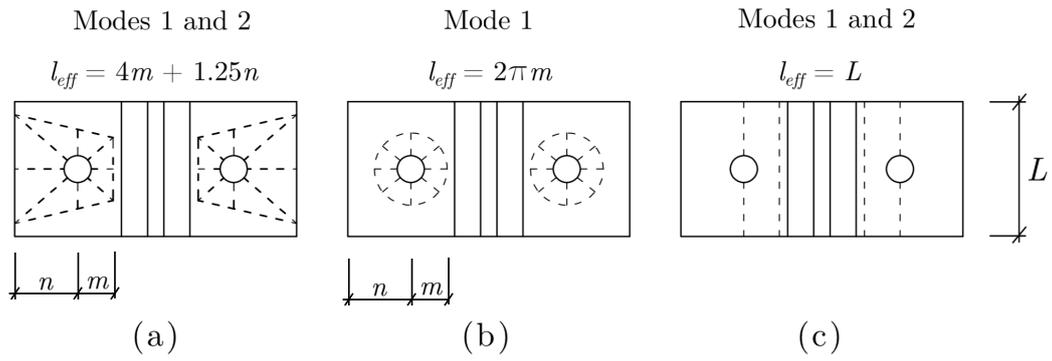
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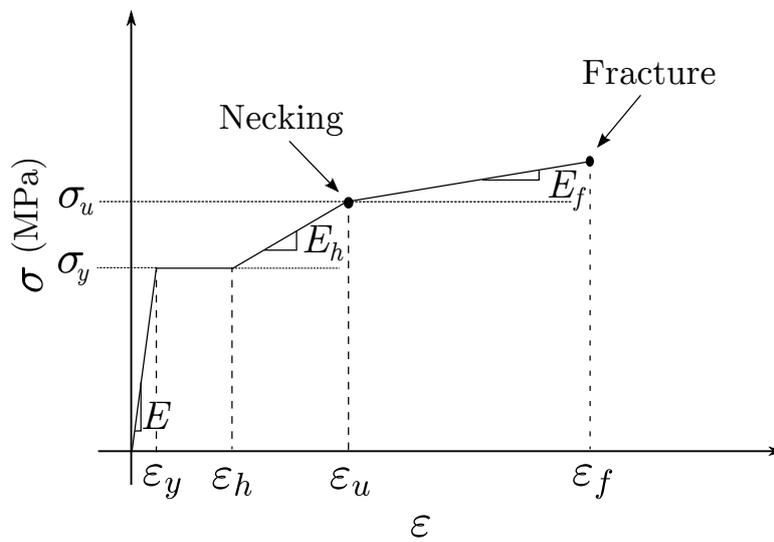
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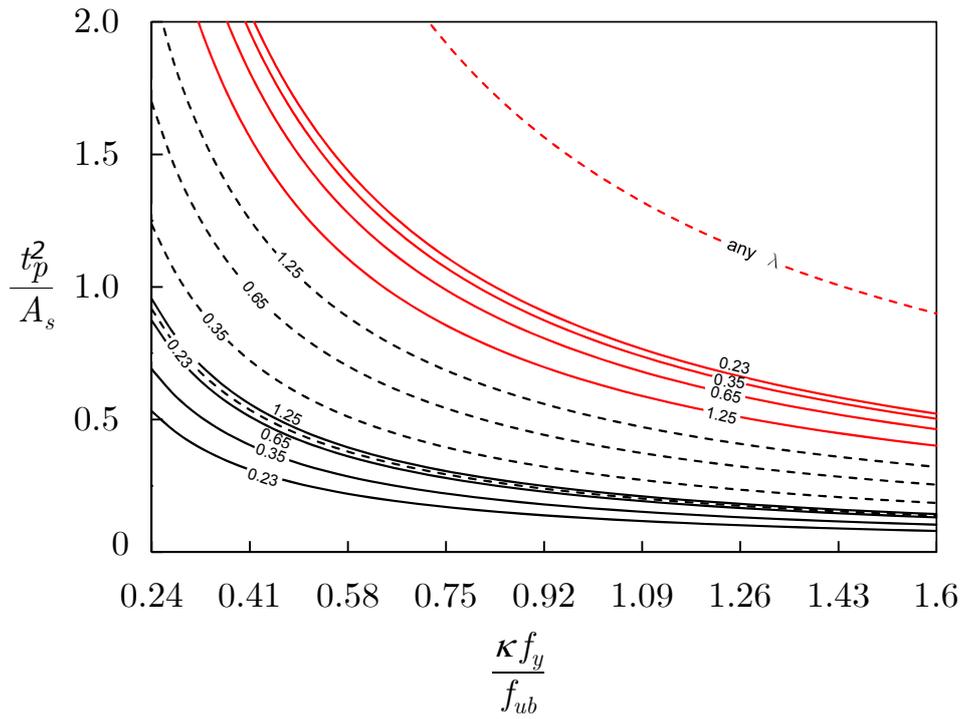
**FIG. 1. Schematic of failure modes adapted from Ribeiro et al. (2015).  $Q$  is prying force and  $B$  is tensile force in the bolt.**



