

1 **Physical properties controlling water repellency in synthesized**
2 **granular solids**

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13 *Running title: Synthesized water repellent granular solids*

14

15 **Summary**

16

17 The wettability of granular solids such as soil is known to depend primarily on two
18 factors: their inherent chemistry and their physical properties such as their particle
19 size, particle shape and surface roughness. Nevertheless, the distinctive physical
20 properties of such materials have not been fully explored to gauge their wettability. In
21 this study, the difference in wettability between a flat solid (microscope slide) and
22 three granular solids, namely: glass beads (GB), Leighton Buzzard Sand (LBS) and
23 crushed Glass (CG) which have different physical properties were examined. The
24 effect of chemistry was isolated by strongly hydrophobizing the above materials by
25 treatment with dimethyldichlorosilane. Wettability measurements were made by
26 measuring the water–solid contact angle (CA) by the sessile drop method after
27 adhering one-layer of uniformly-oriented granular solids onto double-sided adhesive
28 tape initially attached to a microscope slide. Techniques for particle characterization
29 included sieving for particle size, dynamic image analysis for particle shape and
30 confocal laser microscopy to determine surface roughness. Results show that all CAs
31 of the granular solids exceeded that of the hydrophobized microscope slide (103°).
32 The crushed glass had the largest CA (125°). With all three granular solids, there was
33 an increase in CAs as particle size decreased. In addition, as particles became more

34 angular, CAs increased. The influence of shape on wettability became more
35 predominant as particle size decreased. The surface roughness parameter, R_a , was
36 investigated and shown to be sensitive to both the size and shape of the particles. A
37 decrease in R_a from 95.4 to 34.1 μm increased CAs from 107 to 125°. A similar
38 change in CA was shown to correspond to an increase in void fraction from 40.7 to
39 77.4 %. Our results have practical implications for the optimum use of soil by
40 enhancing or suppressing water repellency.

41

42 *Keywords: wetting, contact angle, hydrophobicity*

43

44 **Highlights**

45

- 46 • **How do the physical properties of granular solids such as soil influence**
47 **their wettability?**
- 48 • **Effects of physical properties of particles on CAs were investigated when**
49 **isolated from effect of chemistry.**
- 50 • **Particle shape becomes more important in gauging wettability as particle**
51 **size decreases.**

- 52 • **Wettability of granular solids may be physically controlled by specific**
- 53 **particle characteristics.**

54 **Introduction**

55

56 Soil water repellency is known to influence hydrological processes such as infiltration
57 (de Jonge *et al.*, 2007). A small to medium degree of water repellency in soil is
58 important for the stability of aggregates (Hallett & Young, 1999; Goebel *et al.*, 2012),
59 but severe water repellency has been reported to lead to erosion by water (Shakesby *et*
60 *al.*, 2000) and wind (Ravi *et al.*, 2009). The wetting properties of solids are
61 characterized by their contact angle (CA), with a lower and upper boundary of 0° and
62 180° respectively. Granular solids such as soil show relatively larger CAs than flat
63 solids. Some reported values of maximum CAs of granular solids hydrophobized by
64 treatment with silane compounds were 126° (Bachmann & McHale, 2009), 127° (Liu
65 *et al.*, 2012) and 131° (Chan & Lourenço, 2016). McHale *et al.* (2005) also
66 demonstrated that CAs of synthesized water repellent sand by fluorine-based
67 compounds can reach values > 130°. In contrast, CAs of flat surfaces treated with
68 silane-based compounds do not exceed 108° in general (Gao & McCarthy, 2006).
69 Contact angles reported within the past decade on such surfaces were 108° (Cheng *et*
70 *al.*, 2014), 103° (Zhang *et al.*, 2015) and can be as small as 96° (Bachmann & McHale,
71 2009).

72 Most of the literature documenting water repellency in soil revolves around the
73 presence of organic matter (e.g. Ellies *et al.*, 2005) originating from vegetation and
74 microbial activity, which alters the chemistry of soil particles, whereas little attention
75 has been given to decipher the role of distinct topographies of soil particles on their
76 wettability. With soil particles, their topographies manifest themselves primarily
77 through their size and shape. These factors (individually or combined), together with
78 the effect of voids in between particles affect the wetting properties of granular and
79 porous media.

80 Particle size has been investigated in many disciplines such as soil science and
81 geotechnical engineering. McGhie & Posner (1980) found that coarser fractions of a
82 sand had a larger CA than the finer fractions (a difference of up to 45°), whereas
83 Bachmann *et al.* (2000) in their investigation of wettability of a soil from a natural
84 soil profile observed a decrease in CAs as particles became coarser. On the other hand,
85 Dang-Vu & Hupka (2005) investigated the effect of particle size of glass beads on
86 wettability and showed no dependence of CAs on their size. The effect of particle
87 shape on wettability has been examined mainly in the field of mineral engineering
88 where the process of froth flotation is used to separate minerals, not only with respect
89 to their particle size but also according to their shape. Lourenço *et al.* (2015) showed
90 in their investigation of the wettability of a range of minerals, which included silicates,

91 carbonates, oxides and sulphides, that spherical particles and those with larger aspect
92 ratios had larger CAs. In contrast, Ulusoy *et al.* (2003) investigated the wettability of
93 talc particles and found that elongated and flat particles were more water repellent
94 than rounded particles. The link between surface roughness and wettability of solids
95 has been explored in several studies to enhance mineral recovery and to identify the
96 range of surface roughness below which this property does not affect CAs (Chau *et al.*,
97 2009).

98 The specific objectives of this study were thus to: (i) analyse the CAs of granular
99 solids compared to a flat solid, and to relate particle size, shape and the influence of
100 voids to differences in CAs, (ii) investigate the effect of surface roughness on water
101 repellency and (iii) compare the experimental results with the outcomes of theoretical
102 models such as the Wenzel model, Cassie–Baxter model and the model proposed by
103 Bachmann & McHale (2009).

104

105 **Background**

106

107 *Contact angle and wetting models*

108 Young's model (Young, 1805) lies at the core of studies that make use of CA as a
109 measure of wettability. The model describes the relation between the interfacial forces
110 along the horizontal contact line leading to Equation (1), where θ_y is the Young's CA
111 and γ_{sl} , γ_{sg} and γ_{lg} correspond to the interfacial forces between the solid–liquid,
112 solid–gas and liquid–gas phases. The model is valid only for ideal surfaces and
113 predicts a unique value of CA for a given three-phase system in thermodynamic
114 equilibrium without considering any CA hysteresis. On real surfaces in nature such as
115 soil, CA hysteresis is caused by chemical heterogeneity and roughness effects.
116 Consequently, Young's model is, *sensu stricto*, not applicable because the CA value
117 obtained experimentally is not equal to θ_y .

118

119
$$\cos \theta_y = \frac{\gamma_{sg} - \gamma_{sl}}{\gamma_{lg}}. \quad (1)$$

120

121 Wenzel (1936) investigated the effect of roughness on surfaces and proposed to
122 modify Young's model by a factor r_f as shown in Equation (2). This factor (r_f) is a
123 geometric parameter that is independent of properties of the material; it is defined as
124 the ratio of the actual area of the surface to its projected area. In this model, the drop

125 of liquid adheres fully to the grooves of the solid as illustrated in Figure 1(a), and
126 intensification of the Wenzel CA (θ_w) can only be achieved if θ_y is larger than 90°
127 because $r_f \geq 1$.

128

$$129 \quad \cos \theta_w = r_f \cos \theta_y. \quad (2)$$

130

131 To account for the inhomogeneity of solids, the Cassie model is used to predict
132 CA, provided that θ_y of the different chemical components constituting the solids are
133 known. The CA recorded is a weighted average of surface area of the various
134 chemical components and their respective CA in contact with the drop of liquid.

135 For porous solids, such as water repellent soil, the drop of liquid can arch over
136 different particles causing air to be trapped between the solid and liquid phase as
137 shown in Figure 1(b). Air is the most water repellent material with a CA of 180° . The
138 Cassie model reduces to the Cassie–Baxter model (Equation (3)) assuming that the
139 solid particles have the same chemistry (Cassie & Baxter, 1944). The Cassie–Baxter
140 model relates the CA (θ_{cb}) to the θ_y of the solid. This model applied to surfaces such
141 as soil necessitates the absence of gravitational effects on the drop shape. To satisfy
142 this condition, the gaps in between the soil particles should be less than the capillary
143 length. The capillary length denotes a characteristic length of 2.7 mm for water at

144 standard temperature and pressure above which gravity markedly distorts the drop
145 shape. The coefficients f_1 and f_2 (equivalent to $1-f_1$) refer to the fractions of the drop
146 of liquid in contact with the solid and the air respectively. As f_2 , the fraction of liquid
147 in contact with the air increases, f_1 decreases, causing the magnitude of θ_{cb} to
148 approach the CA of air. For the extreme case where $f_1 = 1$, $\theta_{cb} = \theta_y$. There are
149 considerable challenges to the quantification of f_1 and f_2 because the exact way in
150 which the sessile drop advances and spreads from one particle to another is difficult to
151 determine (Chau *et al.*, 2009).

