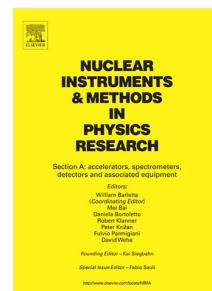


## Journal Pre-proof

Image quality determination of a novel digital detector for X-ray imaging and cone-beam computed tomography applications

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1 Image Quality Determination of a Novel Digital Detector for  
2 X-ray Imaging and Cone-Beam Computed Tomography  
3 Applications

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## 20 Abstract

21 The demand for adequate image quality with low radiation doses for patients has greatly  
22 increased. This is especially true in the case of position verification in radiotherapy which  
23 requires a high number of images per patient. This study presents a physical characterisation  
24 of a new clinical detector named “Lassena (CsI)” based on a thick layer of structured thallium  
25 activated caesium iodide and complementary metal-oxide semiconductor technology with  
26 active pixel sensor architecture for general X-ray imaging and cone-beam computed  
27 tomography (CBCT) applications. We made a critical appraisal of its performance for the  
28 first time and determined its signal transfer property (STP) and its detective quantum  
29 efficiency (DQE) by acquiring the pre-sampling modulation transfer function (pMTF) and  
30 normalised noise power spectrum (NNPS) in addition to the dark current calculation. The  
31 investigation was conducted with the application of three X-ray beam qualities: (50 kV  
32 (RQA3), 70 kV (RQA5) and 90 kV (RQA7)) in compliance with the International  
33 Electrotechnical Commission (IEC 62220-1(2003)) standard. The STP was found to be linear  
34 with the coefficient of determination ( $R^2$ ) more than 0.9995 in all cases. The spatial  
35 resolution and NNPS results led to acceptable DQE values at all energies; in particular the  
36 DQE values at 0.5 line pairs per mm (DQE(0.5)) which were 0.46 for RQA3, 0.52-0.56 for  
37 RQA5 and 0.55-0.59 for RQA7. Lastly, the dark current was 2.51 pA/cm<sup>2</sup> for a 50 µm pixel  
38 pitch. For CBCT applications, Lassena (CsI) showed very promising results.

39  
40 Keywords: signal-to-noise ratio, image quality, DQE, detector characterisation.

## 41 1. Introduction

42 Digital X-ray detectors have gained widespread use in clinical applications. This is due to  
43 their high degree of performance and accuracy which can be quantified using the detective  
44 quantum efficiency (DQE) parameter to indicate the image system performance as it varies  
45 from one system to another [1-5]. In general, the DQE describes the transmission of the  
46 signal to noise ratio (SNR) in X-ray imaging detectors taking into account the spatial  
47 resolution and noise [5,6]. An ideal imaging system would have a DQE value of one at all  
48 spatial frequencies. In all practical cases, however, the DQE decreases as a function of spatial  
49 frequency due to the increased effect of the noise as a function of spatial frequency [6,7]. The  
50 other two criteria of the image quality assessment are pre-modulation transfer function  
51 (pMTF) and normalised noise power spectrum (NNPS). The pMTF indicates the image  
52 resolution and the spatial frequency corresponding to 10% of pMTF is frequently used to  
53

54 express the limiting resolution of the system [7]. In general radiography, the adequate  
55 limiting resolution for a radiation detector ranges between 3 and 5 line pairs per millimetre  
56 (lp/mm) [8]. The NNPS is one of the popular metrics providing a quantitative description for  
57 the noise variance in an imaging system as a function of spatial frequency [9].

58 Digital radiation detectors based on complementary metal-oxide semiconductor (CMOS)  
59 technology are becoming prevalent since they offer low voltage operation, low power  
60 consumption and low cost while maintaining acceptable image quality [10,11]. In particular,  
61 CMOS sensors based on the active pixel sensor (APS) structure as they allow in-pixel  
62 buffering of the signal and therefore they have better image quality [11].

63 The properties of the scintillator and its thickness have an impact on the image quality as  
64 well. X-ray imaging detectors commonly use a powdered scintillator also known as a  
65 phosphor such as gadolinium oxysulfide doped with terbium also known as Gadox or P43 in  
66 a polymer matrix or caesium iodide doped with thallium (CsI(Tl)). Both scintillators have  
67 large conversion gains and their peak emission wavelength is in the green portion of the  
68 visible spectrum at 550 nm for CsI(Tl) and 545 nm for Gadox which matches the peak  
69 quantum efficiency of silicon based sensor resulting in a high signal collection. CsI(Tl) can  
70 be grown to have a micro-columnar structure which reduces the laterally spread of the  
71 scintillation light resulting in greater spatial resolution than phosphor screens. Increasing the  
72 thickness of the scintillator leads to higher X-ray absorption but any scintillation light  
73 generated at the top of the scintillator will spread more resulting in lower spatial resolution  
74 and higher noise [11-14].