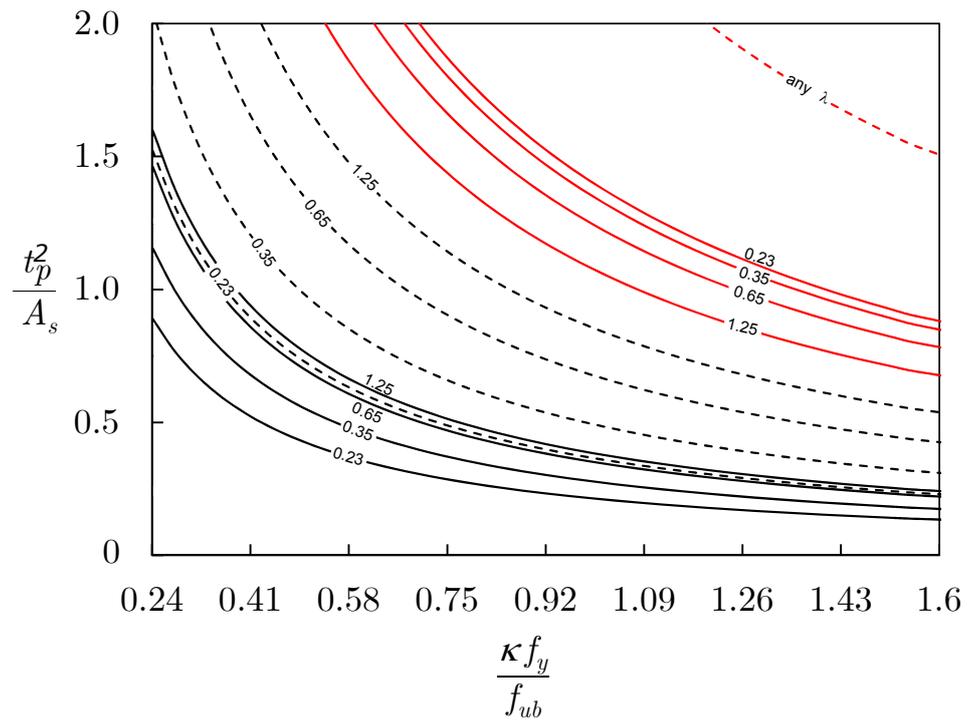
**FIG. 2. Possible yield line patterns in a T-stub: (a) non-circular; (b) circular; and (c) beam. - - - shows hinge line and  $L$  is the width (adapted from Girão Coelho (2004)).**



**FIG. 3. Idealised piece-wise approximation of the true stress-strain curve of the flange material (Piluso et al. 2001).**



**FIG. 4. Plastic failure map. Black and red isolines are régime boundaries at mode 1→2 and mode 2→3 transition, respectively. Solid lines for non-circular and dotted lines for beam yield line patterns.**



**FIG. 5. Ultimate failure map. Black and red isolines are régime boundaries at mode 1→2 and mode 2→3 transition, respectively. Solid lines for non-circular and dotted lines for beam yield line patterns.**

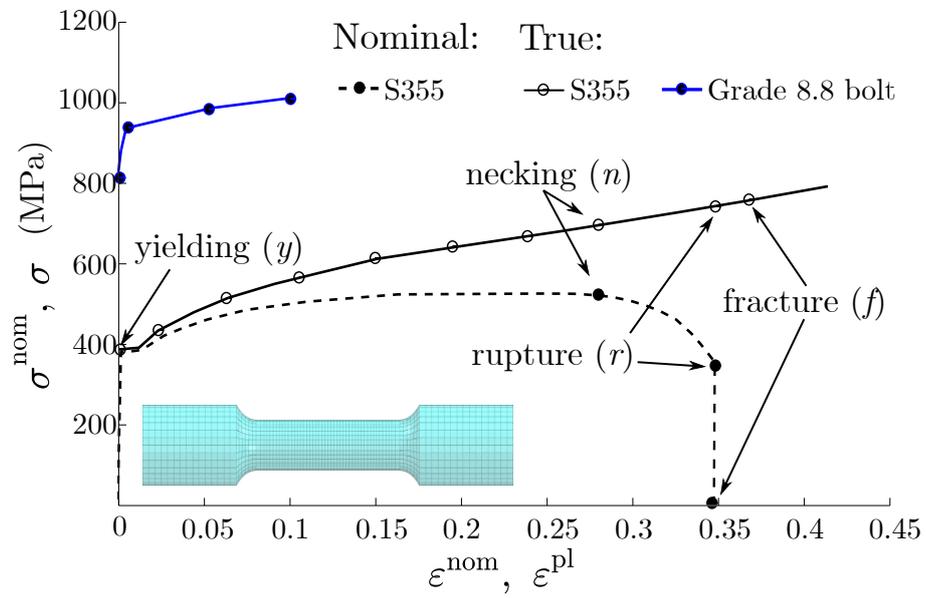
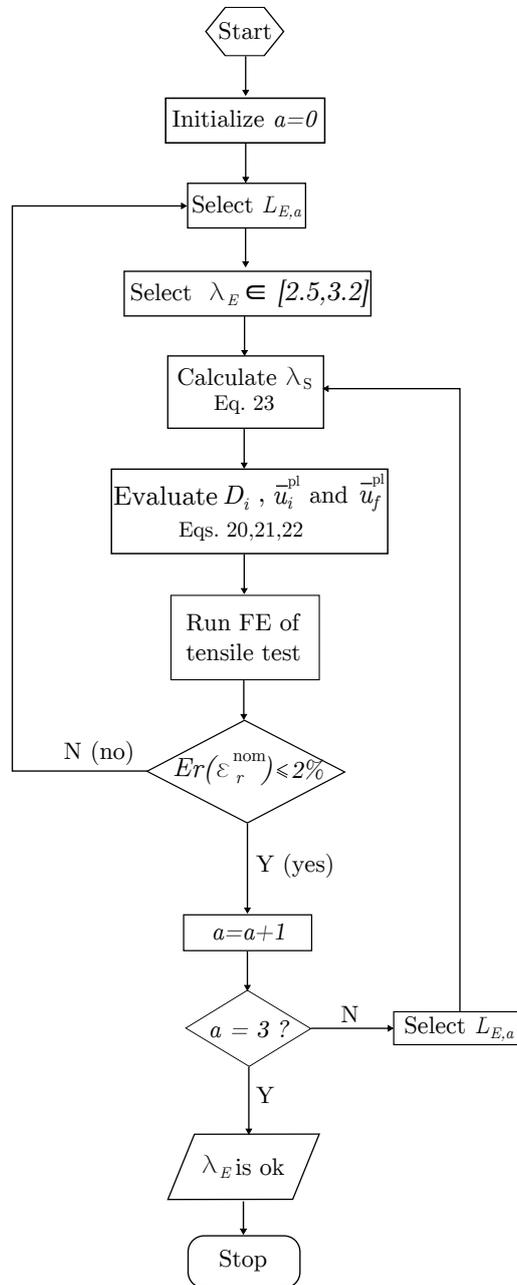


