

First Results of the Search for Neutrinoless Double-Beta Decay with the NEMO 3 Detector

R. Arnold,⁵ C. Augier,⁸ J. Baker,⁴ A. Barabash,⁶ G. Broudin,¹ V. Brudanin,⁷ A. J. Caffrey,⁴ E. Caurier,⁵ V. Egorov,⁷ K. Errahmane,⁸ A. I. Etievre,⁸ J. L. Guyonnet,⁵ F. Hubert,¹ Ph. Hubert,¹ C. Jollet,⁵ S. Jullian,⁸ O. Kochetov,⁷ V. Kovalenko,⁷ S. Kononov,⁶ D. Lalanne,⁸ F. Leccia,¹ C. Longuemare,⁹ G. Lutter,¹ Ch. Marquet,¹ F. Mauger,⁹ F. Nowacki,⁵ H. Ohsumi,¹² F. Piquemal,¹ J. L. Reyss,¹⁰ R. Saakyan,¹³ X. Sarazin,⁸ L. Simard,⁸ F. Šimkovic,¹⁵ Yu. Shitov,⁷ A. Smolnikov,⁷ L. Štekl,³ J. Suhonen,¹⁴ C. S. Sutton,¹¹ G. Szklarz,⁸ J. Thomas,¹³ V. Timkin,⁷ V. Tretyak,⁷ V. Umatov,⁶ L. Vála,³ I. Vanushin,⁶ V. Vasilyev,⁶ V. Vorobel,² and Ts. Vylov⁷

¹CENBG, IN2P3-CNRS and UBI, 33170 Gradignan, France

²Charles University, Prague, Czech Republic

³IEAP, Czech Technical University, Prague, Czech Republic

⁴INL, Idaho Falls, Idaho 83415, USA

⁵IReS, IN2P3-CNRS and ULP, 67037 Strasbourg, France

⁶ITEP, 117259 Moscow, Russia

⁷JINR, 141980 Dubna, Russia

⁸LAL, IN2P3-CNRS and UPS, 91405 Orsay, France

⁹LPC, IN2P3-CNRS and UC, 14032 Caen, France

¹⁰LSCE, CNRS, 91190 Gif-sur-Yvette, France

¹¹MHC, South Hadley, Massachusetts 01075, USA

¹²Saga University, Saga 840-8502, Japan

¹³UCL, London WC1E 6BT, United Kingdom

¹⁴Jyväskylä University, FIN-40351 Jyväskylä, Finland

¹⁵FMFI, Comenius University, SK-842 48 Bratislava, Slovakia

(Received 19 July 2005; published 25 October 2005)

The NEMO 3 detector, which has been operating in the Fréjus underground laboratory since February 2003, is devoted to the search for neutrinoless double-beta decay ($\beta\beta 0\nu$). The half-lives of the two neutrino double-beta decay ($\beta\beta 2\nu$) have been measured for ^{100}Mo and ^{82}Se . After 389 effective days of data collection from February 2003 until September 2004 (phase I), no evidence for neutrinoless double-beta decay was found from ~ 7 kg of ^{100}Mo and ~ 1 kg of ^{82}Se . The corresponding limits are $T_{1/2}(\beta\beta 0\nu) > 4.6 \times 10^{23}$ yr for ^{100}Mo and $T_{1/2}(\beta\beta 0\nu) > 1.0 \times 10^{23}$ yr for ^{82}Se (90% C.L.). Depending on the nuclear matrix element calculation, the limits for the effective Majorana neutrino mass are $\langle m_\nu \rangle < 0.7\text{--}2.8$ eV for ^{100}Mo and $\langle m_\nu \rangle < 1.7\text{--}4.9$ eV for ^{82}Se .

DOI: 10.1103/PhysRevLett.95.182302

PACS numbers: 23.40.-s, 14.60.Pq, 21.10.Tg

The positive results obtained in the last few years in neutrino oscillation experiments [1–4] have demonstrated that neutrinos are massive particles and that lepton flavor is not conserved. In parallel, tritium β -decay experiments [5,6] have established a very low limit on the electron neutrino mass of $m_{\nu_e} < 2.2$ eV (95% C.L.). Grand unified theories can provide a natural framework for neutrino masses and lepton number violation. In particular, the seesaw model [7], which requires the existence of a Majorana neutrino, naturally explains the smallness of neutrino masses. The existence of Majorana neutrinos would also provide a natural framework for the leptogenesis mechanism [8] which could explain the observed baryon-antibaryon asymmetry in the Universe. The observation of neutrinoless double-beta decay ($\beta\beta 0\nu$) would prove that neutrinos are Majorana particles and that global lepton number is not conserved. It would also constrain the mass spectrum and the absolute mass of the neutrinos.