75 The purpose of this work is a characterisation of a new radiation detector through physical  
76 figures of merit (pMTF, NNPS, and DQE) and proposing it as a radiation detector for general  
77 radiographic imaging and cone-beam computed tomography (CBCT) applications which play  
78 an essential role in image-guided radiation therapy (IGRT) and adaptive radiotherapy (ART)  
79 aspects.

## 80 **2. Materials and Methods**

81 The detector under investigation which is referred to as “Lassena (CsI)” is a three  
82 transistor (3T) wafer-scale CMOS APS designed by The Science and Technology Facilities  
83 Council’s Rutherford Appleton Laboratory (RAL, Oxford, UK). It is coupled to a 1000  $\mu\text{m}$   
84 CsI(Tl) scintillator to convert incident X-ray photons into optical light. CsI (Tl) is used in  
85 medical imaging due to its high resolution and low noise [6, 7]. A thicker layer of CsI (Tl)  
86 was used with this Lassena sensor to improve its efficiency at low doses. However, this  
87 comes at the expense of spatial resolution. The detector consists of two sensors tiled next to  
88 each other with a 1-pixel dead area. Each sensor has an active area of 12 cm  $\times$  14 cm to give  
89 a total area of 24 cm  $\times$  14 cm with an effective resolution of 2786 $\times$ 2400 pixels and the pixel  
90 pitch is 50  $\mu\text{m}$ . The sensor has a quantum efficiency of 50 % at 540 nm, the image depth of  
91 analogue-to-digital converter (ADC) is 14-bit and the noise of the sensor with new dedicated  
92 readout electronics was improved from the original reported value of 70 electrons rms (root  
93 mean square) to be 40 electrons rms, as measured by the photon transfer curve (PTC) method  
94 [15].

95 The characterisation was obtained according to guidelines of the published International  
96 Electrotechnical Commission (IEC 62220-1 (2003)) standard that contains a standardised  
97 methodology for digital detector characterisation. The measurement uncertainty for DQE,  
98 MTF and NNPS was calculated by repeating the measurements three times.

### 99 **2.1. Beam Quality**

100 The characterisation was performed at University College London (UCL) laboratory with  
101 an X-ray source (HS-MP1, Ago X-ray limited, England) of focal spot 1 mm, and a tungsten  
102 target with aluminium filtration (W/Al). The measurements were completed using three

103 different standard beam qualities: RQA3 (50 kV), RQA5 (70 kV) and RQA7 (90 kV) as  
 104 RQA3 is suitable for pediatric extremities imaging, while RQA5 is applicable for adult  
 105 extremity radiography and RQA7 is commonly employed for CBCT imaging. The test  
 106 geometry was compliant with IEC 62220-1 (2003) standard. The detector was placed at a  
 107 distance of 150 cm from the X-ray source. This was to ensure beam uniformity on the  
 108 detector surface. The half value layer (HVL) was measured to determine the beam energy  
 109 required for measurements for each beam condition using our source [16].

## 110 2.2. Signal transfer property (STP)

111 The STP describes the relationship between the detector mean pixel value (MPV) and air  
 112 kerma ( $K_a$ ) to determine how the detector responds to the input signal. Ideally, the response  
 113 should be linear without any image processing apart from non-uniformity correction and  
 114 pixel defect calibration [16]. The mean pixel value was studied, in addition to the detector's  
 115 response fit using Eq 1:

$$116 \text{MPV} = B K_a + A \quad (1)$$

117 Where A and B are offset and STP gradient of the fit parameters respectively.

## 118 2.3. Normalised Noise Power Spectrum (NNPS) Determination

119 The NNPS determines the relative noise properties in detector response [17]. For this  
 120 measurement, 30 dark and 30 bright images (across a range of tube currents) were acquired  
 121 below the saturation level. A second-order polynomial fit was applied to correct the beam  
 122 non-uniformity. The noise power spectrum (NPS) analysis was conducted according to the  
 123 IEC protocol by dividing an image into a number of squares referred to as regions of interest  
 124 (ROIs). Each of these measured  $256 \times 256$  pixels with overlapping of 128 pixels. The NPS  
 125 was acquired as a function of spatial frequency by applying the fast Fourier transform (FFT)  
 126 [18,19] using Eq 2:

$$127 \text{NPS}(u,v) = \frac{\Delta x \Delta y}{M N_x N_y} \sum_{m=1}^M |FFT\{I(x_i, y_i) - S(x_i, y_i)\}|^2 \quad (2)$$

128 Where u and v are the spatial frequencies reflecting x and y,  $\Delta x$  and  $\Delta y$  are pixel pitches in  
 129 x and y directions,  $N_x$  and  $N_y$  express the ROI size in x and y directions, M is the ROIs  
 130 number which is used in averaging and  $S(x,y)$  and  $I(x,y)$  are the fitted 2D function and  
 131 corrected flat field image respectively. The NNPS was obtained by Eq 3:

$$132 \text{NNPS} = \frac{\text{NPS}}{(\text{large area signal})^2} \quad (3)$$

133 The large area signal corresponds to the MPV in the image for each dose obtained from  
 134 STP.

135 The coefficient of variation (CoV(%)) was calculated by dividing the standard deviation of  
 136 the pixel values in the image by the mean pixel value. This was included to compare the  
 137 detector to other commercially available detectors for CBCT applications.