FIG. 6. Material properties for S355 steel and Grade 8.8 bolt.



**FIG. 7. Flow-chart on the procedure to calibrate  $\lambda_E$ .**

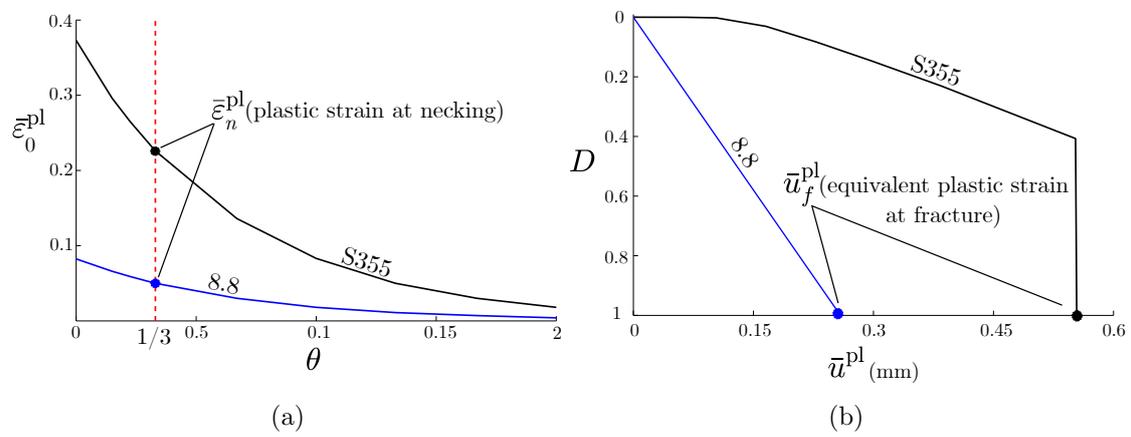