152

$$153 \quad \cos \theta_{cb} = f_1 \cos \theta_y - f_2. \quad (3)$$

154

155 The combination of the Wenzel and Cassie–Baxter models has also been applied
156 in several studies to predict CAs of a surface. Recently, Bachmann & McHale (2009)
157 derived a model based on uniform spheres where it was assumed that a drop of liquid
158 on one of the spheres completely wets a fraction of the curved surface area (Wenzel
159 model) and curls over to the adjacent sphere with air enclosed beneath (Cassie–Baxter
160 model). The model relates the predicted CA (θ_m) to θ_y by considering the packing of
161 the particles with a spacing parameter, ε , according to Equation (4). A relatively close
162 compact arrangement of particles attached to a plane and rigid surface will have a

163 value of ε close to 0 (i.e. a smaller liquid–air interface) than for particles with looser
164 configurations where the liquid–air interface is considerably larger. When $\varepsilon > 0.3$,
165 Equation (4) predicts an increase in θ_m for all values of θ_y as follows:

166

167
$$\cos \theta_m = \frac{\pi(1 + \cos \theta_y)}{\sqrt{3}(1 + \varepsilon)^2} \cos \theta_y - \left(1 - \frac{\pi \sin^2 \theta_y}{2\sqrt{3}(1 + \varepsilon)^2} \right). \quad (4)$$

168

169

170 *Particle characterization*

171

172 Particle size is a fundamental property used for the characterization of granular solids.

173 A sieve analysis is often used to determine the particle-size distribution. In

174 comparison, the measurement of particle shape requires more subtle approaches,

175 consequently descriptions are often qualitative such as angular, sub-angular and flaky.

176 Semi-quantitative methods to characterize particle shape include visual comparison of

177 particles against charts developed in the mid-twentieth century by Powers (1953). The

178 chart uses two of the most commonly used shape parameters in the literature:

179 sphericity and roundness. The former describes the extent to which a particle deviates

180 from a perfect sphere and the latter is a measure of how sharp the edges of a particle

181 are. In the chart, two different classes of sphericity and six categories of roundness are

182 defined; values between 0.12 and 1.00 are associated with roundness. Such techniques

183 are time consuming given the large number of particles that need to be assessed

184 individually and can also be very user-dependent. To overcome subjectivity and

185 improve accuracy in particle shape measurements, 2-D images may be obtained from

186 conventional optical microscopy techniques and analysed by image-processing in

187 packages such as ImageJ, an open source software of the National Institutes of Health

188 (<http://rsb.info.nih.gov/ij/>). Although these static image analyses offer improvements

189 over visual comparisons, characterization of a large number of particles in a wide
190 range of orientations remains time consuming. In addition to the complexity of
191 characterizing particle shape, is the large number of shape parameters described in the
192 literature. These parameters often have the same connotation, but different terms are
193 used to describe them such as convexity ratio (defined as the ratio of the whole area of
194 a particle to its convex area), which was used in Mora & Kwan (2000) and has the
195 same definition as solidity used in ISO 9276-6 (International Organization for
196 Standardization ISO, 2008). In this study, the shape parameters were investigated by a
197 dynamic image analyser; they are sphericity, aspect ratio, convexity and roundness.

198 With granular solids, two scales of surface roughness are recognized, namely the
199 surface roughness of single particles as investigated by Otsubo *et al.* (2015) on glass
200 beads and surface roughness taking into consideration a series of particles as used in
201 Ulusoy *et al.* (2003). In both cases, surface roughness measures the fluctuations in
202 heights with respect to a reference line or plane. Measurements of surface roughness
203 are typically reported by statistical parameters such as centre-line averages (R_a), root
204 mean square (R_q), maximum peak height (R_p) and maximum valley depth (R_v).
205 Techniques to quantify the surface roughness include mechanical methods such as
206 atomic force microscopy and optical methods such as confocal instruments.
207 Compared to the mechanical methods, optical ones can differentiate better between

208 the edges and outlines of features that tend to give a larger magnitude of surface
209 roughness for the same area considered (Whitehouse, 2002).

210 Techniques for measuring the void fraction include 2-D quantification
211 techniques such as optical light microscopy and scanning electron microscopy.
212 Images obtained from such instruments can then be analysed by image processing
213 packages such as ImageJ following noise removal. Three-dimensional techniques to
214 quantify voids include non-destructive techniques such as X-ray computed
215 tomography and confocal instruments. These devices enable a 3-D reconstruction for
216 a given area from which a void fraction based on volume may be determined.

217 In this study, the measurement of surface roughness by R_a values and
218 quantification of the void fraction was possible by making use of a confocal laser
219 scanning microscope. Figure 2(a) depicts a schematic representation of the particle
220 characteristics defined including particle size, particle shape and surface roughness.

221

222 **Materials and methods**

223 *Materials*

224

225 Silica sand from a quarry in Leighton Buzzard, Bedfordshire, UK., and referred to as
226 Leighton Buzzard Sand (LBS), was used. A fundamental part of this study involved
227 investigation of the effect of particle shape on wettability. The use of artificial
228 particles facilitates this task because they provide more control over their
229 characterization. Glass beads (GB) and crushed Glass (CG) particles have distinctive
230 shapes and were used to compare the effect of shape on wettability. The GB were
231 characterized by the manufacturer according to their size (mm) as follows: <0.1, 0.2–
232 0.4, 0.4–0.6, 0.6–0.8, 0.8–1.0, 1.0–1.5 and 1.5–2.0. To ascertain that the artificial
233 particles had the same initial chemistry, the three largest sizes of GB were crushed
234 with a 2.5-kg hammer in a steel mould, therefore, particles of considerably different
235 shapes and sizes resulted. The flat solid sample base, a reference for an ideal surface,
236 was a microscope glass slide (dimensions 76 mm by 26 mm and thickness of 1 mm),
237 made of soda lime-silica glass because its chemistry was comparable to the granular
238 solids (consisting predominantly of silica).

239

240

241 *Coating of solids*

242

243 The three granular solids were sieved under dry conditions with samples of materials
244 of at least 100 g; the following different particle sizes (μm) were isolated: 63–212,
245 212–300, 300–425 and 425–600. The granular solids were air-dried and all solids
246 considered were originally wettable (CA close to 0°). To alter the wettability of the
247 solids, dimethyldichlorosilane (DMDCS), an organo silicon compound (molecular
248 weight of 129.1 g mol^{-1} and density 1.06 g cm^{-3}) obtained from Acros Organics,
249 Morris Plains, NJ, USA in its liquid form was used. The reaction between the granular
250 solids and DMDCS produces a structure consisting of strong covalent siloxane bonds
251 between the silane functional group and the original surface. The reaction differs
252 considerably from coatings induced by organic substances which are more likely to be
253 abraded and thus less stable over time.

254 To minimize the effect of chemistry, i.e. insufficient coating, synthesis of the
255 granular solids was done beyond the critical concentration, which is considered the
256 smallest concentration of DMDCS required to achieve the maximum CA. Beyond this
257 concentration, no further increase in CA is possible solely by the addition of DMDCS.
258 Chan & Lourenço (2016) identified this concentration to be 0.00175% for an air-dried
259 clean sand treated with DMDCS. In this study, we added 140 μl of DMDCS from a

260 single channel pipette (Pipetman P100 from Gilson[®], Villiers-le-Bel, France) to a
261 40-g sample of granular solids (10 g of each of the aforementioned particle sizes).
262 This corresponded to a concentration of 0.37%. The microscope glass slide was made
263 water repellent by dispensing a total volume of 20 μ l of DMDCS on to its surface.
264

265 *Wettability measurements*

266

267 The wettability was assessed by the static contact angle, also referred to as apparent
268 contact angle because of effects of chemical heterogeneity and roughness (Drelich,
269 2013). The measurements of apparent contact angle (here referred to as contact angle
270 (CA)) were made by the sessile drop method using a goniometer (Drop Shape
271 Analyser 25 from KRÜSS GmbH, Hamburg, Germany). A 10- μ l drop of deionized
272 water was dispensed by an automated syringe. The sample preparation technique
273 involved sprinkling a mass of granular solids over a microscope slide to which
274 double-sided adhesive tape had been initially attached. A 1 N-weight was then applied
275 on the resulting lump of material such that one layer of the granular solid was
276 obtained. The application of a load also eliminates elongated particles with their
277 smallest dimension parallel to the microscope glass slide and ensures a more or less
278 uniform orientation of particles. The resulting surface can be described as a
279 quasi-plain surface with closely packed and uniformly shaped particles (Bachmann *et*
280 *al.*, 2000). The motion of the drop of liquid as it reached the sample was recorded by a
281 charged-coupled device camera. For each of the solids in the study, regardless of their
282 physical properties, the initial frame recorded after placing the drop and ending of
283 mechanical perturbances was selected for the measurement of CA. This corresponded

284 to an average time of 50 ms. The CA measurements carried out within this short time
285 lapse, limit infiltration of sessile drops to a minimum and prevent contact line
286 movement that might change the values of the CAs (Shang *et al.*, 2008). In addition,
287 the influence of gravity on drop shape can be disregarded because the gaps between
288 the granular solids are less than the capillary length. The deposition of sessile drops
289 on the particles can also lead to contact line pinning caused by the hindering effect of
290 different types of edges in contact with the sessile drop. Pinning of the contact line is
291 accompanied by an increase in CAs with no movement of the contact line and a
292 constant drop base radius. After extracting the initial frame, CAs were then evaluated
293 with the image processing technique proposed by Saulick *et al.* (2017). Ten CA
294 measurements were made on each sample 24 hours after coating the solids, and from
295 which the mean value and standard deviation were calculated. Measurements of the
296 CAs were carried out at a temperature of between 22 and 24°C and a relative
297 humidity between 60 and 70%.