## 138 2.4. Pre-sampling Modulation Transfer Function (pMTF) Determination

139 The pMTF quantifies the resolution of an X-ray detector [20]. To measure the vertical  
 140 pMTF the polished edge of a tungsten test device, tilted by an angle of  $2^\circ$  relative to the pixel  
 141 rows was placed between two thick lead plates and attached directly to the digital detector.  
 142 Thirty images were captured at the highest current before saturation for each RQA to  
 143 decrease the statistical noise. Afterwards, the test device was rotated  $90^\circ$  clockwise to  
 144 measure the horizontal pMTF. Finally, the pMTF was obtained in the frequency domain by  
 145 fast Fourier transform (FFT) of the line spread function (LSF) which is the derivative of the

146 edge spread function (ESF). The pMTF values were calculated from zero to the Nyquist  
147 frequency and the pMTF at zero spatial frequency was normalised to one [18,19].

### 148 2.5. Detective Quantum Efficiency (DQE) Measurement

149 DQE is defined as the ability of an imaging system to transfer the input signal to an image  
150 [21,22]. It was computed using the following equation (Eq (4)):

$$151 \quad DQE(f) = \left( \frac{SNR_{Out}}{SNR_{In}} \right)^2 = \frac{pMTF^2(f)}{\frac{\Phi}{K_a} * K_a * NNPS(f)} \quad (4)$$

152 Where  $SNR_{Out}^2$  is signal-to-noise ratio square of the output signal on the image and is  
153 measured from the acquired images by dividing the  $pMTF^2$  by the averaged NNPS of the  
154 digital X-ray imaging device. The  $SNR_{In}^2$  is signal to noise ratio square of the input signal to a  
155 detector which can be estimated by multiplying the photon fluence per exposure ratio  $\left( \frac{\Phi}{K_a} \right)$   
156 in photons per  $mm^2/\mu Gy$  by air kerma in  $\mu Gy$  where  $\left( \frac{\Phi}{K_a} \right)$  values were provided by IEC  
157 62220-1 (2003). It was assumed that the detector behaves as an ideal photon counter  
158 [4,23,24].

### 159 2.6. Accumulation of Dark Current

160 This indicates the accumulation of dark charge in the pixel as a function of the integration  
161 time. It was measured using the following equation (Eq (5)):

$$162 \quad i_d = \frac{\bar{S}_d \times K \times q_e}{A \times T_{int}} \quad (5)$$

163 Where  $i_d$  ( $A/cm^2$ ) is the accumulation of dark current within the pixel,  $\bar{S}_d$  (digital number  
164 (DN)) is the mean dark signal at different integration times,  $K$  ( $e^-/DN$ ) is the conversion gain,  
165  $q_e$  is the electron charge which equals to  $1.6 \times 10^{-19}$  coulombs,  $A$  is the pixel area ( $cm^2$ ) and  $T_{int}$   
166 is the integration time (s) [24].

167 The effectiveness of dark frame subtraction, which is the first step in flat field correction,  
168 was assessed in terms of fixed pattern noise removal by subtracting two consecutive images  
169 having the same exposure time [26].

## 170 3. Results and Discussion

### 171 3.1. Signal Transfer Property (STP)

172 Fig. 1 describes the relationship between the  $K_a$  and MPV for all RQAs. The detector  
173 responded linearly at least within the range of investigated exposures ( $0.26-2.17 \mu Gy$  for  
174 RQA3 and  $0.29-1 \mu Gy$  for RQA5,7) with the coefficient of determination ( $R^2$ ) more than  
175 0.9995 in all cases. We observe a signal increase as the beam energy increases. As reported  
176 by [23,27], one explanation can be that the photon fluence per exposure ratio increases as the  
177 radiation energy increases, therefore, more signal carriers (X-ray photons) are travelling  
178 towards the detector as the energy increases. In our case, the signal per unit air kerma  
179 increased from 3387.5 X-rays/ $mm^2/\mu Gy$  at RQA3 to 9602.8 X-rays/ $mm^2/\mu Gy$  at RQA7.  
180 Alternatively, it could be that the production of the optical light in the scintillator is taking  
181 place closer to the sensor surface as the beam energy increases resulting in more light  
182 collection [28]. It is very noticeable that Lassena (CsI) has a high degree of sensitivity to  
183 radiation compared to other detectors [11,23,29]. It saturates at 2.17, 1.02 and 0.93  $\mu Gy$  for  
184 54 (RQA3), 74 (RQA5) and 92 kV (RQA7) respectively whereas in general radiography,  
185 using digital X-ray detectors, the  $K_a$  levels usually range from 0.8 to 8  $\mu Gy$  [30]. However,  
186 Lassena (CsI) can provide a satisfactory image quality at low exposures as will be explained  
187 in section 3.4. This, therefore, offers the potential to reduce the patient dose.

