

1 **Acoustic Emission Waveform Picking with Time Delay Neural Networks During**
2 **Rock Deformation Laboratory Experiments**

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11 **Abstract**

12 We report a new method using a Time Delay Neural Network to transform Acoustic
13 Emission (AE) waveforms into timeseries of instantaneous frequency content and permutation
14 entropy. This permits periods of noise to be distinguished from signals. The model is trained
15 in sequential batches, using an automated process that steadily improves signal recognition as
16 new data are added. The model was validated using AE data from rock deformation
17 experiments, using Darley Dale Sandstone in fully drained conditions at a confining pressure
18 of 20 MPa (approximately 800m simulated depth). The model is initially trained by manual
19 picking of five high amplitude waveforms randomly selected from the dataset (experiment).
20 This is followed by semi-supervised training on a subset of 300 waveforms. When compared
21 with the standard amplitude-threshold picking methodology the source locations obtained from
22 the time difference of arrivals show ~400% increase in the number of events located, ~150%
23 increase in signal to noise ratios and a ~30% reduction in location residuals.

24

25 **Introduction**

26 Acoustic emissions (AE) are the transient elastic waves produced by the sudden
27 redistribution (release) of stress when a material fractures. These signals are considered to
28 represent the laboratory analogue of natural earthquakes with wider applications spanning from
29 non-destructive testing to vibration monitoring. In deformation experiments, AE are related to
30 the initiation and growth of fractures (e.g. Lockner et al., 1992), matrix cracking (Scholey et
31 al., 2010), fluid flow through fractures (Benson et al., 2019; Fazio et al., 2019), and fibre
32 breakage and various debonding processes (Bohse, 2000). Detection of these signals is a
33 valuable asset as they provide a non-destructive and immediate feedback to dynamically
34 evolving systems without the need for interference. However, a limitation of AE analysis is
35 that results are often dependent on the successful discrimination of signals from the background
36 noise. As datasets are often very large, therefore any approach requires an automated approach
37 that is robust enough to handle to waveforms in a medium that is continually evolving.

38 The simplest processing method to extract basic data uses stepwise calculation along the
39 waveform timeseries of energy ratios (“signal-to-noise” ratio). This type of analysis is often
40 used to extract discrete time series from a much longer (quasi-continuous) signal, with
41 triggering dependent on the arrival of large amplitude spikes (such as the first arrival of an
42 earthquake). Whilst intuitive and less computationally demanding than other techniques, this
43 method is prone to errors, as the largest signal-to-noise ratio may not refer to the first arrival
44 of a genuine signal (Guoping et al., 2004; Pomponi & Vinogradov, 2013). To reduce these false
45 triggers, methods that take advantage of the inherent characteristics of microseismic signals,
46 i.e., frequency content, have been proposed (Zhang et al., 2013). These methods use an
47 automated approach via the Hilbert-Huang Transform (HHT, Huang et al., 1998) to identify
48 low-frequency AE signals from high-frequency background noise. A further refinement of this
49 approach combines the HHT with the Akaike Information Criterion (Akaike, 1973) to pick up

50 the first arrival of seismic waves (Jia et al., 2015).

51 Entropy-based techniques, which identify the incoherent nature of the noise signal by
52 relating data to past and previous values, are also well documented. In medicine, Ródenas et
53 al. (2015) reported the accurate detection of atrial fibrillation from electrocardiograms. A
54 similar approach uses intrinsic mode functions (IMF) to recognise focal electroencephalogram
55 signals (Sharma et al., 2015). Similar techniques have found their use at the field scale. For
56 example, fluctuations in entropy were detected prior to the Mexican earthquake of 2017
57 (Ramírez-Rojas et al., 2018). Entropy is highly sensitive to seismic amplitude, frequency and
58 phase changes (Jia et al., 2019). They demonstrate immunity to noise under a wide range of
59 environmental conditions due to a sensitivity to the coherency of a real signal, making them
60 uniquely suited for seismological analyses. The technique can also be applied to active seismic
61 data (Zoukaneri & Porsani, 2015) by measuring the entropy of the instantaneous frequency
62 content, increasing the level of information available for sub-surface mapping.