298 *Dynamic image characterization of particle shape*

299

300 Particle shape was characterized with a dynamic image analyser, QicPic™ (Sympatec
301 GmbH, Clausthal-Zellerfeld, Germany). The device has been used to assess the shape
302 of granular solids such as soil, in several studies e.g. Altuhafi *et al.* (2012). The
303 QicPic™ enables samples of granular solids to be analysed by first dispersing them
304 by gravity: a 3-g mass was mounted on a feeding unit (VIBRI™) set at a specified
305 feed rate (15%) and allowed to fall through a gravity dispenser (GRADIS™) of height
306 50 cm and a gap width of 5 mm. A frame rate of 250 Hz was used enabling 40 000
307 2-D binary images at most to be evaluated within a couple of minutes. The resolution
308 of the lenses in the camera was 10 µm. The QicPic™ generates cumulative
309 distributions of shape parameters with the images, from which a median value (50%)
310 is then obtained. It was assumed that the number of images analysed was large
311 enough so that analysis of more images would not affect the median value. The shape
312 parameters obtained are sphericity, aspect ratio and convexity. With reference to
313 Figure 2(b–d), the following conclusions about the shape parameters may be drawn: (i)
314 sphericity, given as the ratio of P_{eqpc} to P_{real} , is a 2-D parameter with values that vary
315 between 0 and 1, particles close to unity resemble a circle more, (ii) aspect ratio of a
316 particle has a lower and upper limit of 0 and 1, respectively, and (iii) amongst the

317 three shape parameters evaluated by QicPic™, convexity is the only one that is
318 calculated as a direct function of area. Convexity is measured by comparing A_{real} to
319 A_{cvx} . The plots of cumulative distributions of the shape parameters (sphericity, aspect
320 ratio and convexity) for the 63–212 μm size of LBS are illustrated in Figure 3(a).
321 Evaluation of the median values was considered representative of the sample. For
322 example, a sphericity of 0.8819 (obtained with the LBS, with particle size 63–212 μm)
323 is the sphericity of the particle below which 50% of the particles lie.

324 In addition to the three above shape parameters, roundness was calculated using
325 the 2-D images generated. Roundness was defined as the ratio of average diameter of
326 the inscribed circles at the edges of the particle to the maximum inscribed circle
327 within the 2-D image as shown in Figure 2(e), which is similar to Wadell's (1932)
328 definition. A sample of 20 images for LBS and CG was investigated. The corners of
329 each of these particles were identified and circles were inserted manually to obtain
330 roundness. The roundness values of GB were taken as 1.0.

331 *Surface roughness and void fraction*

332

333 The granular particles were fixed on a microscope slide before being examined by a
334 3-D confocal laser scanning microscope (CLSM: Keyence corporation, Osaka, Japan).

335 The device consists of a microscope (VK-9710) connected to a violet laser (VK-9700)
336 colour scanning unit. A lens with 20× magnification was used to obtain the images.

337 Laser light intensity was set to a value of 1459 (adjustable on a scale of 0 (low light
338 intensity) to 16 383 (high light intensity)). The image size obtained in the x–y plane

339 was 0.716 mm by 0.537 mm (1024 × 768 pixels) corresponding to an area of
340 approximately 0.385 mm² (region of interest). The resolution in the x and y direction

341 was 0.7 μm. The resolution in the z-direction was 0.5 μm, i.e. optical sections were
342 acquired in 0.5 μm increments in the vertical direction. After a series of sections was

343 gathered, a 3-D reconstruction of the surface was generated to obtain the surface
344 profile. Ten measurements were made systematically on each sieved fraction of the

345 three granular solids. Figure 4 illustrates output images from the confocal laser
346 scanning microscope with particles of LBS. The proprietary software within the

347 CLSM, VK Viewer (version 2.4.0.1) enabled calculation of the two parameters of
348 interest:

349

350 1. Surface roughness

351 Surface roughness is evaluated by first identifying a baseline (or reference surface) by
352 the least squares method on the height data within the region of interest. The distance
353 from the baseline to the height of each point is then evaluated by surface roughness
354 parameters defined by ISO 4287 (International Organization for Standardization
355 ISO,1997). The surface roughness parameter used with the CLSM was the mean
356 height, R_a . The parameter is defined as the arithmetic mean of the absolute values of
357 the height of each point within the region of interest. It is expressed mathematically
358 according to Equation (5), where R_n is the height of each pixel with respect to the
359 baseline and N is the total number of pixels in the region of interest. In contrast to
360 particle shape, R_a is a direct reflection of the particles on the slide and quantifies the
361 slide area by considering their orientations and the spacing between particles.

362

363
$$R_a = \frac{1}{N} \sum_{n=1}^N |R_n|. \quad (5)$$

364

365 2. Void fraction

366 The integrated software within the CLSM, VK Viewer enabled the volume of
367 granular solids (v) on the slide to be calculated. The void fraction was calculated as
368 the difference between the total volume (V) and v expressed as percentage of V . To
369 obtain V , the heights of the lower and upper limits were adjusted accordingly from the
370 3-D reconstructed surface in the region of interest. The height of the lower limit was
371 set as the base of the slide. To obtain an average diameter that is representative of the
372 particles within the region of interest, Feret diameters (D) obtained from the QicPic™
373 were compared to the sieve fraction. Figure 3(b) illustrates such a comparison for the
374 LBS with the particle size of 63–212 μm . The median values (D_{50}) for the minimum,
375 mean and maximum Feret diameters defined by QicPic™ were 186.3, 219.4 and
376 265.6 μm , respectively. These values indicate that the median value obtained from the
377 Feret minimum diameter was closest to the particle size obtained by the sieve analysis.
378 Similar observations have also been made with the other granular solids, regardless of
379 particle size. Therefore, to calculate V , the upper height limit (i.e. the average particle
380 diameter in the region of interest) was set to the median value of the Feret minimum
381 diameter.

382 .

383 **Results and discussion**

384

385 To investigate the effect of wettability on particle size, the particle shape was isolated
386 by considering the different sizes of the same granular solid. Similarly, to examine the
387 effect of particle shape on wettability, the particle size was isolated and a given
388 particle size of the three granular solids was investigated.

389

390 *Wettability and particle size*

391

392 The mean CAs of the granular solids all exceeded the CA of the flat hydrophobized
393 microscope slide, which was 103° (SD = 2°). The standard deviations of the CAs
394 were 5° at most and accord with previous studies such as by Bachmann *et al.* (2000).

395 The effect of particle size on the CAs is illustrated in Figure 5(a). The largest CA
396 recorded (125° (SD = 4°)) was obtained with the 63–212 μm fraction of CG and the
397 smallest one (107° (SD = 5°)) was obtained with the 425–600 μm fraction of LBS.

398 With all three granular solids, there was a general decrease in CAs as the particle size
399 increased. For example, with particle sizes of 63–212, 212–300, 300–425 and 425–
400 600 μm for CG the CAs obtained are: 125° (SD = 4°), 119° (SD = 5°), 114° (SD = 4°)
401 and 110° (SD = 5°). The difference in CAs between the extreme particle sizes (63–
402 212 and 425–600 μm) for CG, LBS and GB are 14° , 16° and 9° , respectively.