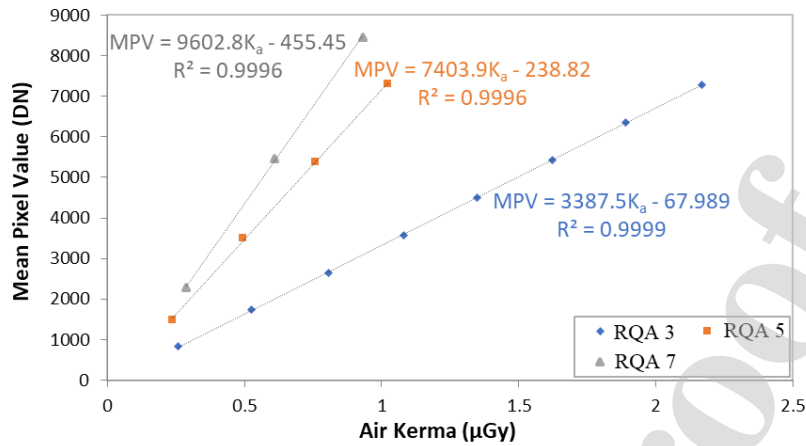


Fig. 1 The relationship between the MPV and  $K_a$  for all three energies.

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189

### 190 3.2. Pre-sampling Modulation Transfer Function (pMTF)

191 It was found that the pMTFs are independent of the beam quality within the investigated  
192 range (Fig. 2). Increasing the radiation energy slightly improved the resolution at low  
193 frequencies between 1 and 3 lp/mm. This is attributed to the longer mean free path of the  
194 higher energy X-rays resulting in a greater number of interactions closer to the sensor, which  
195 limited the spread of the scintillation photons [23]. The pMTF reaches 50% at 0.9, 1.08 and  
196 1.1 lp/mm for 54, 74 and 92 kV beam qualities respectively (Fig.2). As mentioned before, the  
197 frequency corresponding to 10 % MTF describes the limiting resolution of a system and it is  
198 around 3 lp/mm for all three beam qualities. In general radiography, the adequate limiting  
199 resolution for a detector ranges between 3 and 5 lp/mm [8,23]. Lassena (CsI) was compared  
200 to other available CBCTs on the market, the results are illustrated in table 1 [31-35]. The  
201 small pixel size gives Lassena (CsI) better resolution. For radiographic imaging comparison,  
202 table 2 [11,23,29] shows that Lassena (CsI) has modest pMTF resulting from the scintillator  
203 thickness that increases volumetric space for light to spread and scatter [17].