63 Nonetheless, each of these techniques is limited by a user-defined ‘critical value’ that
64 defines when a time-period is “signal” or “noise”. Often, this threshold is set based on the user
65 experience and knowledge of the recording equipment. When tools are not calibrated, as is
66 often the case in AE data (e.g. Høgsberg & Krenk, 2015), such an approach remains imprecise
67 and will lead to further systematic errors regardless of the technique. Nevertheless, advances
68 in machine learning and artificial intelligence tools are now ushering a new era of data analysis,
69 which seeks to minimise biases enforced by user-defined parameters. Time Delay Neural
70 Networks (TDNN) are an automated signal classification technique that are designed with the
71 purpose of identifying patterns and trends in shift-invariant time-series data without explicitly
72 knowing the beginning or end of a signal (Derakhshani & Schuckers, 2004; Peddinti et al.,
73 2015; Waibel et al., 1995). Analogous to a 1D Convolutional Neural Network (CNN), TDNNs
74 are a form of Recurrent Neural Network (RNN) that models the propagation characteristics of

75 time series data similar to entropy techniques. By constructing models of the key elements of
76 audio, or elastic vibrations in the case of seismology, they can recognise different speakers
77 when employed for speech recognition (Haffner & Waibel, 1992) and perform robustly in the
78 presence of reverberations (Snyder et al., 2015).

79 Here, we seek to take advantage of TDNNs by applying a simple noise or signal
80 classification routine to AE data from a typical rock-deformation experiment using Darley Dale
81 Sandstone, where the TDNN is used to define the onset of energy as the boundary between the
82 two classifications. By iteratively training a neural network on calculated timeseries of seismic
83 envelopes, instantaneous frequency content and estimated signal entropy, the method picks AE
84 waveforms and automatically incorporates them into the evolving (updated) model, thus
85 improving later picks. To validate our approach, AE source locations are solved using the pick
86 times obtained from the trained model using a Time Difference of Arrival method (e.g.
87 Comanducci et al., 2020). AE distribution and error parameters are then compared to those
88 obtained from picking signals with a simple amplitude threshold method.

89 **Data and Method**

90 *Acoustic Emission Data*

91 A 40x100 mm cylindrical sample of Darley Dale Sandstone was deformed using a
92 conventional triaxial deformation apparatus (Fazio, 2017) at a confining pressure of 20 MPa,
93 representing approximately 800m depth. To ensure the experiment collected an AE dataset that
94 contains a large range of amplitudes (Nakamura et al., 1972) and fracture mechanisms
95 (Stanchits et al., 2006), the sample was deformed at a deformation rate of 3.6 mm/hr until brittle
96 failure following the protocol of Fazio et al., (2017). For a detailed overview of the AE data
97 acquisition methodology, please refer to **Appendix 1**.

98 *Model Parameters*

99 A previous study had proposed using only the waveforms as input data (Derakhshani &
100 Schuckers, 2004). After extensive testing, that approach was found to perform less robustly
101 with the high noise conditions of the AE measured here. Instead, we focus on the simple
102 attributes of signal amplitude and frequency content as *training* parameters to classify the
103 waveform. The former is modelled through the seismic envelope, whilst the latter by
104 instantaneous frequency content obtained through the HHT. A third parameter, permutation
105 entropy (Unakafova & Keller, 2013), encompasses both aspects. A fourth timeseries, defined
106 as the *target* model output, is used to classify the AE waveforms through binary separation of
107 noise and signal data, assigned -1 and 1 respectively. Signal is defined as the time period
108 between the calculated onset of energy and the point at which energy drops below a pre-pick
109 noise threshold. Due to uncertainties in the waveform content following the main arrival (i.e.
110 reflections) the model is only trained on noise identified during the pre-signal period.

111 The seismic envelope is the most intuitive and useful parameter to apply as it represents
112 the instantaneous energy of the signal (**Figure 1a**). Typically, noise is represented by low
113 amplitude values and signal by high amplitude values. The boundary between these two is often
114 defined by a rapid increase in values as a signal arrives. However, in the presence of low
115 amplitude data, or strong scattering where the envelope of the signal can become broadened,
116 the increase becomes shallower or simply not present (De Siena et al., 2016). To address this,
117 the HHT provides a measure of instantaneous frequency content along the entire time series
118 (**Figure 1b**) that is independent of amplitude, therefore any boundary between classifications
119 is still preserved when energies are low. Noise is identified through a dominance of high-
120 frequency energy, whilst the signal is identified through a consistent dominance of low
121 frequency waves. However, on its own, this method is likely to fail in the presence of high-
122 frequency signals that are similar in content to the background.