403 Reasons for the larger CAs of the finer particles include a relatively larger void
404 fraction with such particles. The plot of void fraction against particle size is shown in
405 Figure 5(b). The smallest void fraction (37.0%) was obtained with the particle size of
406 425–600 μm of LBS, whereas the void fraction recorded with the particle sizes of
407 300–425, 212–300 and 63–212 μm of LBS were 52.0, 62.1 and 68.8%, respectively.
408 Similar trends were also observed with the particles of GB and CG.
409

410 *Wettability and particle shape*

411

412 The results of the effect of particle shape for the different particle sizes of the three
413 granular solids are shown in Figure 6. All the shape parameters considered are
414 sensitive to the different granular solids. Of these parameters, convexity had the
415 smallest range (from 0.8944 with the 63–212 μm particle size of CG to 0.9777 with
416 the 425–600 μm particle size of GB), whereas the roundness showed the largest range
417 (from 0.3046 with the 425–600 μm particle size of CG to unity with the particles of
418 GB). The largest values of sphericity (0.9429), aspect ratio (0.9633) and convexity
419 (0.9777) were all for GB with a particle size of 425–600 μm and the smallest values
420 of sphericity (0.8322), aspect ratio (0.6656), convexity (0.8944) and roundness
421 (0.3046) were obtained with CG. With the exception of convexity, the variation in
422 particle shape across the range of particle sizes shows little to no change compared to
423 the variation in relation to the different materials where the difference was noticeably
424 more pronounced. The largest standard deviation of the shape parameters was for the
425 aspect ratio of CG (0.13). This is partly attributed to the crushing process;
426 nevertheless most standard deviations of the shape parameters of the three granular
427 solids were comparable, for example those of sphericity for CG, LBS and GB are 0.06,
428 0.06 and 0.04 respectively for the particle size of 425–600 μm . Comparison of
429 particle shape before and after reaction with DMDCS showed no change, indicating

430 that the thickness of the coating induced by DMDCS at a concentration of 0.37% was
431 smaller than the resolution of the lenses in QicPic™.

432 Although the general trend of the shape parameters is for one of them to increase
433 as the other increases, the extent to which they are positively correlated varies.
434 Analysis of the shape parameters showed a relatively poor linear correlation
435 (determined by the correlation coefficient r) between convexity and the remaining
436 three shape parameters. The values of the correlations between convexity and
437 roundness, sphericity and aspect ratio were 0.612, 0.743 and 0.748, respectively.
438 Results reported by Altuhafi *et al.* (2012) showed an even weaker correlation between
439 convexity and aspect ratio ($r = 0.469$). By comparison, the linear correlations between
440 sphericity, aspect ratio and roundness were much stronger ($r > 0.96$). A strong
441 correlation was noted between roundness and aspect ratio ($r = 0.969$) despite the
442 difference in sample size.

443 Figure 7 shows the plot of contact angle against the four shape parameters. There
444 is a reduction in CA with an increase in the value of the shape parameters. For
445 example, for the particle-size range of 63–212 μm , a change in CA from 125° (SD =
446 4°) to 117° (SD = 4°) corresponded to an increase from 0.8369 to 0.9374, 0.6656 to
447 0.9263, 0.8944 to 0.9426 and 0.4747 to 1 in sphericity, aspect ratio, convexity and
448 roundness, respectively. The relatively steeper slopes for the finest fractions show that

449 as particle size decreases, the effect of particle shape becomes more predominant in
450 dictating wettability. Qualitatively, the angular particles were the most water repellent.
451 These results can be attributed to pinning effects caused by the edges of particles in
452 contact with the sessile drop. For sharp edges like those present with CG, the sessile
453 drop adheres and becomes pinned to the edge rather than moving over it. The process
454 of pinning with the sharp edges causes CAs to increase, as observed by Oliver *et al.*
455 (1977). As a result, a larger force is required for the sessile drop to move to the next
456 particle. This force can be in the form of a capillary force, which prevents the sessile
457 drop from entering the grooves of the granular solids and increases the fraction of air
458 between the solid and liquid interface (Cao *et al.*, 2007). In contrast, with the
459 relatively smoother edges of GB, the pinning effect was less significant resulting in a
460 smaller CA.

461

462

463 *Wettability and surface roughness*

464

465 The relation between surface roughness, characterized by R_a on the CAs is illustrated
466 in Figure 8. For all three granular solids, the particle-size range of 425–600 μm has
467 the largest R_a value and it is smallest for the range 63–212 μm . Figure 8 also shows
468 the increase in standard deviation of R_a as particle size increases: with CG, the
469 standard deviation for the particle sizes 63–212, 212–300, 300–425 and 425–600 μm
470 are 5.36, 8.65, 10.97 and 26.65 μm , respectively. Comparison of R_a values for the
471 different materials shows that particles with a small shape parameter as for CG
472 (sphericity = 0.8369, aspect ratio = 0.6656, convexity = 0.8944 and roundness =
473 0.4647 for particle size 63–212 μm) generally, results in a smaller R_a . For the finest
474 particle size, the R_a values of CG, LBS and GB are as follows: 34.1 μm (SD = 5.36
475 μm), 44.8 μm (SD = 2.07 μm) and 58.9 μm (SD = 3.42 μm). These data show that R_a
476 is not only sensitive to changes in particle size but also to particle shape. A fine
477 particle size with a small shape parameter corresponds to a small value of R_a . An
478 increase in surface roughness from 34.1 to 95.4 μm corresponds to a decrease in CA
479 from 125 to 107°. There seems to be little to no change in CAs of the granular solids
480 considered beyond an R_a value of 95.4 μm .

481 *Comparison with wetting models*

482

483 To investigate differences between measured CAs and theoretical models for the flat
484 and granular solids, the measurements of CAs, surface roughness and void fractions
485 were compared with predicted values.

486 The values of the roughness factor, r_f , Equation (2) for θ_y of 103° is given in
487 Table 1. The roughness factor increases as particle size decreases; the smallest value
488 was 1.304 obtained with the 425–600 μm particle size of LBS and the largest value
489 was 2.531 obtained with the 63–212 μm particle size of CG. Although the values of
490 R_a measured in this study with CLSM cannot be translated directly to a roughness
491 factor, the increase in surface roughness leading to a decrease in CAs for the water
492 repellent granular solids in this study qualitatively contradict the Wenzel equation.

493 The variation in void fraction of the granular solids with the CAs is shown in
494 Figure 9. The largest void fraction is generally associated with the finest sieve fraction.
495 There was a maximum difference in void fraction of around 40.4% between the
496 particle size of 212–300 μm for CG and that of 425–600 μm for LBS. Figure 9 shows
497 that an increase in void fraction corresponds to an increase in CA. Figure 9 also
498 shows the Cassie–Baxter model (Equation (3)) plotted with θ_y of 103° . This model,
499 which fails to take into account the physical properties of granular solids (particle size,

500 particle shape and their distributions) overestimates the data obtained experimentally.
501 This comparison was made with the assumption that for very water repellent granular
502 solids as used in this study, the bottom interface of the sessile drop bridges from one
503 particle to another in the case of maximum void fraction leading to additional liquid–
504 air interfaces adjacent to the solid–liquid interfaces. This effect increases the
505 measured CAs.

506 The experimental data obtained for GB were also compared with the model
507 proposed by Bachmann & McHale (2009). The explicit determination of the spacing
508 parameter, ε , using the 2-D images from CLSM gave an average value of 0.261 for
509 the particle sizes of 63–212 and 212–300 μm . This value is smaller than 0.319, the
510 value obtained by Bachmann & McHale (2009) with GB particles, but it is close to
511 the lower boundary of the stable range (0.260 to 0.320) reported by the same authors.
512 Because of difficulties in obtaining the exact location of the centroid of the LBS and
513 CG for all particle sizes, the analysis could not be done on these particles. Figure 10
514 shows the predicted contact angle by the model (θ_m) plotted against θ_y for ε equal
515 to 0.261. For $\theta_y > 42^\circ$, the model predicts an increase in CA. The predicted CA, θ_m ,
516 for $\theta_y = 103^\circ$ is 131° . This value exceeds the maximum CA reported for GB in this
517 study (117° (SD = 4°)) which was recorded for the 63–212 μm fraction. The model is
518 very dependent on the void fraction, determined by the spacing parameter. A probable

519 reason for the overestimation of CA is because of the relatively closer and compact
520 arrangement of the particles used in this study. Figure 10 also shows the plot of ε
521 equal to 0 and 0.060, values adjusted to fit the maximum CA measured with GB,
522 taking into consideration the range of experimental error. The θ_m value for ε equal to
523 0 and 0.060 were 117° and 121° , respectively.

524

525 **Conclusions**

526

527 All granular solids had larger CAs than the microscope slide. Finer particle sizes had
528 larger CAs than their coarser counterparts, irrespective of the material considered. By
529 isolating the different sieve fractions, we showed that angular particles represented by
530 relatively smaller values (i.e. values diverging from unity) of sphericity, aspect ratio,
531 convexity and roundness had the largest CAs. The surface roughness, R_a , had small
532 values for finer particle sizes and more angular particles.

533 Changes in CAs of granular solids have often been attributed exclusively to a
534 change in chemistry. In this study, we assumed that the coatings were similar for all
535 size classes and materials because the surfaces were treated with larger amounts of
536 DMDCS than required to achieve the maximum CA. This study has shown that
537 granular solids with a particular size and shape are able to control their wettability.
538 Rather than relying on coatings, taking the topography of granular solids into
539 consideration can lead to (i) reduction in the amount of chemicals (e.g. silanes) used
540 to achieve the desired level of water repellency, (ii) a longer durability in water
541 repellency compared to coatings that may be abraded over time and (iii) help to
542 source suitable granular solids from natural soils and waste materials for use in
543 applications such as slope covers.

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545

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550

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644 _____

645 **FIGURE CAPTIONS**

646

647 **Figure 1** Schematic illustrations of a drop of liquid in (a) Wenzel state (b) Cassie–
648 Baxter state.