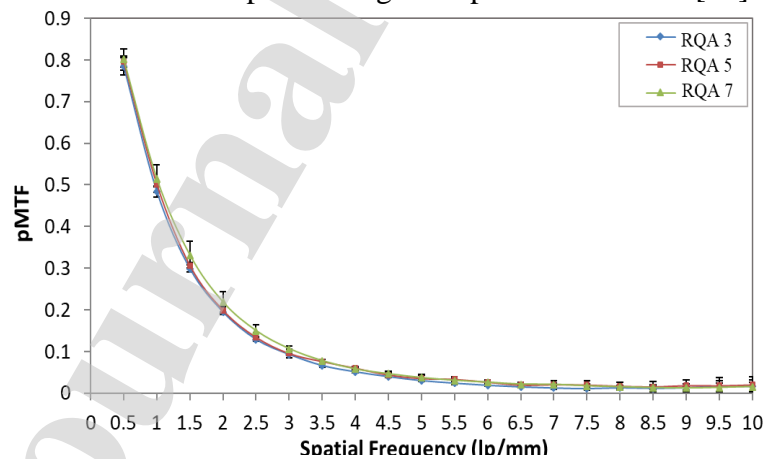


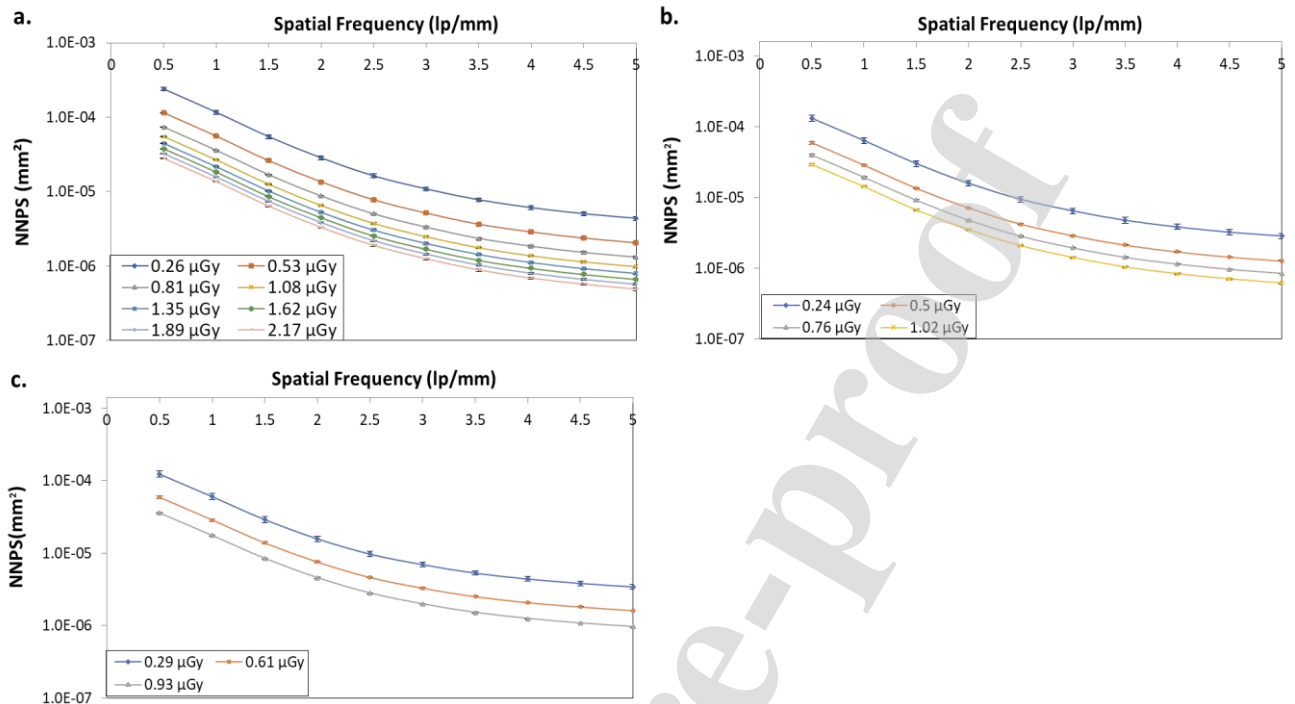
Fig. 2 The averaged pMTFs for Lassena (CsI) at three different energies.

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205

### 206 3.3. Normalised Noise Power Spectrum (NNPS)

207 For the detector under investigation, the NNPS at different energies demonstrated that the  
208 NNPS decreased as the radiation energy and dose (or  $K_a$  value) increased as shown in Fig. 3.  
209 This reduction is due to the intensification of the signal due to the higher number of photons  
210 interacting with detector material. Therefore, the number of absorbed photons increases  
211 [26,27]. This finding implies that NNPS heavily depends on exposure, consequently, it is  
212 expected that the DQE will increase at a higher dose since the DQE is inversely proportional  
213 to NNPS.

214 Lassena (CsI) has a CoV of 0.11% and it is the lower than other commercially available  
 215 detectors for CBCT applications as displayed in table 1.



216  
 217 Fig. 3 1D NNPS of a. RQA3 (54 kV) b. RQA5 (74 kV) c. RQA7 (92 kV) at different  $K_a$  values.

### 218 3.4. Detective Quantum Efficiency (DQE)

219 We found that the DQE at the three radiation energies reduced as a function of spatial  
 220 frequency and at high frequencies, the DQE became less exposure dependent due to the  
 221 increase of the photon shot noise and decrease of the pMTF [26,27]. On the other hand, DQE  
 222 improved when the radiation current (dose) and voltage (energy) rose as shown in Fig. 4. The  
 223 DQE (0.5) values are around 0.46 for all doses for RQA3 and they range from 0.52-0.56 for  
 224 RQA5, lastly, the values of DQE (0.5) for RQA7 are about 0.55-0.59. We can observe that  
 225 the investigated system demonstrates a quantum-limited condition for RQA3, in other words,  
 226 it is less exposure dependent at low energies due to remnant fixed pattern noise (FPN) or due  
 227 to the increased effect of CMOS APS inherent non-linearity as proved by [18]. However, the  
 228 detector shows higher DQE values at higher energies [36-38]. Looking at table 2, Lassena  
 229 (CsI) provides acceptable DQE values at low exposures as a detector for general radiography.