123 The final *training* parameter, permutation entropy, describes the uncertainty and the

124 degree of irregularity of a random timeseries (**Figure 1c**). Ordinal patterns are calculated that
125 represent the relationship between past and future values at each time step of the waveform
126 (i.e. an increasing trend). A probability density function of the number of patterns within a
127 moving window is then calculated, thus removing any dependence on the amplitude of the
128 original waveform (Unakafova & Keller, 2013). The larger the value of permutation entropy,
129 the higher the diversity of ordinal patterns is and the more complex the input data are. Thus,
130 noise is characterised by high values and signal by low, however, the boundary is slightly
131 shifted to that of the seismic envelope. A more detailed overview on the *training* data is
132 provided in **Appendix 2**.

133 ***Training Routine***

134 Similar to the previously described entropy method, a TDNN models the temporal trend
135 (or trajectory) of the *training* data, classifying patterns in the data according to the *target*
136 timeseries. A key difference of a TDNN from other multi-layer neural networks is that patterns
137 are classified with shift-invariance; a specific pattern may occur at any point in the timeseries.
138 With each successive layer in the network, increasingly coarse trends in the *training* data are
139 identified, therefore ensuring the model learns the most important features of the chosen
140 classifications.

141 Five high-amplitude waveforms are randomly selected from an AE pool to create an
142 initial model that is trained under supervised conditions. Their onsets are manually picked, and
143 the model iteratively trained on the concatenated input parameters (*training* dataset) after each
144 waveform. At this stage, the model already has a relatively high degree of picking accuracy.
145 Following this, the model undergoes semi-supervised training. A waveform of any amplitude
146 is randomly selected from the available pool. The input parameters are both calculated and
147 simulated in the neural network model generating an output.

148 Whilst easier than picking the signal directly from the waveform, it is still difficult to
149 pick the signal onset from the model output due to the uncertain transition from noise to signal.
150 With knowledge of the sample velocities, an effective approach was to define a window
151 between the origin time (or first arrival time when picking subsequent waveforms in a single
152 AE group) and the maximum possible arrival time for energy traversing the length of the
153 sample. Within this window we calculate the instantaneous frequency content of the model
154 output in order to characterise the rapid change in model values as a signal is detected. The
155 onset of energy is then set when the frequencies and model output exceed 14 KHz and -0.95,
156 respectively.

157 Once the onset is identified, two ‘picking quality’ ratios are calculated around this value.
158 A short-term SNR of the original waveform and a SNR of the model output. After extensive
159 testing, we set their values to 2.0 for the former and between 0.3 and 0.9 for the latter as they
160 identified a reliable onset when the training dataset was still small. When both ratios are
161 exceeded, the input parameters are added to the training data and the neural network is updated.
162 To improve computation speed, the model is trained in batches of 10 waveforms. Waveforms
163 can be further down-sampled to increase computation speed, although this will decrease data
164 quality. As this is an iterative process (**Figure 2**), the model is steadily improved over time,
165 which allows for the later inclusion of data that may not have been initially included. The model
166 ceases training once the *training* dataset exceeds 300 waveforms as the quality of picking did
167 not improve significantly after this stage.

168 To provide a robust comparison and best highlight the improvements offered by the
169 proposed neural network approach, AE data are also processed using a simple amplitude-
170 threshold method to pick data. The root mean square (RMS) envelope of each waveform is
171 calculated and the pick time is chosen when the envelope exceeds 1.1 times the level of the
172 background noise (using the same windowing approach as before).

173 *Source Location*

174 Pick times are inverted for source location using a Time Difference of Arrival (TDOA)
175 method (Tobias, 1976). Although this method is only robust in the case of weak velocity
176 anisotropy, the TDOA method is well established in micro-seismic source location (e.g.
177 Comanducci et al., 2020) and works well for the near-isotropic Darley Dale sandstone as used
178 here. These data are calculated by pairwise subtraction of the time of arrival values to each
179 sensor from a single source. A fixed velocity is used in the calculation that is updated as time
180 progresses (see **Appendix 1: Rock Deformation Experiment**). The calculated TDOA values
181 are determined through iterative estimation of the source location. The process iterates the
182 location using the L2 norm of the location residual, comparing between calculated and TDOA
183 values. It arrives to a local minimum at the best determined source location. To be valid,
184 sources are located with a minimum of 6 arrival times, although the accuracy greatly improves
185 with the addition of more sensors.