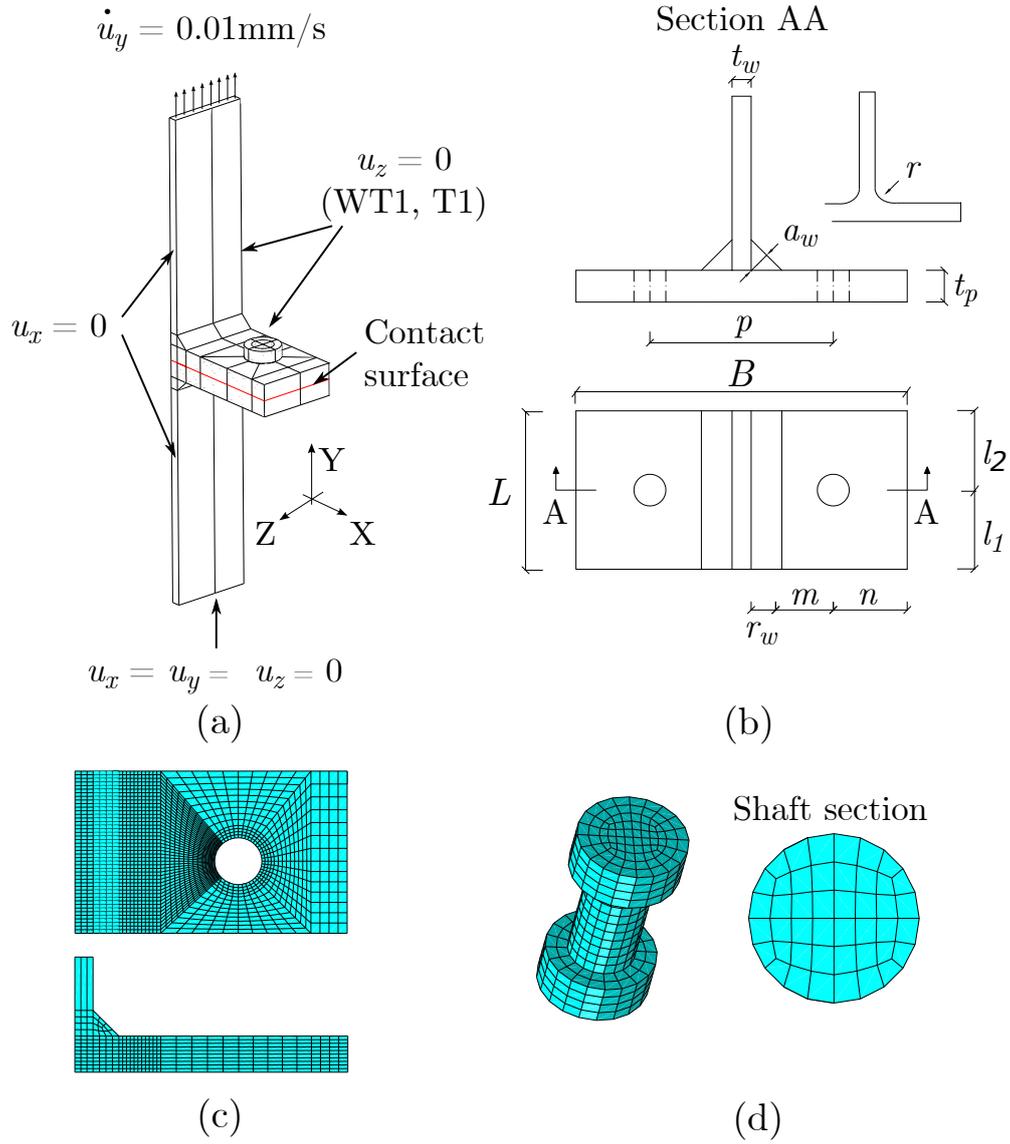
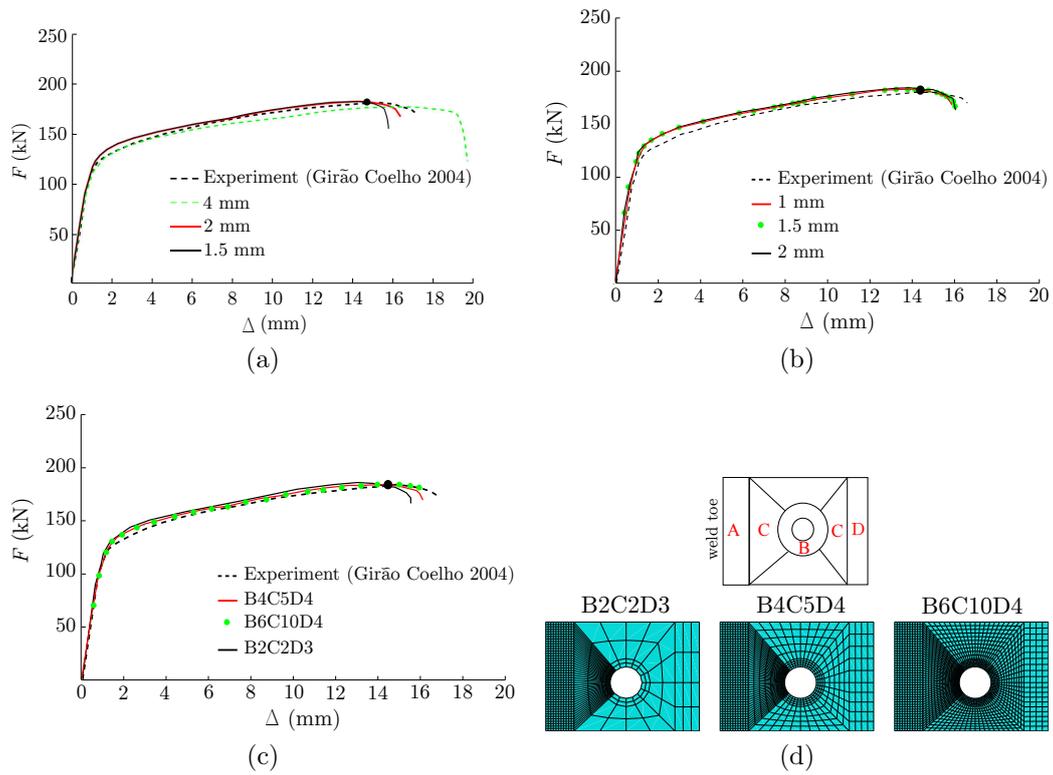


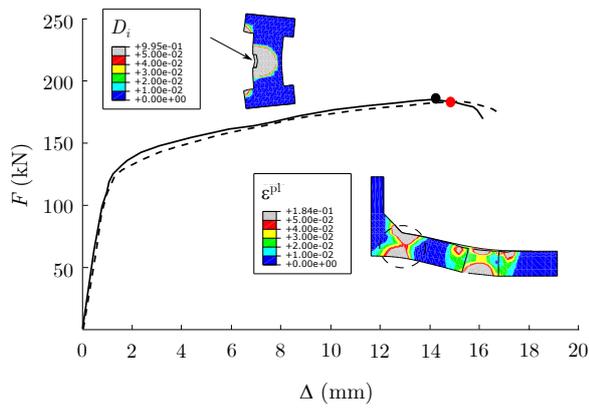
FIG. 8. (a) Damage initiation criterion and (b) damage evolution law.