The NEMO 3 detector [9], installed in the Fréjus underground laboratory (LSM, France) is searching for $\beta\beta 0\nu$ decay by the direct detection of the two electrons with a combination of tracking and calorimeter information. The

two main isotopes present inside the detector in the form of very thin foils (40–60 mg/cm²) are ^{100}Mo (6914 g, $Q_{\beta\beta} = 3034$ keV) and ^{82}Se (932 g, $Q_{\beta\beta} = 2295$ keV). On both sides of the sources, there is a gaseous tracking detector which consists of 6180 open drift cells operating in the Geiger mode allowing three-dimensional track reconstruction. To minimize the multiple scattering, the gas is a mixture of 95% helium, 4% ethyl alcohol, 1% argon, and 0.1% water. The wire chamber is surrounded by a calorimeter which consists of 1940 plastic scintillator blocks coupled to very low radioactive photomultipliers (PMTs). The energy resolution (FWHM) of the calorimeter is 14% at 1 MeV for the scintillators equipped with the 5" PMTs on the external wall and 17% for the 3" PMTs on the internal wall. The resolution of the summed energy of the two electrons in the $\beta\beta 0\nu$ decay is mainly a convolution of the energy resolution of the calorimeter and the fluctuation in the electron energy loss in the foil source which gives a non-Gaussian tail. The FWHM of the expected two-electron energy spectrum of the $\beta\beta 0\nu$ decay is 350 keV. Absolute energy calibrations are carried out every 40 days using ^{207}Bi sources. A daily laser survey controls

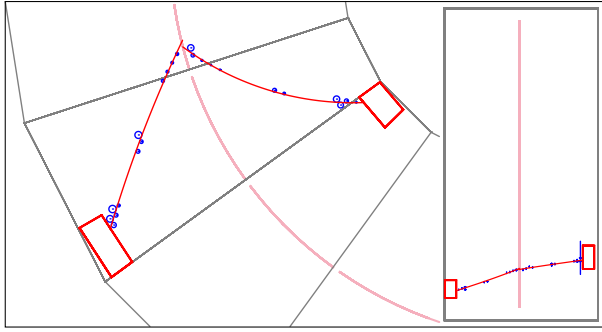


FIG. 1 (color online). Transverse (left) and longitudinal (right) view of a reconstructed $\beta\beta$ event selected from the data with a two-electron energy sum of 2812 keV.

the gain stability of each PMT. A solenoid surrounding the detector produces a 25 gauss magnetic field in order to distinguish electrons from positrons. An external shield of 18 cm of low radioactivity iron, a water shield, and a wood shield cover the detector to reduce external γ and neutrons.

A two-electron ($2e^-$) event (see Fig. 1) candidate for a $\beta\beta$ decay is defined as follows: two tracks come from the same vertex on the source foil, each track must be associated with a fired scintillator, its curvature must correspond to a negative charge, and the time of flight must correspond to the two electrons being emitted from the same source position. For each electron an energy threshold of 200 keV for ^{100}Mo and 300 keV for ^{82}Se is applied. Figures 2(a) and 3 show the two-electron energy sum spectra after background subtraction obtained after 389 effective days of data collection with ^{100}Mo and ^{82}Se , respectively. The angular distribution of the two electrons and the single energy spectrum are also presented in the case of ^{100}Mo in Figs. 2(b) and 2(c). All these spectra are in good agreement with the $\beta\beta 2\nu$ simulations. The values of the measured half-lives are $T_{1/2}(\beta\beta 2\nu) = [7.11 \pm 0.02(\text{stat}) \pm 0.54(\text{syst})] \times 10^{18}$ yr for ^{100}Mo (with a single state dominance decay) and $[9.6 \pm 0.3(\text{stat}) \pm 1.0(\text{syst})] \times 10^{19}$ yr for ^{82}Se . These values are in agree-

ment with (but have a higher precision than) the previous measurements [10].

A complete study of the background in the $\beta\beta 0\nu$ energy window has been performed. The level of each background component has been directly measured using different analysis channels in the data.