186 **Results and Discussion**

187 At high amplitudes, both the neural network (NN) and the amplitude threshold method
188 (AT) performed equivalently. This is unsurprising as the seismic envelope is a key indicator of
189 an AE arrival in both methods. However, as amplitudes decrease and reach the level of the
190 background noise, discrepancies in pick times become increasingly apparent (**Figure 3**). Whilst
191 it still performed relatively well with impulsive arrivals, the AT method (**Figure 3**, blue line)
192 struggled with the strongly scattered or emergent arrivals that are characteristic of low
193 amplitude AE data. Such data are challenging for picking methods, as the first arrival of the P-
194 wave becomes more discriminatory due to the convolution of many wave-modes that extend
195 into the coda (e.g. Grosse & Ohtsu, 2008). The NN on the other hand, was much more
196 consistent in capturing the first arrival regardless of amplitude, wave-mode or waveform

197 character. By modelling characteristic sequences of data, the TDNN is able to classify the
198 patterns that uniquely define signal or noise in the timeseries. From the model output (**Figure**
199 **3**, red line), the arrival time may then be obtained from a far simpler timeseries, even for data
200 that was not picked by the AT method. It should be noted that the windowing scheme was
201 crucial to avoid picking late high energy coda arrivals or even the coda from a prior AE in the
202 pre-pick noise.

203 Although relevant in all forms of waveform analysis (e.g. seismic, microseismic), the
204 occurrence of external (outside the sample) or electronic noise (often high amplitude, high
205 frequency) with temporal and frequency characteristics similar to real events is a characteristic
206 element of AE data processing. Waveform frequency data has already been shown to be a
207 reliable parameter to characterise waveform data (96% accuracy), but still leads to false-alarms
208 in the pre-pick noise without additional constraints on arrival time (Jiang & Zheng, 2020).
209 Similar results are obtained when only amplitude data is used (Chen et al., 2019), but such an
210 approach is likely to fail in high-noise environments. The method presented here is distinct
211 from the previously mentioned studies in that it combines multiple input parameters to train
212 the model (amplitude, frequency and entropy). This provided a notable advantage over single
213 feature methods as a more robust characterisation of noise (rather than only signal) was
214 obtained. However, the TDNN performed poorly in the presence of very high amplitude data
215 ($\text{SNR} > 800$). Due to the way the signal is decomposed during the HHT, the noise window is
216 characterised by low frequency information thus resulting in errors during training. As these
217 data occurred very rarely, it was not considered detrimental to the overall model.

218 Sources that were located using both methodologies reveal the time-dependent evolution
219 of the fault plane through the entire experiment (**Figure 4**Error! Reference source not found.).
220 Although events with high location residuals are displayed here for completeness, the
221 distribution of AE hypocentres is consistent with post-test visual inspection of the final fault

222 (Figure 4a, e.g. Benson et al., 2010; Lockner et al., 1992). Considering the evolution through
223 time (Figure 4b), we note that the event locations are diffuse throughout the sample during the
224 early stages of deformation (windows 1, 2 and 3). As strain increases, during the period of
225 strain-softening (window 4), the distribution of AE locations starts shaping a fault structure
226 (strain localisation). During crack coalescence and brittle failure (window 5) event locations
227 are primarily located along the observed failure plane. Whilst the general distribution of events
228 remain the same for both methods, the inconsistencies in waveform picking that were shown
229 in Figure 3 have resulted in significant differences in source location that can differ by several
230 centimetres for individual AE (Figure 4, coloured dots). A key component of the presented
231 method is that the model is trained on data that came from same the dataset that is due to be
232 picked. As AE datasets are typically very large, it is impractical to manually pick a dataset to
233 truly validate our source locations. Nonetheless, a focusing of events towards a failure plane is
234 much more evident when using picks from the NN compared to the AT method.