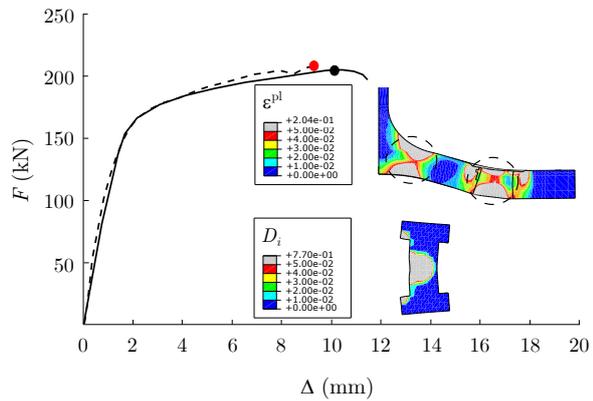
**FIG. 9. (a) FE model of the T-stubs and (b) its corresponding geometric parameters. The FE meshes for bolt and flange are shown in (c) and (d).  $r_w = 0.8a_w\sqrt{2}$  for WT1 and T15;  $r_w = 0.8r$  for T1.**



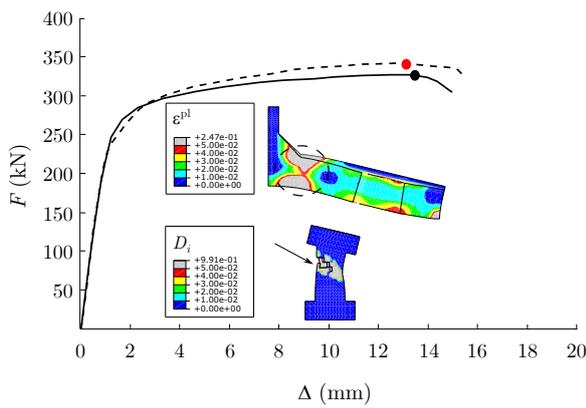
**FIG. 10. Results of mesh sensitivity study for WT1 (Girão Coelho 2004), where black dot indicates ultimate displacement  $\Delta_u$ : (a) effects of element size in the bolt shaft; (b) effects of element size at the weld toe; (c) predicted  $F - \Delta$  curves by modelling the flange as shown in (d).**



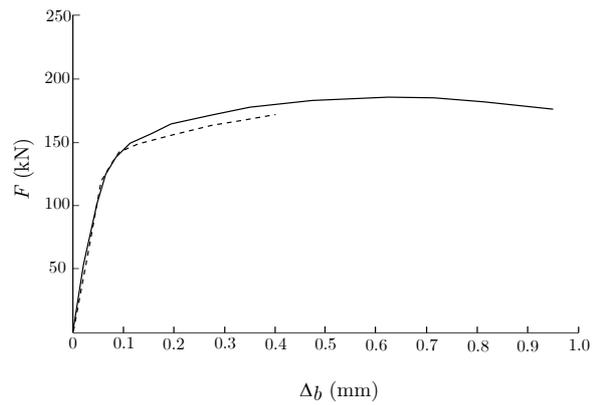
(a) WT1 (Girao Coelho 2004)



(b) T1 (Bursi and Jaspart 1997)