649 **Figure 2** (a) Characterization of particles: particle size, particle shape and surface
650 roughness characterized by sieve analysis, dynamic image analyser and confocal laser
651 scanning microscope respectively, and (b–e) definitions of shape parameters used:
652 sphericity, aspect ratio, convexity and roundness.

653 **Figure 3** (a) Cumulative distribution of particle shape parameters of Leighton
654 Buzzard Sand and evaluation of their median values and (b) particle-size distributions
655 of Leighton Buzzard Sand.

656 **Figure 4** Topographic images of Leighton Buzzard Sand from confocal laser
657 scanning microscope (a) 2-D laser image, (b) 2-D height image and (c) 3-D display.

658 **Figure 5** (a) Contact angle plotted against particle size. The error bars indicate
659 standard error of the mean (ten measurements) and (b) variation of void fraction with
660 particle size. The lines are to guide the eye.

661 **Figure 6** Variation of (a) sphericity, (b) aspect ratio, (c) convexity and (d) roundness
662 for the different particle sizes of the granular solids. The lines are to guide the eye.

663 GB, Glass beads; LBS, Leighton Buzzard Sand; CG, Crushed glass.

664 **Figure 7** Contact angle plotted against (a) sphericity, (b) aspect ratio, (c) convexity
665 and (d) roundness. The error bars indicate standard error of the mean (ten
666 measurements). The lines are to guide the eye.

667 **Figure 8** Contact angle plotted against surface roughness, R_a . The error bars indicate
668 standard error of the mean (ten measurements).

669 **Figure 9** Variation in void fraction with contact angle

670 **Figure 10** Predicted contact angles according to model proposed by Bachmann &
671 McHale (2009) for the experimentally measured value of $\varepsilon = 0.261$ and fitted values
672 of $\varepsilon = 0$ and $\varepsilon = 0.060$

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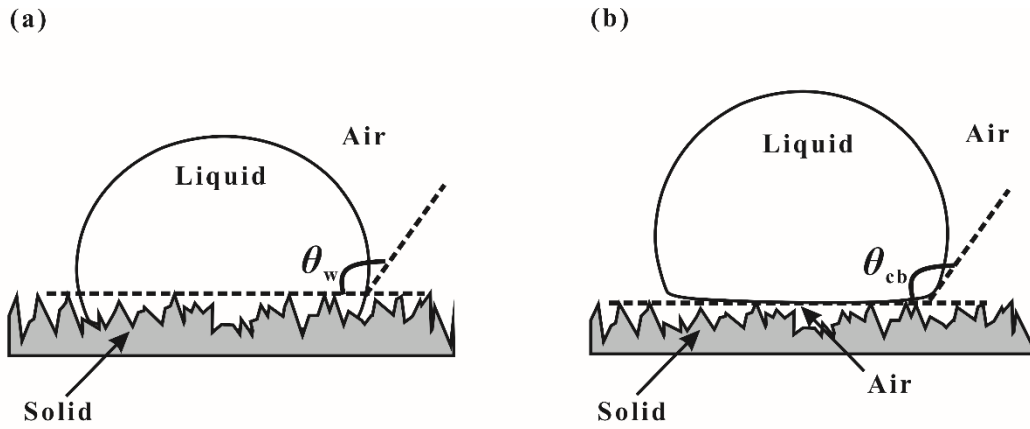
674 **TABLES**

675 **Table 1** Calculated Wenzel roughness factor, r_f from θ_y

Particle Size/ μm	Granular Solids		
	GB	LBS	CG
63–212	2.004	2.428	2.531
212–300	1.981	2.244	2.173
300–425	1.580	1.592	1.783
425–600	1.378	1.304	1.543

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677 GB, Glass beads; LBS, Leighton Buzzard Sand; CG, Crushed glass. _____