235 In terms of overall performance TDNNs (or RNNs) are generally considered to be
236 equivalent to the more powerful Convolutional Neural Network (CNN) for use in signal
237 detection (G et al., 2018; Guo et al., 2020; Yao et al., 2020). Moreover, the number of training
238 data required to create a well performing model also remains relatively low for both approaches
239 (e.g. Chen et al., 2019). An important difference, however, is that the TDNN presented here
240 outputs a transitional classification rather than the absolute output of a CNN. This provides a
241 measure of uncertainty that is time dependent and thus physically much more appropriate for
242 data generated by time-dependent processes, such as waveform data and for the vast majority
243 of pre-failure processes.

244 The semi-supervised approach presented here may be considered advantageous due to
245 the low workload requirements when training the model on new datasets. To initiate a new
246 training routine, only 5 high amplitude waveforms needed to be manually picked before the

247 model could train itself. In addition, as the model improves more low amplitude data is
248 incorporated into the model automatically, thus removing the requirement to manually pick
249 difficult data. The choice of network, CNN or TDNN, is therefore tuned to the end-user
250 requirements rather than on any standard performance metrics.

251 In any case, the models derived using the method proposed here may be applied to new
252 datasets or modified as a triggering tool to use with continuous or real-time data. An important
253 advantage of the TDNN over CNNs is that the input timeseries may be of any length, regardless
254 of those used in the training routine. Rather than identifying individual elements (such as the
255 first arrival), TDNNs characterise sequences of features in order to classify patterns. Thus, a
256 continuous dataset may first be down-sampled to generate a quick approximation of arrival
257 times (trigger) and the exact same model may then be used on the higher resolution recording
258 to obtain a more accurate pick.

259 **Conclusions**

260 This study has presented a workflow and application of a machine learned waveform
261 picking tool. A time delay neural network is trained to recognise the onset of AE energy using
262 instantaneous frequency, seismic envelope and entropy measurements. Statistical results
263 demonstrate the reliability of the method and highlight the potential of using multiple
264 waveform characterisation techniques to determine the arrival of acoustic energy.

265 We conclude that an automated process may be readily applied to AE datasets, as
266 recorded here from a triaxial rock deformation experiment. Aside from the 5 waveforms used
267 to initially train the model, no user intervention is required. A high degree of picking accuracy
268 and the inclusion of low-amplitude data that may be missed by traditional single-parameter
269 threshold methods results in datasets with a high source location recovery rate and a reduction
270 in the propagation of errors in further analysis of the data.

271 **Data and Resources**

272 Acoustic Emission data are obtained at the Rock Mechanics Laboratory, University of
273 Portsmouth. Raw data files, pick times, source locations and errors are available at
274 <https://zenodo.org/record/3958910>. Programming codes were developed in MATLAB®
275 version 2018a and are available at <https://github.com/thomaskingunito/programming>.

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279 and David Dempsey for their helpful comments during the development of the code.

280 **Appendix 1: Rock Deformation Experiment**

281 Darley dale sandstone (DDS) is ubiquitous in rock mechanics testing. It is a brown-
282 yellow, feldspathic sandstone with a modal composition of quartz (69%), feldspars (26%), clay
283 (3%) and mica (2%) (Heap et al., 2009). Previous studies report a connected porosity of 13.3%
284 \pm 0.8% with grain sizes varying from 100-800 μ m (Zhu & Wong, 1997). Pore sizes measured
285 by mercury porosimeter have a mean diameter of 10 μ m, but with significant microporosity
286 (Read et al., 1995). The unconfined compressive strength is 160 MPa (Baud & Meredith, 1997).
287 At the scale analysed here, no distinct layering or laminations were present. A cylindrical rock
288 sample was cored using a diamond tipped hollow coring drill to prepare a 4 cm diameter sample
289 that was then trimmed to 10 cm length with a diamond saw. End faces are accurately ground
290 using a lathe fitted with a cross-cutting diamond grinding disk with surfaces flat and parallel to
291 within 0.01 mm.

292 Deformation was performed using a conventional triaxial deformation cell installed at
293 the Rock Mechanics Laboratory, University of Portsmouth (Fazio, 2017). The sample

294 presented here was deformed until brittle failure at a confining pressure of 20 MPa at a constant
295 deformation rate of 3.6 mm/hr. Experimentation was performed under fully drained conditions
296 to avoid any fluid-driven effects on AE frequency content (Benson et al., 2010). These
297 environmental conditions ensure that a high number of AE are obtained and any time-
298 dependent variations in the signal waveform are predominantly due to the scattering effects of
299 microfractures, thus allowing for the sampling of a diverse range of deformation structure.
300 Axial displacement is measured with a non-contact Eddy Displacement system mounted to the
301 apparatus. It comprises of three sensors that accurately (sub-micron) measure the distance to a
302 target steel plate attached to the driving piston. These readings are averaged and are used to set
303 the target deformation rate via feedback to an axial stress intensifier.