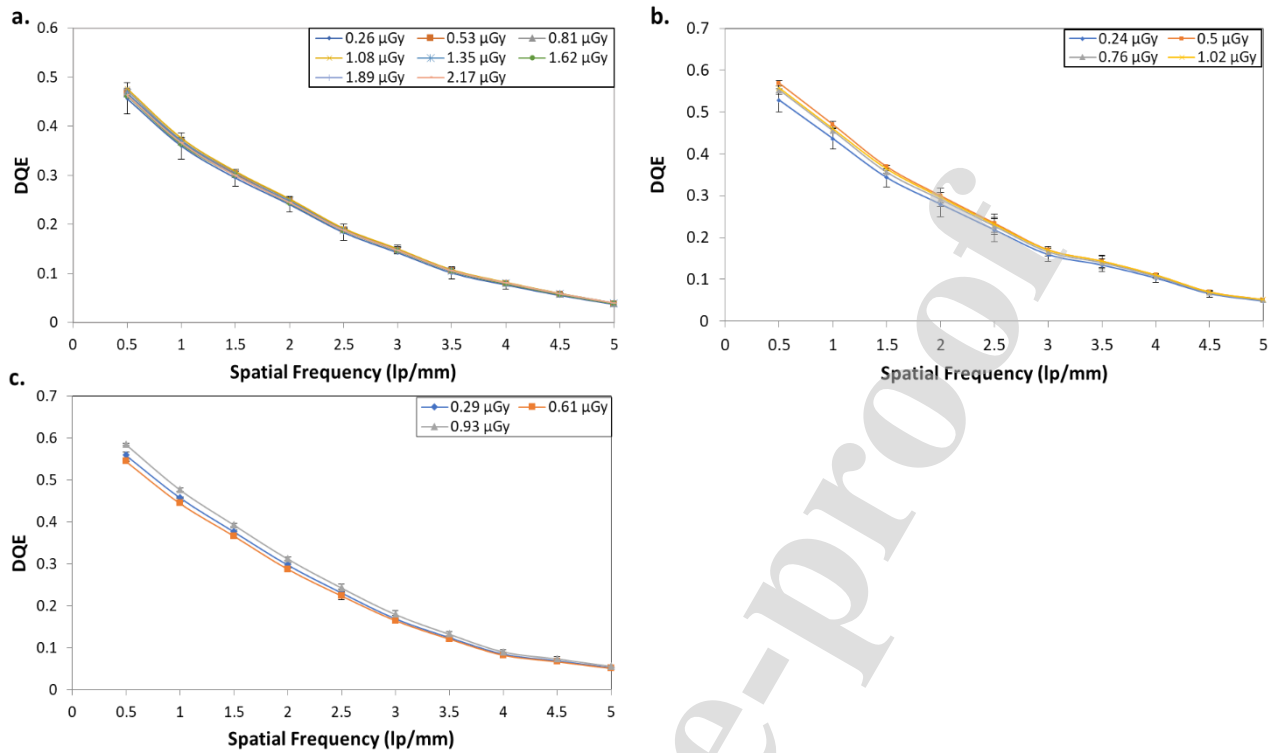


Fig. 4 DQE of a. RQA3 (54 kV) b. RQA5 (74 kV) c. RQA7 (92 kV) at different  $K_a$  values.

230  
231  
232

233 **Table 1 Comparison between the Lassena (CsI) and other available detectors that can be used for CBCT in the**  
234 **radiotherapy departments.**

|                                  | Varian       | Elekta      | Siemens     | Lassena (CsI) |
|----------------------------------|--------------|-------------|-------------|---------------|
| Resolution (Pixels)              | 2048 × 1536  | 1024 × 1024 | 1024 × 1024 | 2786 × 2400   |
| Physical size (cm <sup>2</sup> ) | 39.73 × 29.8 | 41 × 41     | 41 × 41     | 24 × 14.4     |
| Image depth (bit)                | 16-bit       | 16-bit      | 12/16-bit   | 14-bit        |
| Pixel pitch ( $\mu\text{m}$ )    | 388          | 500         | 400         | 50            |
| Max frame rate (fps)             | 30           | 5.5         | 25          | 30            |
| Tube voltage                     | 30-140 kV    | 70-150 kV   | 6 MV        | 54-92 kV      |
| MTF 50%/10% (lp/mm)              | 0.548/0.939  | 0.28/0.45   | 0.3/0.5     | 1.5/3         |
| Coefficient of variation (%)     | 0.7%         | 1.4%        | 2.7%        | 0.11%*        |

235 \*Obtained at the energy of 92 kV and  $K_a$  value of 0.93  $\mu\text{Gy}$ .

236 **Table 2** List of Studies that evaluated a new detector and the main findings compared to Lassena (CsI)

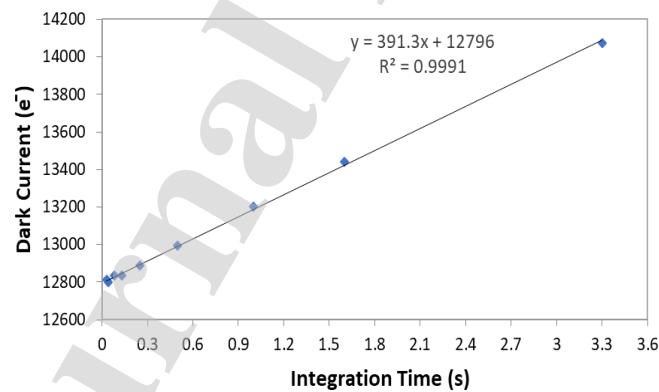
|   | Konstantinidis et al. (2012) | Michail et al. (2015) | Michail et al. (2016) | Lassena (CsI) |
|---|------------------------------|-----------------------|-----------------------|---------------|
| Resolution (Pixels)   | 3888 × 3072                  | 1200 × 1600           | 1200 × 1600           | 2786 × 2400   |
| Physical size (cm <sup>2</sup> )                                | 29 × 23                      | -                     | -                     | 24 × 14.4     |
| Image depth (bit)   | 14-bit                       | -                     | -                     | 14-bit        |
| Pixel pitch (μm)  | 74.8                         | 22.5                  | 22.5                  | 50            |
| Max. frame rate (fps)   | 26                           | -                     | -                     | 30            |
| Scintillator material   | CsI:TI                       | CsI:TI                | CsI:TI                | CsI:TI        |
| Scintillator thickness (μm)                                     | 200                          | 170                   | 490                   | 1000          |
| MTF 50%/10% (lp/mm)   | 1.2/4.5                      | 3.6/9.6               | 1.9/5.8               | 1.5/3         |
| <i>RQA3</i>   | 0.53-0.68                    | -                     | -                     | 0.45-0.47     |
| <i>DQE at 0.5 lp/mm</i><br><i>Range of dose</i><br><i>(μGy)</i> | 0.14-3.09                    | -                     | -                     | 0.26-2.17     |
| <i>RQA5</i>   | 0.68-0.75                    | 0.1                   | 0.8                   | 0.52- 0.56    |
| <i>Range of dose</i><br><i>(μGy)</i>                            | 0.13-6.45                    | 31.05                 | 8.39                  | 0.24-1.02     |