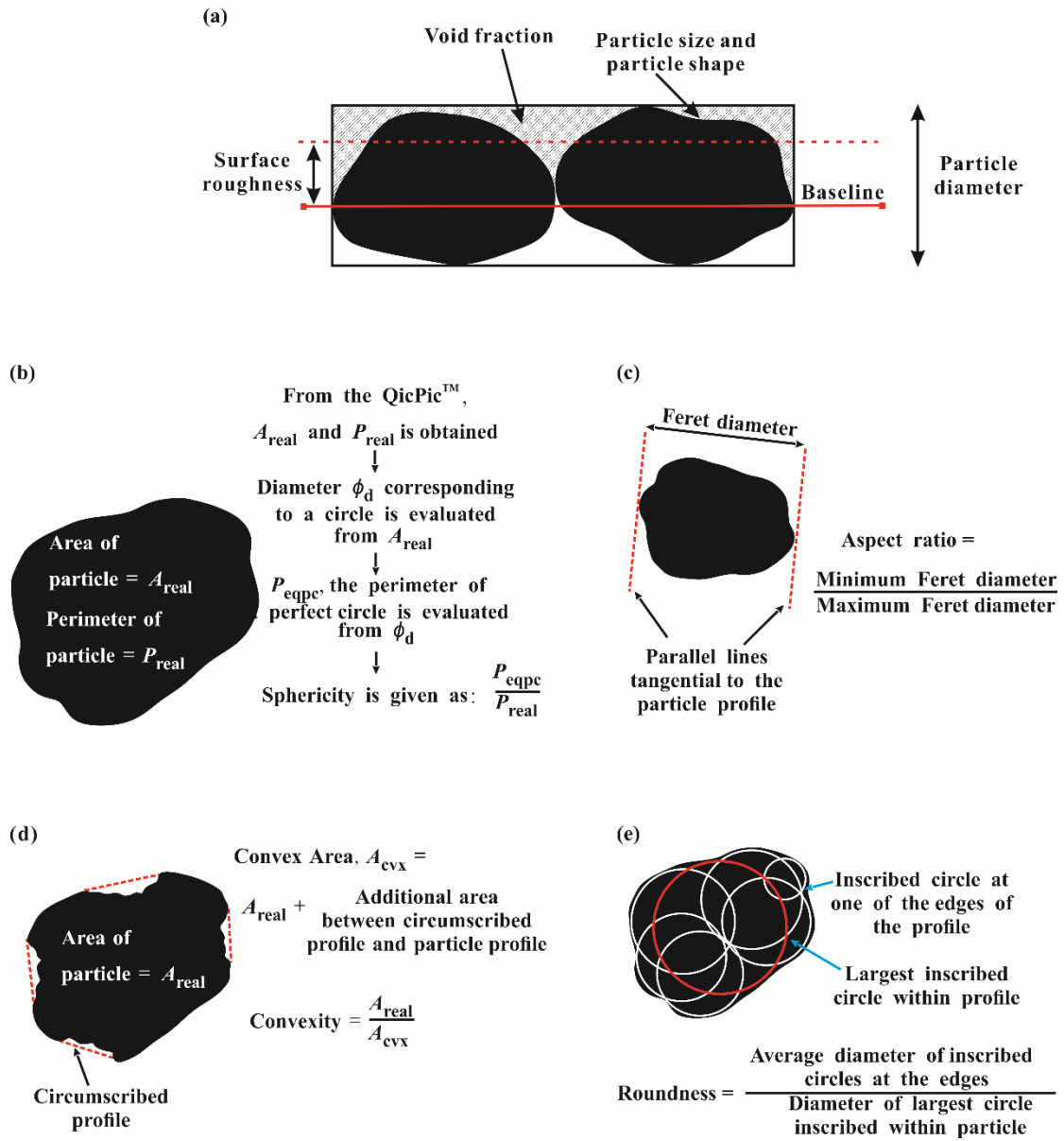
678

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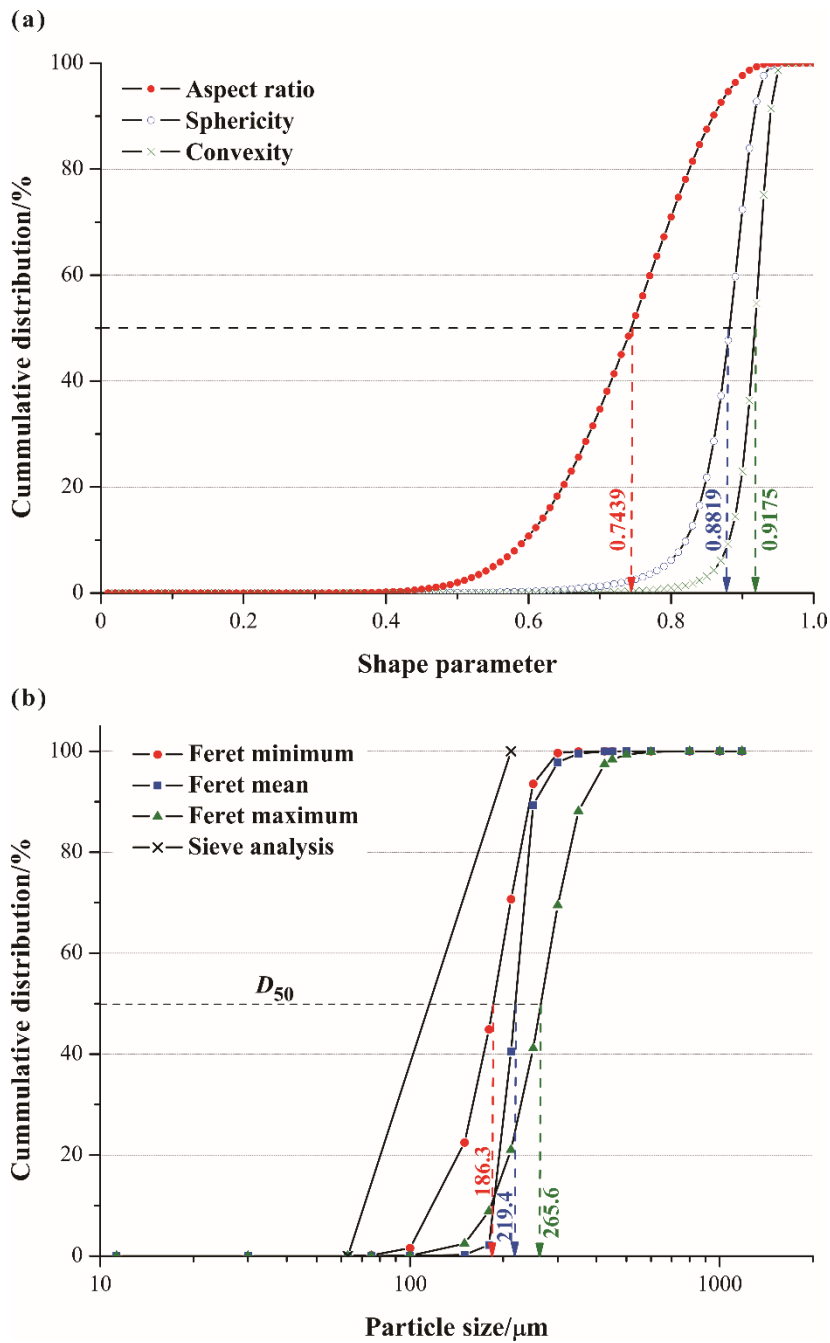
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686 scanning microscope respectively, and (b–e) definitions of shape parameters used:

687 sphericity, aspect ratio, convexity and roundness.

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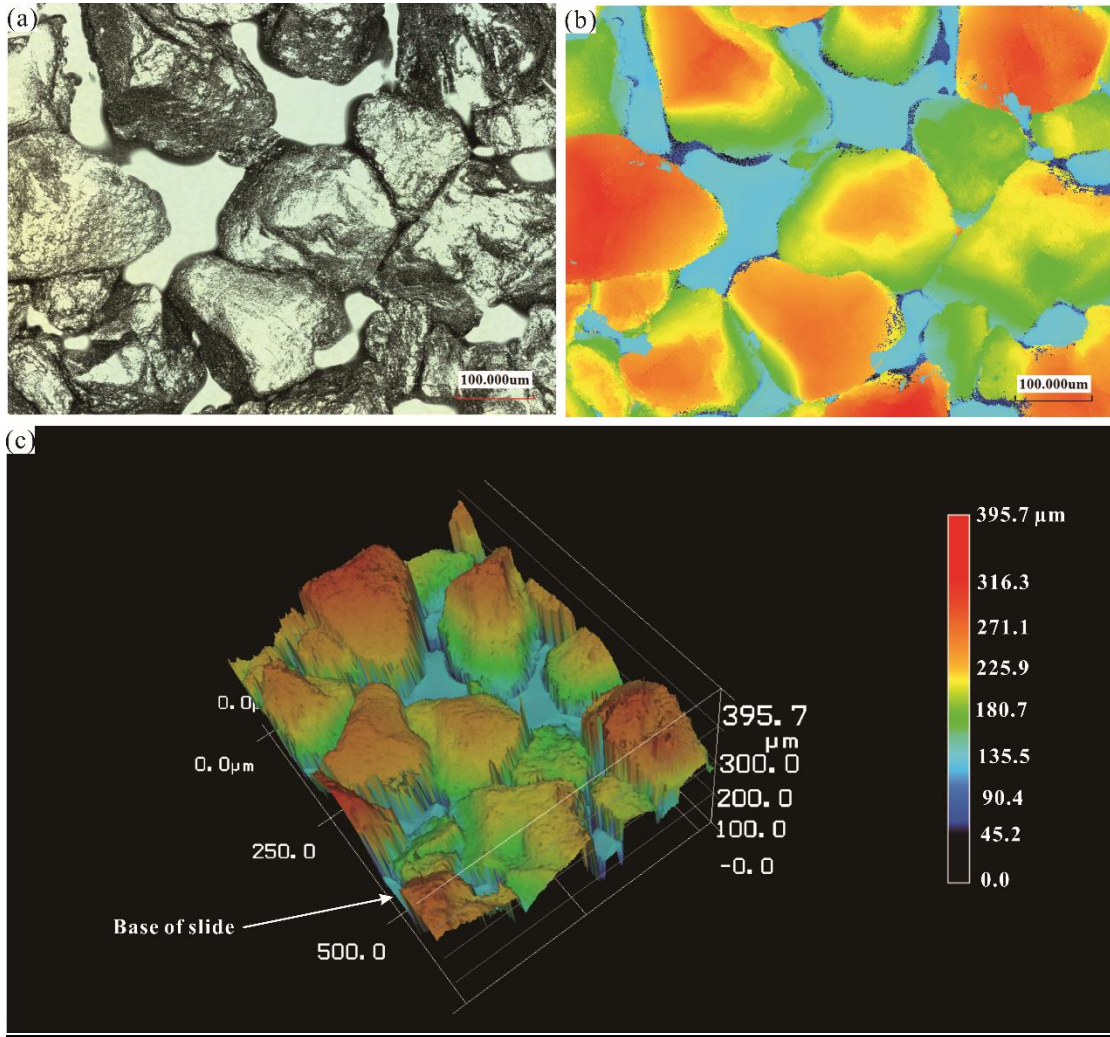
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692 Buzzard Sand and evaluation of their median values and (b) particle-size distributions

693 of Leighton Buzzard Sand.

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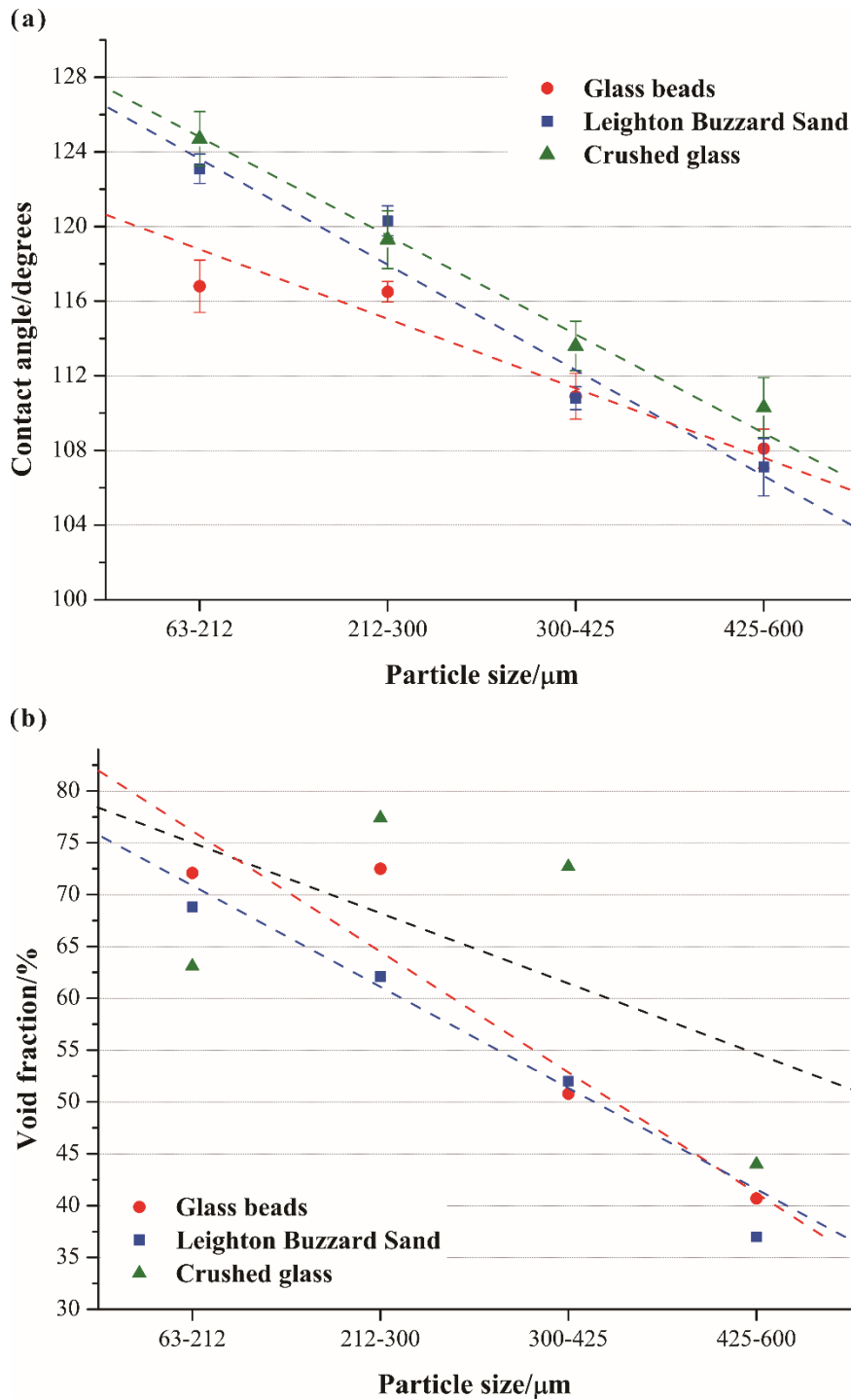
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697 **Figure 4** Topographic images of Leighton Buzzard Sand from confocal laser

698 scanning microscope (a) 2-D laser image, (b) 2-D height image and (c) 3-D display.