304 For AE data acquisition the protocol of Benson et al. (2007) was followed. The dry
305 sample was positioned inside an engineered rubber jacket fitted with ports for an array of
306 twelve 1 MHz single-component Piezo-Electric Transducers (PZTs, model PAC Nano30) were
307 embedded (Error! Reference source not found. Error! Reference source not found.). These
308 sensors have a relatively flat frequency response between 125-750 KHz. Sensor output is
309 connected to preamplifiers set to 40 dB, focusing on data quality over quantity. An ITASCA-
310 Image “Milne” recorder operate in a standard ‘trigger’ model, downloading all twelve channels
311 when any single channel passes a set 100 mV threshold (e.g. Gehne, 2018). During
312 experimentation, in addition to ‘passive’ mode (recording of fracturing events) the sensors are
313 also used in ‘active’ mode for P-wave elastic velocity measurements in order to derive a
314 velocity model for source locations (**Figure A2**). Periodically, each PZT was triggered in
315 sequence with a high voltage (200 V) pulse, with the energy recorded by the remaining 11
316 sensors in the array. Signal to noise (SNR) was further improved by pulsing each sensor
317 multiple times (16) and stacking the received waveforms, with each survey of 12-sensors
318 generating 144 non-unique raypaths and taking approximately 30 seconds to complete.

319 Mechanical data from the deformation testing reveals a typical response of DDS. A total
320 of 24360 AE were detected by the ‘trigger’ recording system (**Figure A2**, histogram 0.01%
321 strain bins), however, it is unknown what percentage of these signals are noise related.
322 Differential stress increases gradually (**Figure A2**, black line), entering the elastic phase of
323 deformation at approximately 0.3% strain, coinciding with a rapid increase in the number of
324 AE recorded. During the period of strain-softening from 0.7-0.9% strain the rate of AE
325 dramatically accelerates and peaking at ~1000 events/0.01% strain. As the sample goes through
326 peak stress (145 MPa) several pulses in AE count correspond with minor drops in differential
327 stress. However, the overall rate decreases with time. At ~1.4% strain the sample fails, with
328 **Figure 4a** Error! Reference source not found. highlighting the formation of a fracture plane
329 marked in red.

330 **Appendix 2: Model Training Parameters**

331 *Seismic Envelope*

332 The Root Mean Square (RMS) envelope provides a scaled amplitude estimate of the AE
333 trace. The envelope represents the instantaneous energy of the signal and is computed in a
334 sliding window of 10 samples. Such a narrow window is selected to avoid smoothing of the
335 low amplitudes at signal onset. In low-noise environments the amplitude of the signal greatly
336 exceeds that of the pre-signal noise and provides a good marker of the onset of energy.
337 However, as noise increases, this onset becomes increasingly difficult to identify resulting in a
338 high degree of uncertainty.

339 *Instantaneous Frequency Content*

340 Intrinsic modal functions represent simple oscillatory modes of complex signals. Unlike
341 harmonic signals, IMF can have variable frequency and amplitude content through time.
342 Known as empirical mode decomposition (EMD), IMF of individual AE are obtained through

343 a continuous screening process called “sifting” and must satisfy the following requirements:
 344 The number of extrema and the number of zero-crossings must either be equal or differ by a
 345 minimum of one, whilst the mean of the upper and lower envelopes is zero. The procedure
 346 follows as thus (Huang et al., 1998):

- 347 1. The local extrema in the time series $X(t)$ are identified
- 348 2. Local maxima are connected by a cubic spline line to produce an upper envelope
- 349 3. This is repeated for the local minima to produce a lower envelope

350 The difference between $X(t)$ and the mean of the two envelopes, m_1 , is the first
 351 component h_1 :

$$352 \quad h_1 = X(t) - m_1.$$

353 Using h_1 as the new time series, this process is repeated k times until the standard deviation of
 354 h_{1k} , σ_k , is less than 0.8. This value was selected as it allows for a relatively small number of
 355 IMF components, whilst still capturing signal complexity:

$$356 \quad \sigma_k = \sum_{t=0}^T \frac{|h_{k-1}(t) - h_k(t)|^2}{h_{k-1}^2(t)}.$$