237 **3.5. Accumulation of Dark Current**

238 The accumulation of dark current was calculated using Eq (5) and it is 2.51 pA/cm<sup>2</sup> for 50  
 239 μm pixel pitch system; equivalently, the detector accumulates 391.3 e<sup>-</sup>/s in the absence of the  
 240 illuminations. In this case, we measured 12798 e<sup>-</sup> at 0.04 s and 14074 e<sup>-</sup> at 3.3 s (see Fig. 5).  
 241 The integration time of the current study was selected at 0.13 s [25,26].

242 Furthermore, the effectiveness of the fixed-pattern noise correction using dark signal noise  
 243 (DSN) subtraction was assessed. It was found that the correction removed 99.7 % of the dark  
 244 fixed-pattern noise as expected [26]. The mean noise dropped from 2053 DN to 4.3 DN after  
 245 the correction at the selected integration time (0.13 s) (see Fig. 6). This proves the  
 246 effectiveness of the DNS subtraction for fixed-pattern noise correction.

247



248

249

**Fig. 5** The relationship between the integration time and the dark current in the absence of illumination.

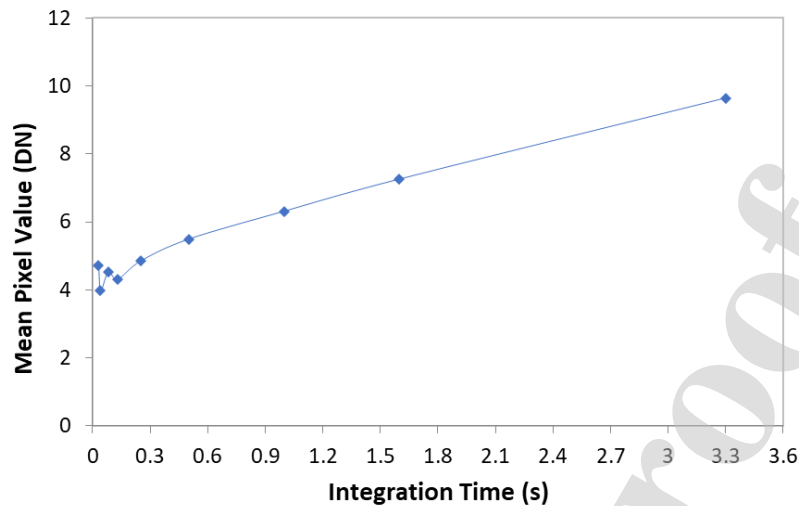


Fig. 6 The difference of two subsequent DNS images as a function of exposure time.

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251

### 252 3.6. Role of Scintillator (CsI) in Lassena Performance

253 The performance of this detector system should be considered in terms of the  
 254 scintillator (CsI) and the sensor (Lassena) because there are additional factors could  
 255 degrade the image quality if the sensor is not optimally designed. The Lassena CMOS  
 256 sensor has 40 e- rms noise while the more ubiquitous a-Si:H detectors have typical noise  
 257 values of 800-1000 e- rms [39]. The Rose criterion states a resolvable signal needs to be 5  
 258 times the noise level to be resolved [40]. For Lassena which has a quantum efficiency of  
 259 50%, a point source would need to emit 400 optical photons to be resolved while an a-  
 260 Si:H detector would need 8000 optical photons. This low noise performance comes at the  
 261 cost of dynamic range, Lassena has a full well capacity of 112,000 e- and as shown in  
 262 section 3.1 a bright scintillator such as CsI(Tl) will result in Lassena quickly saturating.  
 263 The dynamic range of Lassena could be significantly increased by using a dimmer  
 264 scintillator, but if the absorption remains the same then the noise at lower flux will  
 265 increase. Ideally, a scintillator with higher absorption with a lower light yield would be  
 266 preferred. Since low noise digital sensors such as Lassena are relatively new to the  
 267 market, no such scintillator is available in a suitable form for imaging. There are many  
 268 alternative scintillators in common use for other applications which have higher X-ray  
 269 absorption properties and lower light yields, but the K-edge location of the elements that  
 270 compose the scintillator is crucial for efficiency (Fig. 7). The Fig. 7 shows that energies  
 271 above 70 keV alternative scintillators begin to show a significant increase in SNR  
 272 compared to CsI (Tl) in terms of X-ray absorption.