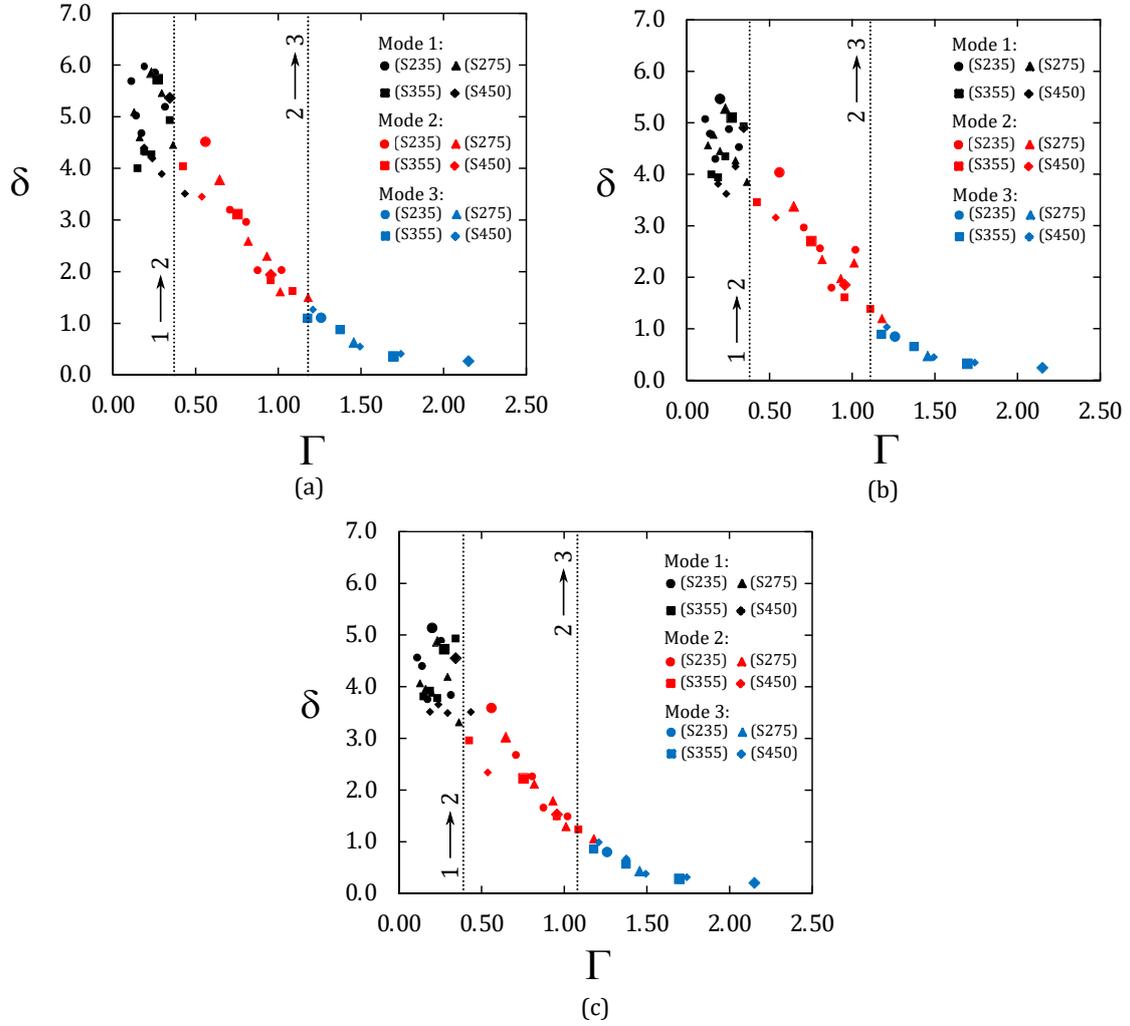


(c) T15 (Ribeiro et al. 2015)



(d) Force - bolt elongation (Girao Coelho 2004)

**FIG. 11. Comparison between FE (—) and experiments (- -). Arrow indicates deleted elements in the bolt and dotted circle indicates where plastic hinge develops.**



**FIG. 12. Results of parametric study: (a)  $\lambda_{\min} = 0.9$ , (b)  $\lambda_{\text{inter}} = 1.11$  and (c)  $\lambda_{\max} = 1.25$ . Vertical lines indicate boundaries between ultimate failure modes.**

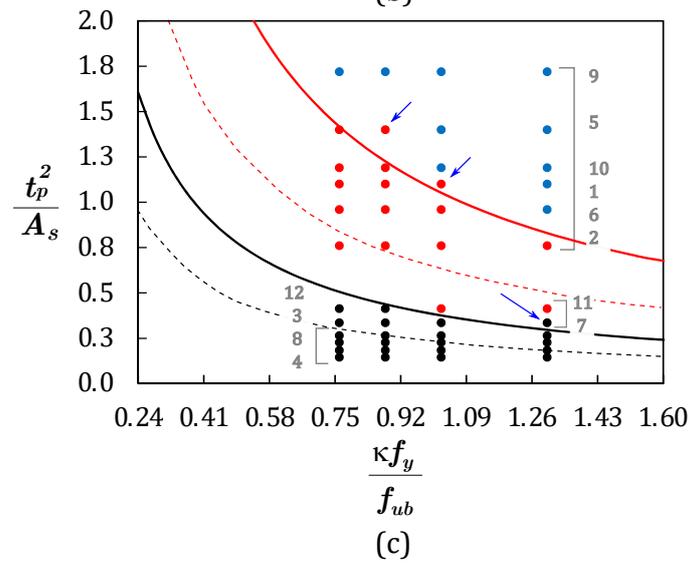
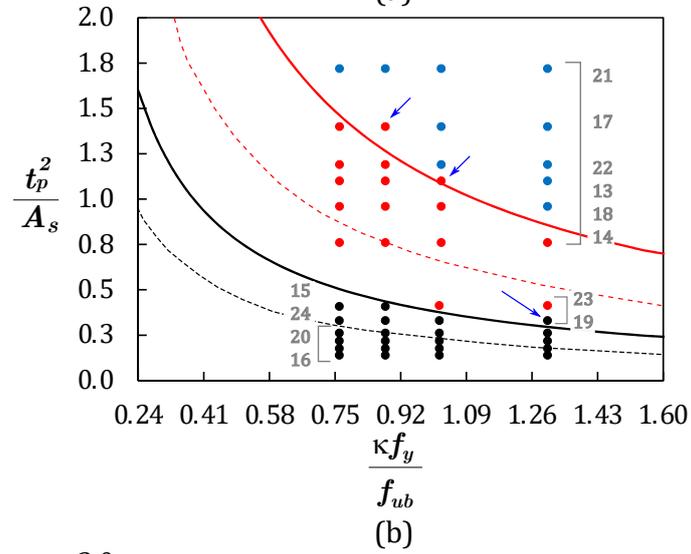
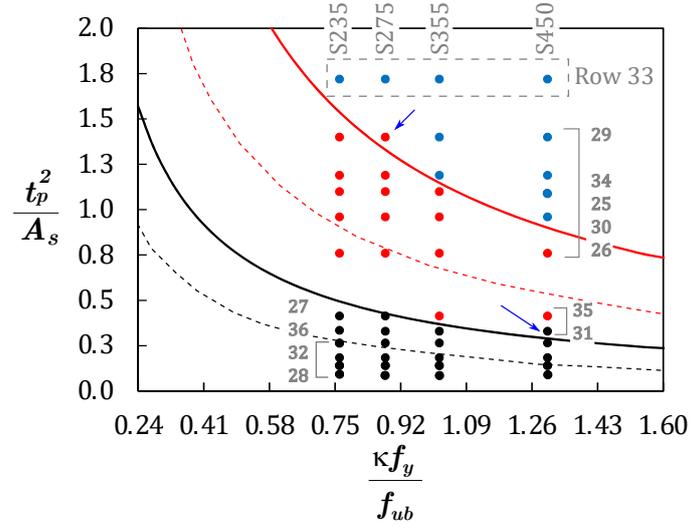