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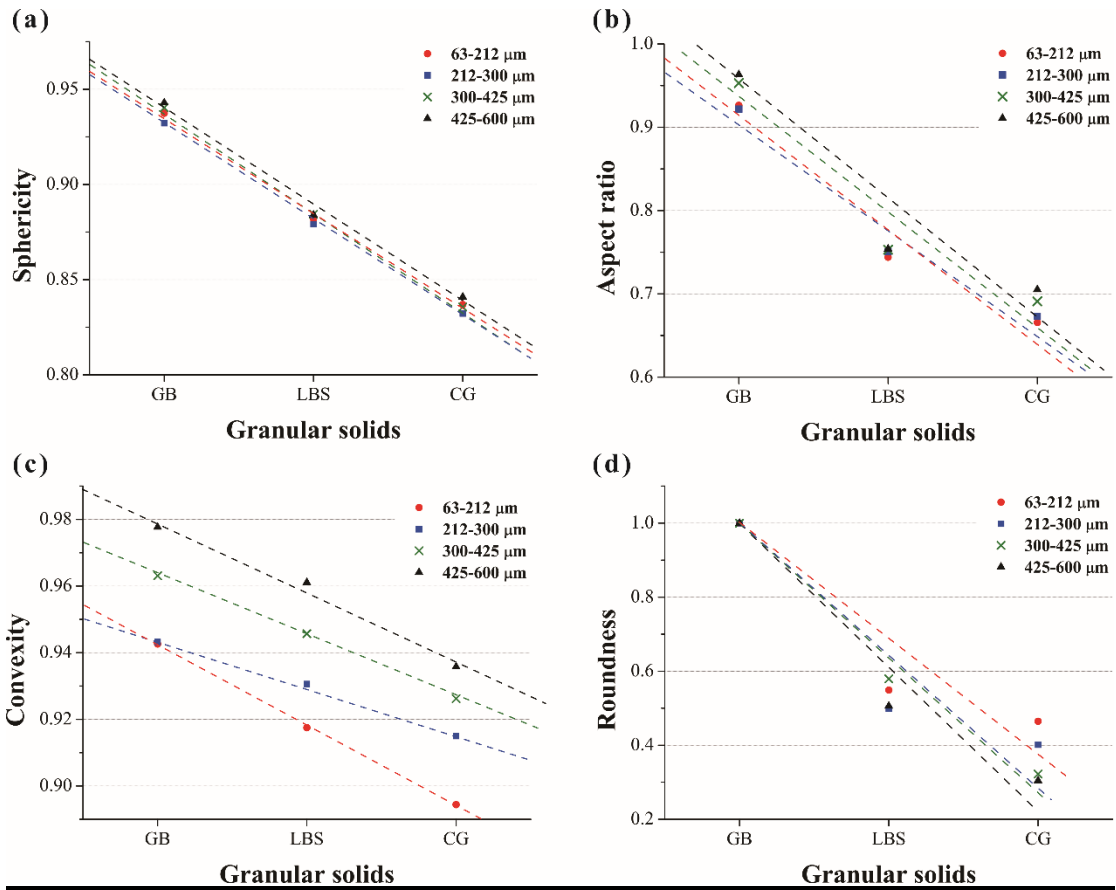
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702 **Figure 5** (a) Contact angle plotted against particle size. The error bars indicate

703 standard error of the mean (ten measurements) and (b) variation of void fraction with

704 particle size. The lines are to guide the eye.

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706

707 **Figure 6** Variation of (a) sphericity, (b) aspect ratio, (c) convexity and (d) roundness

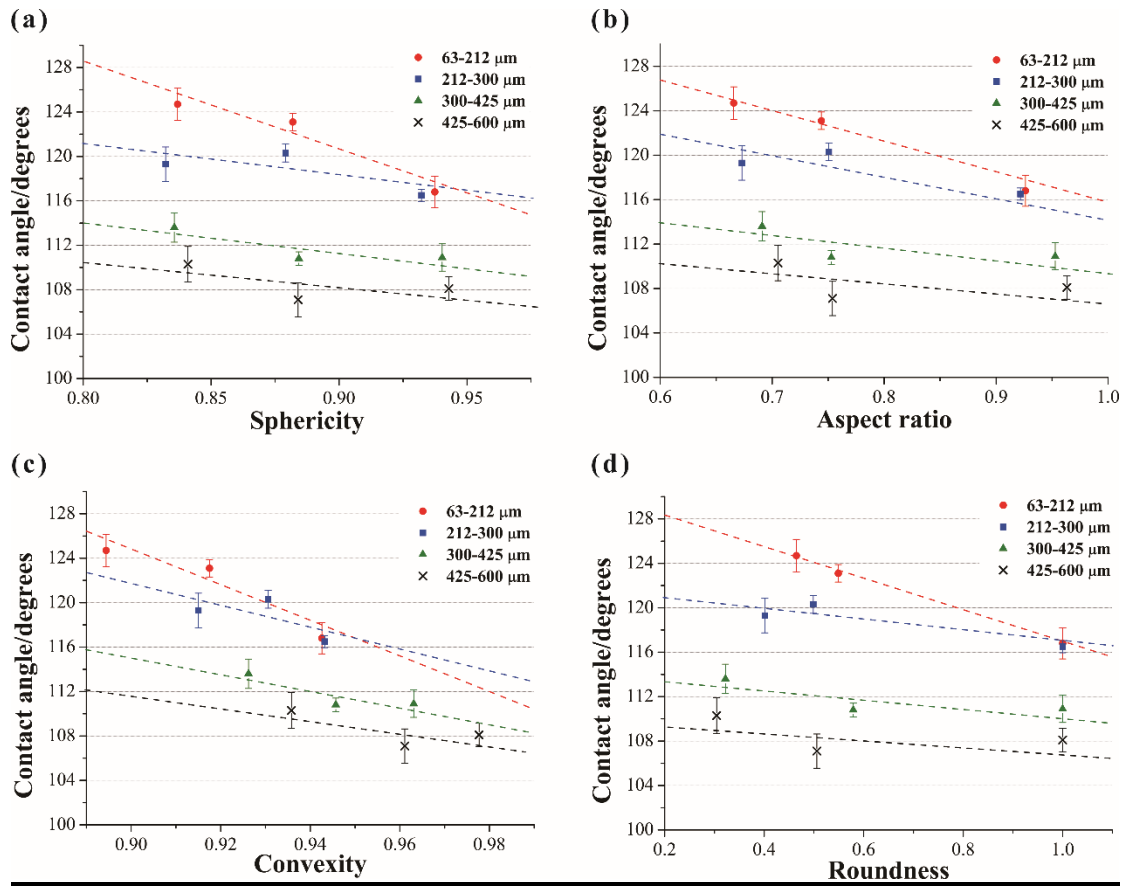
708 for the different particle sizes of the granular solids. GB, LBS and CG are

709 abbreviations for Glass beads, Leighton Buzzard Sand and Crushed glass respectively.