357 Thus, h_{1k} is defined as the first IMF component of the data:

$$358 \quad c_1 = h_{1k}.$$

359 Typically, c_1 contains the highest frequency component of $X(t)$. Subsequent IMF, which
 360 contain longer period data, are then calculated from the residual r_n where:

$$361 \quad r_n = r_{n-1} - c_n.$$

362 The sifting process is stopped when r_n becomes a monotonic function from which no new IMF
 363 can be extracted. Consequently, the test data is decomposed into n empirical modes where

364
$$X(t) = \sum_{j=1}^n c_j + r_n.$$

365 Once obtained, the Hilbert transform H , or instantaneous amplitude A_{inst} , of each IMF
 366 component is calculated as:

367
$$H(c_n)(t) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{c_n(\tau)}{t - \tau} d\tau.$$

368 Where τ is the width of the analysis window. The phase angle θ ,

369
$$\theta = \text{Im}(\log(H(c_n))),$$

370 is then used to compute the instantaneous frequency content f_{inst} as

371
$$f_{inst} = \frac{\Delta(\theta) / Ts}{2 * \pi},$$

372 where Ts is the sampling period. To characterise the data for use in classification, each step of
 373 the vector $freq(t)$ is then defined as the $f_{inst}(t)$ of the highest $A_{inst}(t)$ of all IMF components
 374 in $X(t)$. Thus, $freq$ is simply a high-resolution vector of the dominant frequency content
 375 through time.

376 ***Permutation Entropy***

377 The Permutation Entropy method describes the uncertainty and the degree of irregularity
 378 in a random series. To compute the empirical PE (refer to Unakafova and Keller, (2013) for
 379 more details), ordinal patterns of the AE are obtained with delay of $\tau = 1$, indicating a distance
 380 of 1 between points in patterns and an order of $d = 5$, meaning patterns contain 6 points ($d+1$).
 381 With a high degree of overlap between points, it is possible to use all of the information about
 382 order relations between points in the AE waveform. The distribution of ordinal patterns are
 383 obtained by the simple enumeration of the type of pattern. The PE, computed in a sliding time
 384 window of 30 points to reflect the complexity in the time series, is calculated at time t as:

385
$$PE_d^{\tau}(t) = - \sum_{j=0}^{(d+1)!-1} \frac{q_j}{M} \ln \frac{q_j}{M} = \ln M - \frac{1}{M} \sum_{j=0}^{(d+1)!-1} q_j \ln q_j,$$

386 where $q_j = \#\{k \in \{t, t-1, \dots, t-M+1\} \text{ has the ordinal pattern } j\}$ (with $0 \ln 0 := 0$) and M
 387 is defined as the number of ordinal patterns in the window.

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516 **Figure Legend**

517 Figure 1: Input parameters used in model training. Data classifications, indicated above, are noise and signal. a)
518 Seismic envelope. High values typically denote signal, however, low amplitude data is poorly characterised. b)
519 Dominant frequency content of the AE waveform. Red and blue indicates high and low frequency content
520 respectively. Pre-signal noise is characterised by high frequency and the signal itself with low values. c)
521 Permutation energy behaves similarly to seismic envelope but the trend is opposite with low values identifying
522 signal. Even at low amplitudes, permutation entropy detects coherency in frequency content.

523 Figure 2: Flowchart detailing key elements of the training loop.

524 Figure 3: AE signal (black), neural network model output (NN, red) and amplitude threshold (AT, blue) highlight
525 the discrepancies in picking quality between the two methods at low amplitudes. Strongly scattered or emergent
526 waveforms are frequently misspiked due to late arrivals in the coda. The AT method also misses many low
527 amplitude arrivals that the NN is still able to detect.

528 Figure 4: a) Post-failure imaging of the sample highlights a singular failure plane propagating to the lower left. b)
529 AE sources that are located by both methodologies (black dots). Example locations (corresponding coloured dots)
530 highlight the discrepancy in source location due to the different picking methodologies.