FIG. 13. Predicted failure mode by FE for (a)  $\lambda_{\min}=0.9$ , (b)  $\lambda_{\text{inter}}=1.11$  and (c)  $\lambda_{\max}=1.25$ . Mode 1 (●), Mode 2 (●) and Mode 3 (●). Grey numbers identify the geometry (Table 4) and flange materials are indicated.

Mode transition	$\beta$
1→2	$\frac{2\lambda}{1+2\lambda}$
2→3	2

TABLE 1: Values of  $\beta$  delineating the transition between failure modes given by Piluso et al. (2001).

		$h(\lambda)$		
Mode transition		Non-circular	Beam	$\kappa$
Plastic limit state	1→2	$\frac{3.57\lambda}{(1+2\lambda)(4+1.25\lambda)}$	$\frac{1.44\lambda}{(1+2\lambda)}$	1
	2→3	$\frac{3.57}{4+1.25\lambda}$	1.44	1
Ultimate limit state	1→2	$\frac{6\lambda}{(4+1.25\lambda)(1+2\lambda)}$	$\frac{2.4\lambda}{1+2\lambda}$	Eq. 10
	2→3	$\frac{6}{(4+1.25\lambda)}$	2.4	Eq. 10

TABLE 2: Expressions for  $h(\lambda)$  in Eq. 13.

Model	$m$	$n$	$t_p$	$d_b$	$l_1$	$l_2$	$L$	$B$	$p$	$a_w/r$	$t_w$	$r_w$
WT1	34.34	30	10	12	20	25	45	150.08	50	5	10	5.7
T1	29.45	30	10.7	12	20	20	40	150	90	15	7.1	12
T15	42.1	30	15	20	52.5	52.5	105	170	110	7	10	7.9

TABLE 3: Geometric dimensions corresponding to Fig. 9b. All dimensions are in mm except  $a_w/r$ .

Row	Geometry					Predicted $\delta$			
	m	n	$\lambda$	$t_p$	$d_b$	S235	S275	S355	S450
1	40	50	1.25	8	10	2.3(5.0,2.0)	1.8(3.8,1.6)	1.2 (1.4,0.7)	0.7(0.9,0.5)
2	40	50	1.25	8	12	3.6(5.0,2.3)	3.0(3.8,2.1)	2.2(3.6,2.6)	1.5(2.6,1.57)
3	40	50	1.25	8	20	5.1(5.2,5.1)	4.9(4.0,3.8)	4.7(3.9,3.7)	4.6(3.34,3.53)
4	40	50	1.25	8	27	4.6(5.2,5.0)	4.1(4.0,3.6)	3.8(3.9,3.5)	3.5(3.34,2.65)
5	40	50	1.25	9	10	1.5(1.9,1.4)	1.1(1.2,0.4)	0.6(0.3,0.3)	0.3(0.4,0.18)
6	40	50	1.25	9	12	2.7(4.0,1.4)	2.1(2.5,1.4)	1.5(1.5,1.4)	1.0(0.92,0.35)
7	40	50	1.25	9	20	4.9(4.1,4.1)	4.2(3.2,4.6)	4.9(3.1,3.7)	3.5(2.6,2.1)
8	40	50	1.25	9	27	4.4(4.1,3.9)	4.0(3.2,2.3)	3.9(3.1,2.7)	3.7(2.7,2.16)
9	40	50	1.25	10	10	0.8(0.3,0.2)	0.4(1.1,0.2)	0.3(0.5,0.3)	0.2(0.3,0.3)
10	40	50	1.25	10	12	1.7(3.2,1.4)	1.3(1.3,1.1)	0.9(0.7,0.3)	0.4(0.1,0.2)
11	40	50	1.25	10	20	3.8(3.4,3.9)	3.3(2.6,4.2)	3.0(3.1,2.5)	2.3(2.5,3.0)
12	40	50	1.25	10	27	4.1(3.4,3.2)	3.9(2.6,2.4)	3.8(2.6,2.3)	3.5(2.2,3.5)
13	44	46	1.1	8	10	2.6	2.01	1.4	0.9
14	44	46	1.1	8	12	4.0	3.4	2.7	2.1
15	44	46	1.1	8	20	5.5	5.3	5.1	4.9
16	44	46	1.1	8	27	5.1	4.6	4.0	3.8
17	44	46	1.1	9	10	2.5	2.3	0.5	0.3
18	44	46	1.1	9	12	3.0	1.6	1.3	0.3
19	44	46	1.1	9	20	5.6	5.4	5.3	5.0
20	44	46	1.1	9	27	4.8	4.8	3.9	3.6
21	44	46	1.1	10	10	1.2	0.9	0.5	0.3
22	44	46	1.1	10	12	1.8	2.3	0.8	0.7
23	44	46	1.1	10	20	4.5	3.9	1.9	3.5
24	44	46	1.1	10	27	4.7	4.5	4.4	4.2
25	47	43	0.9	8	10	3.0	2.3	1.6	1.1
26	47	43	0.9	8	12	4.5	3.8	3.1	1.9
27	47	43	0.9	8	20	6.0	5.9	5.7	5.4
28	47	43	0.9	8	27	5.7	5.1	4.0	4.4
29	47	43	0.9	9	10	2.0	2.6	0.3	0.6
30	47	43	0.9	9	12	3.2	1.8	1.3	0.4
31	47	43	0.9	9	20	5.9	5.5	4.9	3.5
32	47	43	0.9	9	27	5.0	4.6	4.3	4.2
33	47	43	0.9	10	10	1.5	0.6	0.4	1.1
34	47	43	0.9	10	12	2.0	1.6	1.0	0.9
35	47	43	0.9	10	20	5.2	4.5	4.0	3.5
36	47	43	0.9	10	27	4.7	4.3	4.3	3.9