710 The lines are to guide the eye.

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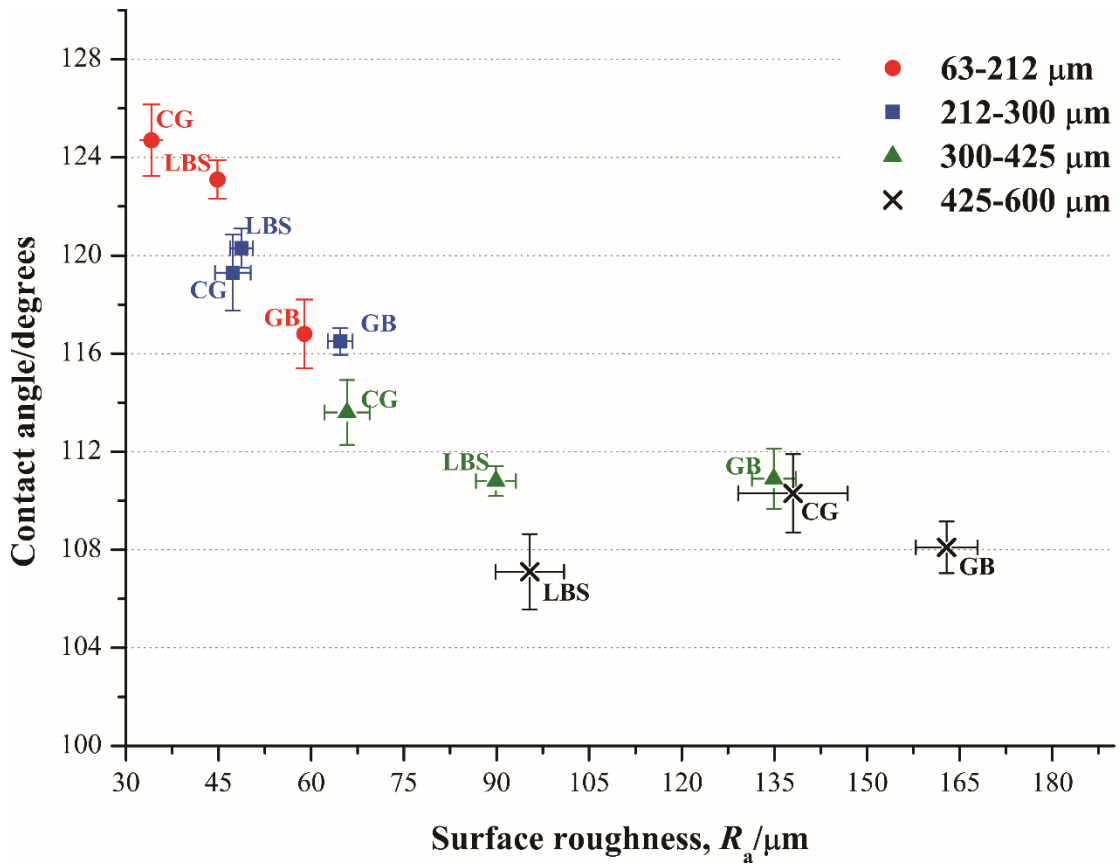
714 **Figure 7** Contact angle plotted against (a) sphericity, (b) aspect ratio, (c) convexity

715 and (d) roundness. The error bars indicate standard error of the mean (ten

716 measurements). The lines are to guide the eye.

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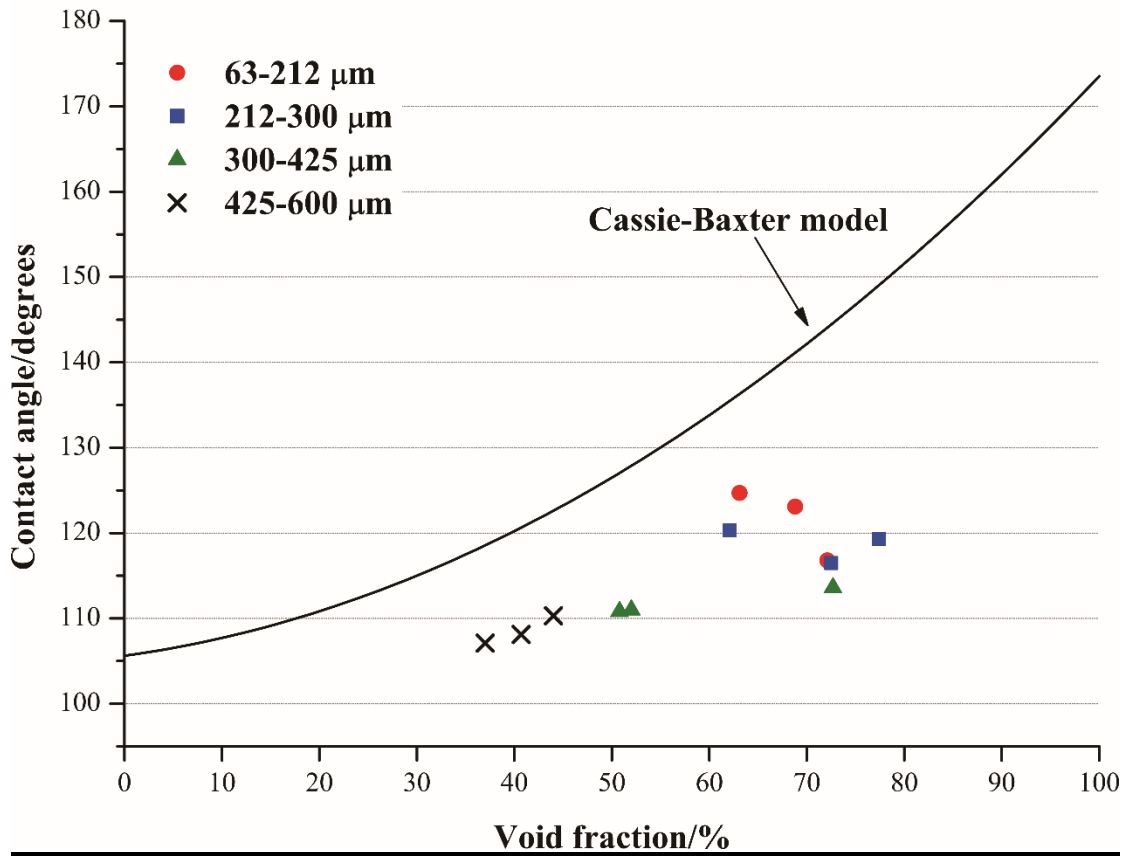
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720 **Figure 8** Contact angle plotted against surface roughness, R_a . The error bars indicate

721 standard error of the mean (ten measurements).

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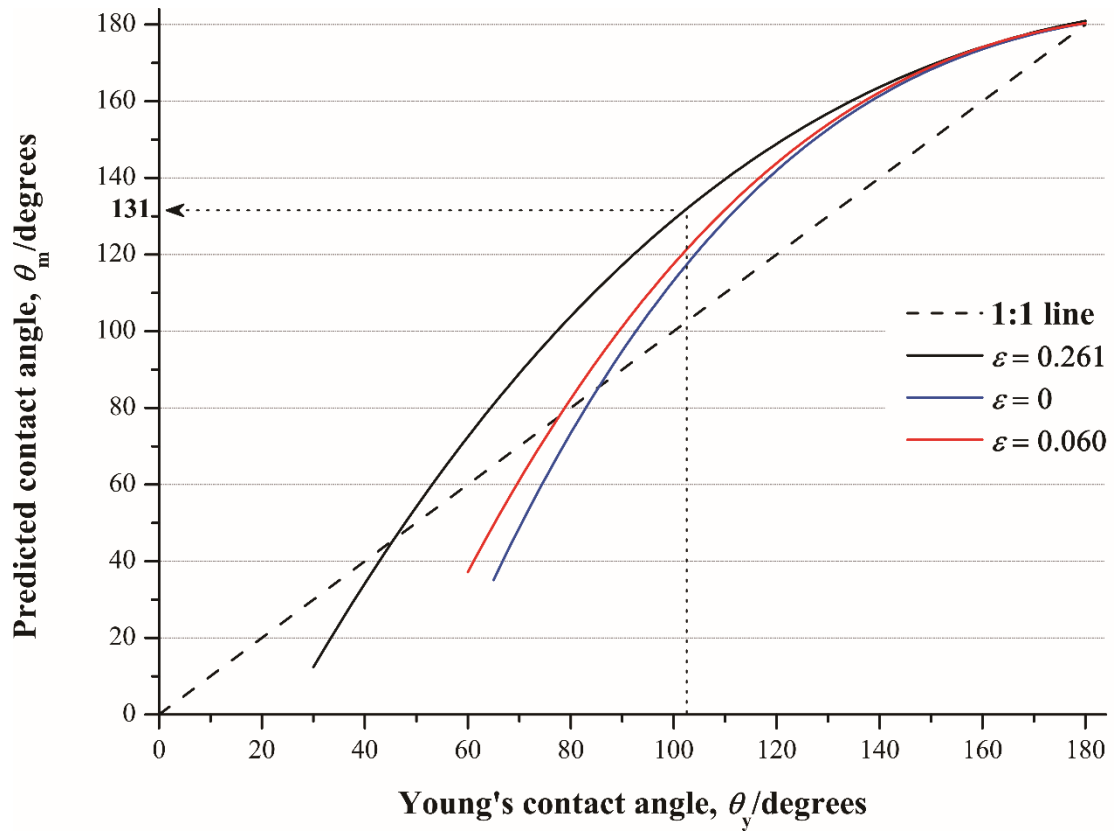


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725 **Figure 9** Variation in void fraction with contact angle

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729 **Figure 10** Predicted contact angles according to model proposed by Bachmann &
 730 McHale (2009) for the experimentally measured value of $\epsilon = 0.261$ and fitted values
 731 of $\epsilon = 0$ and $\epsilon = 0.060$

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733