External backgrounds due to ^{214}Bi and ^{208}Tl contaminants outside the source foils (mostly in the PMTs) have been measured by searching for Compton electrons emitted from the source foils by external γ . For ^{208}Tl , a total activity of ~ 40 Bq has been measured and is in agreement with the previous high purity germanium (HPGe) measurements of samples of the PMT glass. For ^{214}Bi , an activity of ~ 300 Bq has been found, again in agreement with the HPGe measurements of PMTs and also the level of radon surrounding the detector inside the shield. The expected number of $\beta\beta 0\nu$ -like events due to this background is negligible, $\leq 10^{-3}$ counts $\cdot \text{kg}^{-1} \cdot \text{yr}^{-1}$ in the $[2.8\text{--}3.2]$ MeV energy window where the $\beta\beta 0\nu$ signal is expected.

External neutrons and high energy γ backgrounds have been measured by searching for crossing electron events above 4 MeV. This corresponds to a negligible expected level of background of $\sim 3 \times 10^{-3}$ counts $\cdot \text{kg}^{-1} \cdot \text{yr}^{-1}$ in the $\beta\beta 0\nu$ energy window.

The level of ^{208}Tl impurities inside the sources has been measured by searching for internal ($e^- \gamma \gamma$) and ($e^- \gamma \gamma \gamma$) events. The measured activity is 80 ± 20 $\mu\text{Bq}/\text{kg}$ in molybdenum and 300 ± 50 $\mu\text{Bq}/\text{kg}$ in selenium. It is in agreement with the previous HPGe measurements which gave an upper limit of 100 $\mu\text{Bq}/\text{kg}$ for molybdenum and a positive measurement of 400 ± 100 $\mu\text{Bq}/\text{kg}$ for selenium. This corresponds to an expected level of background in the $\beta\beta 0\nu$ energy window of ~ 0.1 counts $\cdot \text{kg}^{-1} \cdot \text{yr}^{-1}$ for molybdenum and ~ 0.3 counts $\cdot \text{kg}^{-1} \cdot \text{yr}^{-1}$ for selenium. The measurement of ^{214}Bi impurities inside the sources could not be achieved in this first period of data due to radon contamination (see later). However, the previous HPGe detectors gave an upper limit of 350 $\mu\text{Bq}/\text{kg}$ for molybdenum and a positive measurement of

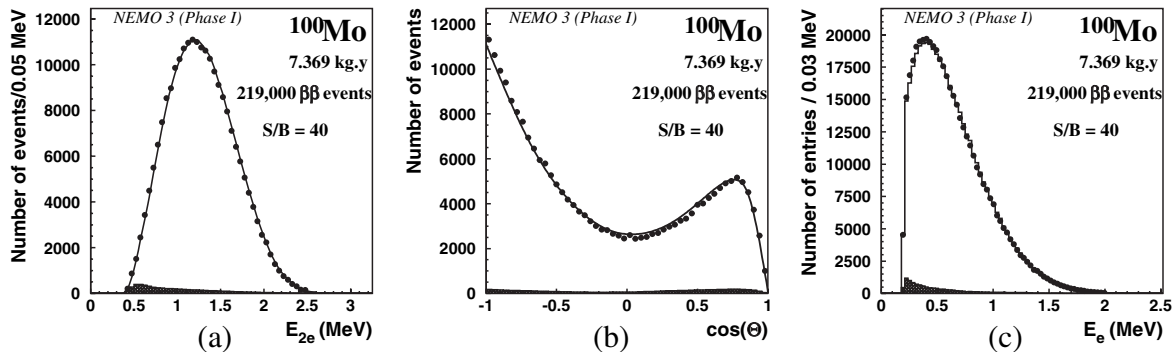


FIG. 2. (a) Energy sum spectrum of the two electrons, (b) angular distribution of the two electrons, and (c) single energy spectrum of the electrons, after background subtraction from ^{100}Mo with 7.369 kg \cdot yr exposure. The solid line corresponds to the expected spectrum from $\beta\beta 2\nu$ simulations and the shaded histogram is the subtracted background computed by Monte Carlo simulations. The signal contains 219 000 $\beta\beta$ events and the signal-to-background ratio is 40.