TABLE 4: Dimensions (mm) of T-stubs - see Fig. 9b - considered in the parametric study.  $l_1 = l_2 = L/2 = 125\text{mm}$ ;  $B = 200\text{mm}$ ;  $t_w = 6.5\text{mm}$ ;  $a_w = 5\text{mm}$  and  $r_w = 7.5\text{mm}$  in all the models. XX (YY, ZZ) refers to FEM (Piluso et al. (2001) , Francavilla et al. (2016)).

Steel	$\sigma_y$ (MPa)	$\sigma_u$ (MPa)	$\varepsilon_y$	$\varepsilon_h$	$\varepsilon_u$	$\varepsilon_f$	$E$ (MPa)	$E_h$ (MPa)	$E_f$ (MPa)	$\kappa$
S235	235	360	0.001	0.014	0.036	0.25	210000	5500	360	2.587
S275	275	430	0.001	0.015	0.047	0.22	210000	4800	430	2.560
S355	355	510	0.002	0.017	0.053	0.2	210000	4250	510	2.305
S450	440	550	0.002	0.017	0.053	0.17	210000	4250	550	2.305

TABLE 5: Key stresses, strains, moduli and  $\kappa$  corresponding to Fig. 3.

S235		S275		S355		S450	
$\theta$	$\bar{\varepsilon}_0^{\text{pl}}$	$\theta$	$\bar{\varepsilon}_0^{\text{pl}}$	$\theta$	$\bar{\varepsilon}_0^{\text{pl}}$	$\theta$	$\bar{\varepsilon}_0^{\text{pl}}$
0.1	0.183	0.1	0.158	0.1	0.144	0.1	0.128
0.13	0.176	0.13	0.156	0.13	0.143	0.13	0.130
0.33	0.130	0.33	0.112	0.33	0.102	0.33	0.090
0.67	0.078	0.67	0.067	0.67	0.061	0.67	0.054
1	0.048	1	0.041	1	0.037	1	0.033
1.33	0.029	1.33	0.025	1.33	0.022	1.33	0.020
1.67	0.114	1.67	0.015	1.67	0.014	1.67	0.012
2	0.010	2	0.009	2	0.008	2	0.007

TABLE 6: Damage initiation criteria.

Model	Geometry			Material			Failure Mode		
	$t_p$	$d_b$	$\lambda$	$\kappa$	$f_y$	$f_{ub}$	Exp.	Plas.	Ult.
1	14.4	20	1.04	3.12	291	800	2	2	2
2	14.6	20	1.04	3.33	265	800	2	2	2
3	13	20	0.52	3.49	273	800	2	2	2
4	12.3	24	1.19	3.68	300	800	1	2	1
6	16.4	24	0.99	3.70	280	800	2	2	2
9	12.5	27	0.99	3.58	301	800	1	1	1
12	12.2	20	0.93	2.95	347	800	1	2	1
WT1	10	12	0.9	2.31	355	800	1	2	2
WT7_M20	10	20	0.9	3.33	355	800	1	2	1
WT7_M16	10	16	0.9	3.33	355	800	1	2	2
WT57_M16	10	16	0.9	1.74	690	800	2	2	2
WT57_M12	10	12	0.9	1.74	690	800	2	2	2
WT7_M12	10	12	0.9	3.33	355	800	2	2	2
T1			1.02				1	2	2
T15			0.7				2	2	2

TABLE 7: Comparison of predicted failure mode (by the deformation maps) against experiment data from (a) Piluso et al. (2001), (b) Girão Coelho (2004), (c) Bursi and Jaspart (1997) and (d) Ribeiro et al. (2015).