273 The calculation of how signal and associated noise propagate through each stage of the  
 274 entire imaging system is known as quantum accounting and allows regions which limit  
 275 the DQE to be identified [41]. This approach has shown the benefit of increased X-ray  
 276 absorption in improving DQE. Traditionally the signal transfer stages that limit the DQE  
 277 also known as the quantum sinks have typically been in the collection of the optical light  
 278 from the scintillator. In other scintillator based detector systems, the quantum sink  
 279 associated with the collection of scintillation light, has been a result of the poor efficiency  
 280 of lens-based systems or high noise of the image sensors. Given the very low noise of  
 281 Lassena, the X-ray absorption of the scintillator is the most significant factor in limiting  
 282 the performance of Lassena and a high gain optical stage i.e. bright scintillator is no  
 283 longer beneficial but instead severely limits the dynamic range of the system.

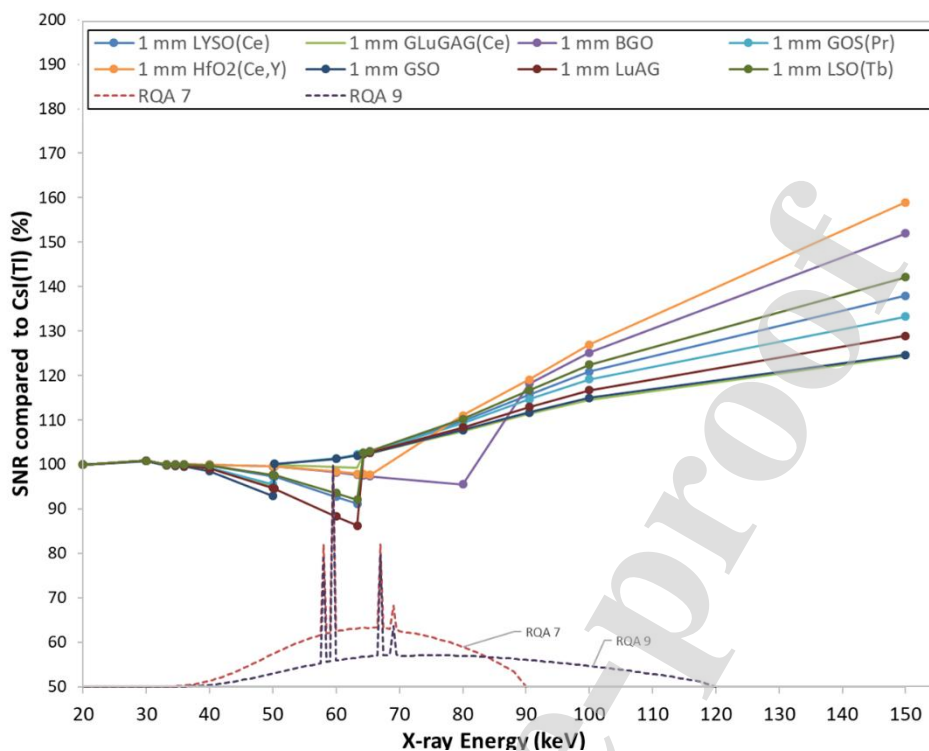


Fig. 7. The percentage difference in signal to noise ratio associated with X-ray absorption in common and uncommon scintillators compared to CsI(Tl). The RQA 7 and 9 spectra are overlaid on the figure for reference.

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#### 288 4. Conclusions

289 In this paper, a characterisation of the detector Lassena (CsI) performance was realised for  
290 the first time over a range of radiation energies and currents to determine its image quality as  
291 a detector for general radiography and CBCT applications. The detector responded linearly  
292 within the investigated dynamic range however it has very high radiation sensitivity which  
293 limited its dynamic range. Despite the scintillator thickness, this system presents a good  
294 limiting resolution (3 mm/lp) for a medical imaging application. The thick scintillator  
295 improves the SNR but also limits the dynamic range. The spatial resolution and NNPS results  
296 led to acceptable DQE at all energies, DQE (0.5) values were 0.46 for RQA3, 0.52-0.56 for  
297 RQA5 and 0.55-0.59 for RQA7 at an integration time of 0.13s. For CBCT applications,  
298 Lassena (CsI) showed very promising results, and the development of new scintillators to  
299 take advantage of low noise sensors, such as Lassena, could drastically improve the  
300 performance of imaging systems based on this type of sensor.

301

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**\*Declaration of Interest Statement**

**Declaration of interests**

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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