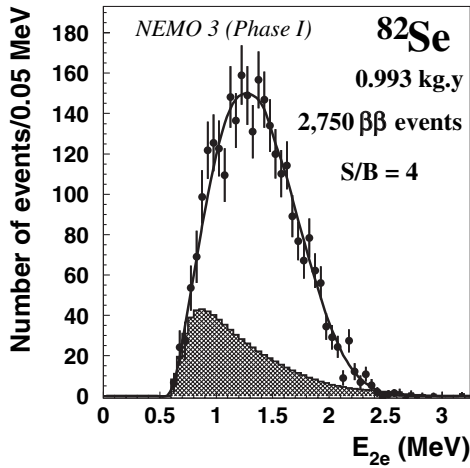


FIG. 3. Energy sum spectrum of the two electrons after background subtraction from ^{82}Se with 0.993 kg · yr exposure (same legend as Fig. 2). The signal contains 2750 $\beta\beta$ events and the signal-to-background ratio is 4.

1.2 ± 0.5 mBq/kg for selenium, corresponding to a negligible expected level of background.

The expected level of background due to the tail of the $\beta\beta 2\nu$ distribution in the $\beta\beta 0\nu$ energy window is ~ 0.3 counts · kg $^{-1}$ · yr $^{-1}$ for molybdenum and ~ 0.2 counts · kg $^{-1}$ · yr $^{-1}$ for selenium.

The dominant background in this first period of data was radon gas inside the tracking chamber due to a low rate of diffusion of radon from the laboratory (~ 15 Bq/m 3) into the detector. Two independent measurements of the radon level in the detector were carried out. The first used a high sensitivity radon detector similar to the one developed by the Super-Kamiokande collaboration [11]. The second was done by searching for (e^- , delayed- α) events in the NEMO 3 data. Indeed the tracking detector allows the detection of the delayed tracks (up to 700 μs later) in order to tag delayed α emitted by ^{214}Po in the Bi-Po process.

Both measurements are in good agreement and indicate a level of radon inside the detector of 25 ± 5 mBq/m 3 . This radon contamination corresponds to an expected level of background in the $\beta\beta 0\nu$ energy window of ~ 1 count · kg $^{-1}$ · yr $^{-1}$.

Figures 4(a) and 4(b) show the tail of the two-electron energy sum spectrum in the $\beta\beta 0\nu$ energy window for ^{100}Mo and for ^{82}Se , respectively. The number of $2e^-$ events observed in the data is in agreement with the expected number of events from $\beta\beta 2\nu$ and the radon simulations. For ^{100}Mo , in the energy window [2.8–3.2] MeV, the expected background is 8.1 ± 1.3 (error dominated by the uncertainty on the radon activity) and 7 events have been observed. For ^{82}Se , in the energy window [2.7–3.2] MeV, the expected background is 3.1 ± 0.6 and 5 events have been observed. In order to independently check the dominant radon contribution above 2.8 MeV, the energy sum spectrum [Fig. 4(c)] has been plotted for the two electrons emitted from the copper and tellurium foils where no background except radon is expected. The data are in agreement with the radon simulations.

In order to make the optimum use of all the information from the NEMO 3 detector, a maximum likelihood analysis [12] has been applied to the $2e^-$ event sample above 2 MeV using the three available variables: the energy sum (E_{tot}) of the two electrons, the energy of each electron (E_{min} is the minimum electron energy), and the angle between the two tracks ($\cos\theta$). A three-dimensional probability distribution function P^{3D} can be written as follows: $P^{3D} = P(E_{\text{tot}})P(E_{\text{min}}/E_{\text{tot}})P(\cos\theta/E_{\text{min}})$, where $P(E_{\text{min}}/E_{\text{tot}})$ and $P(\cos\theta/E_{\text{min}})$ are two conditional probability distribution functions. The likelihood is defined as $\mathcal{L} = \prod_{i=1}^{N_{\text{tot}}} (\sum_{k=1}^8 x_k P_k^{3D})$, where k corresponds to one of the eight contributions: $\beta\beta 0\nu$, $\beta\beta 2\nu$, radon, external and internal ^{207}Bi and ^{208}Tl , and neutrons. Here x_k is the ratio of the number of $2e^-$ events due to the process k relative to the total number of observed events N_{tot} . Finally, P_k^{3D} is built using simulated events of contribution k .

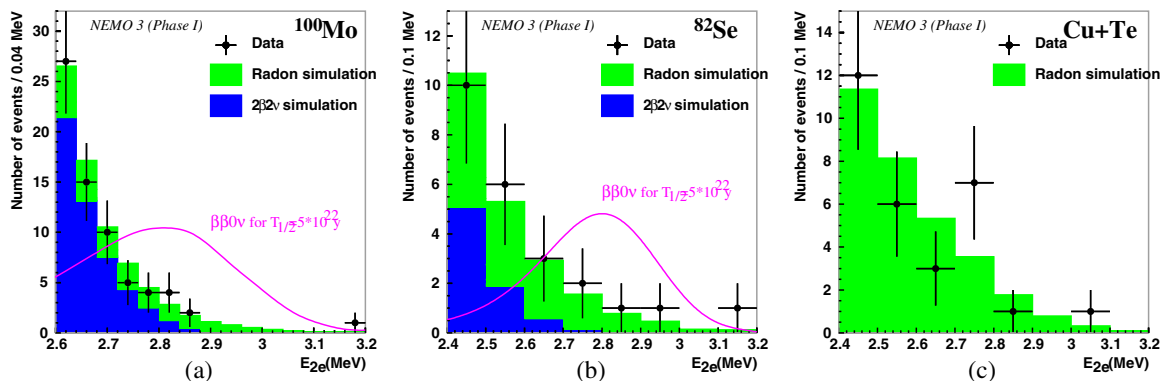


FIG. 4 (color online). Spectra of the energy sum of the two electrons in the $\beta\beta 0\nu$ energy window after 389 effective days of data collection from February 2003 until September 2004 (phase I): (a) with 6.914 kg of ^{100}Mo ; (b) with 0.932 kg of ^{82}Se ; (c) with copper and tellurium foils. The shaded histograms are the expected backgrounds computed by Monte Carlo simulations: dark grey (blue) is the $\beta\beta 2\nu$ contribution and light grey (green) is the radon contribution. The solid line corresponds to the expected $\beta\beta 0\nu$ signal for $T_{1/2}(\beta\beta 0\nu) = 5 \times 10^{22}$ yr.

TABLE I. Limits (in eV) on the effective neutrino mass $\langle m_\nu \rangle$ obtained from different theoretical calculations of nuclear matrix elements with $T_{1/2}(\beta\beta 0\nu) > 4.6 \times 10^{23}$ yr for ^{100}Mo and $T_{1/2}(\beta\beta 0\nu) > 1.0 \times 10^{23}$ yr for ^{82}Se . QRPA is quasi-random-phase approximation.

Nuclear matrix elements		^{100}Mo	^{82}Se
Shell model	Caurier 1996 [15]		<4.9
QRPA	Rodin 2005 [16]	<2.7–2.8	<4.1–4.5
QRPA	Simkovic 1999 [17]	<1.0	<3.3
QRPA	Suhonen 2003 [18,19]	<1.1	<2.8–4.2
QRPA	Stoica 2001 [20]	<0.7–1.1	<1.7–3.7

With 389 effective days of data collection, limits at 90% C.L. obtained with the likelihood analysis are $T_{1/2}(\beta\beta 0\nu) > 4.6 \times 10^{23}$ yr for ^{100}Mo and 1.0×10^{23} yr for ^{82}Se . These limits are about 10 times higher than the previous limits obtained with ^{100}Mo and ^{82}Se [13,14]. The corresponding upper limits for the effective Majorana neutrino mass range from 0.7 to 2.8 eV for ^{100}Mo and 1.7 to 4.9 eV for ^{82}Se depending on the nuclear matrix element calculation [15–20]. Results for each calculation are given in Table I. For ^{100}Mo , since an incorrect value of the phase-space factor has been used in Ref. [20], the value calculated in Ref. [21] has been used. The claim of a positive $\beta\beta 0\nu$ signal observed with ^{76}Ge [22] gives an allowed effective mass range 0.1–0.9 eV. Our limit obtained with ^{100}Mo slightly overlaps this range. In the hypothesis of gluino or neutralino exchange, and using the nuclear matrix elements calculated in Ref. [23], limits obtained on the trilinear R -parity-violating supersymmetric coupling are $\lambda'_{111} < 1.6 \times 10^{-4}$ for ^{100}Mo and 3.0×10^{-4} for ^{82}Se . In the hypothesis of a right-handed weak current, the limits are $T_{1/2}(\beta\beta 0\nu) > 1.7 \times 10^{23}$ yr at 90% C.L. for ^{100}Mo and 0.7×10^{23} yr for ^{82}Se , corresponding to an upper limit on the coupling constant of $\lambda < 2.5 \times 10^{-6}$ for ^{100}Mo and 3.8×10^{-6} for ^{82}Se using the nuclear calculations from Refs. [18,24].

In conclusion, the NEMO 3 detector has been running reliably since February 2003. The $\beta\beta 2\nu$ decay has been measured for ^{82}Se and ^{100}Mo with very high statistics and better precision than the previous measurements. The two-electron energy sum spectrum, the single energy spectrum, and the angular distribution are all in good agreement with the $\beta\beta 2\nu$ simulations. All components of the background in the $\beta\beta 0\nu$ energy window have been measured directly using different analysis channels in the data. After 389 effective days of data collection, no evidence for $\beta\beta 0\nu$ decay has been found in ^{100}Mo or ^{82}Se . The limits at the 90% C.L. are $T_{1/2}(\beta\beta 0\nu) > 4.6 \times 10^{23}$ yr for ^{100}Mo and 1.0×10^{23} yr for ^{82}Se . For this first running period (phase I) presented here, radon was the dominant background at a level of about 3 times higher than the $\beta\beta 2\nu$ background for ^{100}Mo . It has now been significantly reduced by a factor ~ 10 by a radon-tight tent enclosing the

detector and a radon-trap facility in operation since December 2004 which has started a second running period (phase II). After five years of data collection, the expected sensitivity at 90% C.L. will be $T_{1/2}(\beta\beta 0\nu) > 2 \times 10^{24}$ yr for ^{100}Mo and 8×10^{23} yr for ^{82}Se , corresponding to $\langle m_\nu \rangle < 0.3\text{--}1.3$ eV for ^{100}Mo and $\langle m_\nu \rangle < 0.6\text{--}1.7$ eV for ^{82}Se .

A portion of this work was supported by grants from INTAS (03051-3431, 03-55-1689) and NATO (PST CLG 980022).

-
- [1] Y. Fukuda *et al.*, Phys. Rev. Lett. **86**, 5651 (2001).
 - [2] Y. Fukuda *et al.*, Phys. Rev. Lett. **81**, 1562 (1998).
 - [3] Q.R. Ahmad *et al.*, Phys. Rev. Lett. **89**, 011301 (2002).
 - [4] T. Araki *et al.*, Phys. Rev. Lett. **94**, 081801 (2005).
 - [5] C. Weinheimer *et al.*, Nucl. Phys. B, Proc. Suppl. **118**, 279 (2003).
 - [6] V. Lobashev *et al.*, Nucl. Phys. B, Proc. Suppl. **91**, 280 (2001).
 - [7] R.N. Mohapatra *et al.*, Phys. Rev. Lett. **44**, 912 (1980).
 - [8] M. Fukugita *et al.*, Phys. Lett. B **174**, 45 (1986).
 - [9] R. Arnold *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **536**, 79 (2005).
 - [10] K. Hagiwara *et al.* (Particle Data Group), Phys. Rev. D **66**, 010001 (2002).
 - [11] Y. Takeuchi *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **421**, 334 (1999).
 - [12] A. I. Etiennevre, Ph.D. thesis, University Paris-Sud, 2003.
 - [13] H. Ejiri *et al.*, Phys. Rev. C **63**, 065501 (2001).
 - [14] R. Arnold *et al.*, Nucl. Phys. A **636**, 209 (1998).
 - [15] E. Caurier *et al.*, Phys. Rev. Lett. **77**, 1954 (1996).
 - [16] V.A. Rodin *et al.*, Phys. Rev. C **68**, 044302 (2003).
 - [17] F. Simkovic *et al.*, Phys. Rev. C **60**, 055502 (1999).
 - [18] M. Aunola *et al.*, Nucl. Phys. A **643**, 207 (1998).
 - [19] J. Suhonen *et al.*, Nucl. Phys. A **723**, 271 (2003).
 - [20] S. Stoica *et al.*, Nucl. Phys. A **694**, 269 (2001).
 - [21] F. Boehm and P. Vogel, in *Physics of Massive Neutrinos* (Cambridge University Press, Cambridge, England, 1992), 2nd ed.
 - [22] H.V. Klapdor-Kleingrothaus *et al.*, Nucl. Instrum. Methods Phys. Res., Sect. A **522**, 371 (2004).
 - [23] A. Faessler *et al.*, Phys. Rev. D **58**, 115004 (1998).
 - [24] J. Suhonen, Nucl. Phys. A **700**, 649 (2002).