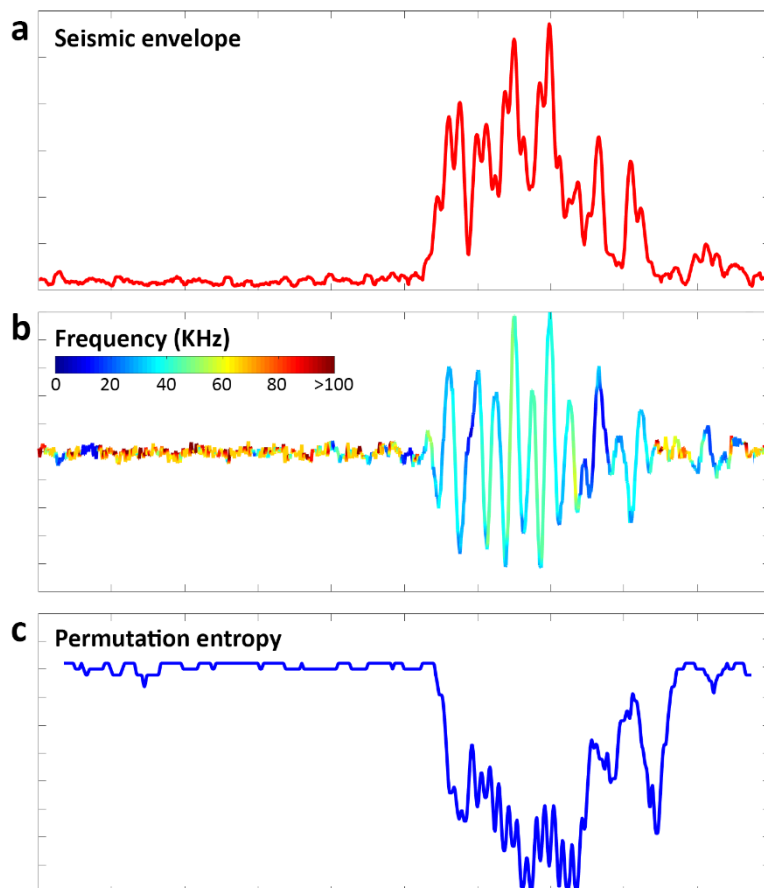
531 Figure A1: Sensor array.

532 Figure A2: Velocity data, stress-strain curve and raw AE count for Darley Dale Sandstone deformed at 20 MPa
533 confining pressure.

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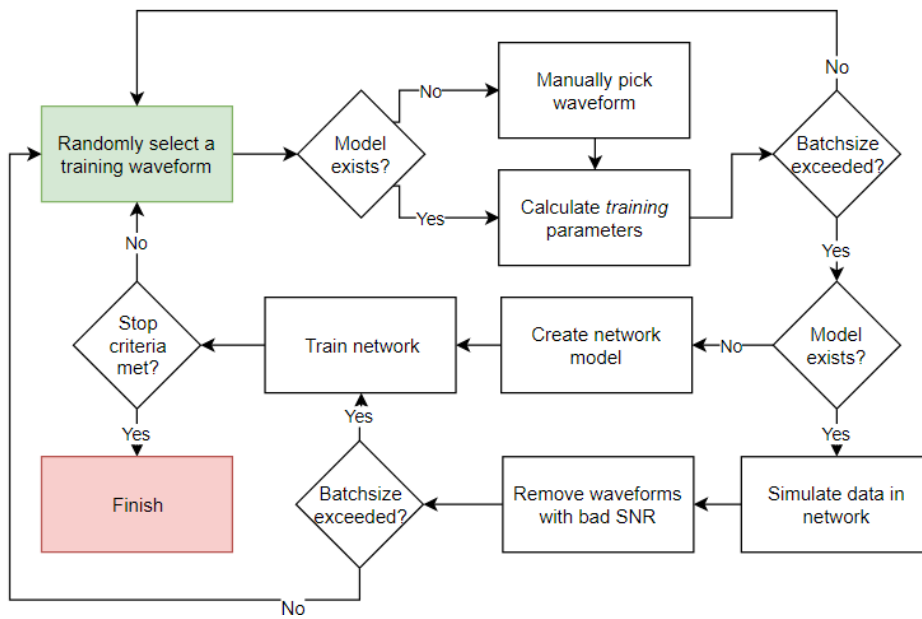
536 **Figures (In order of appearance)**



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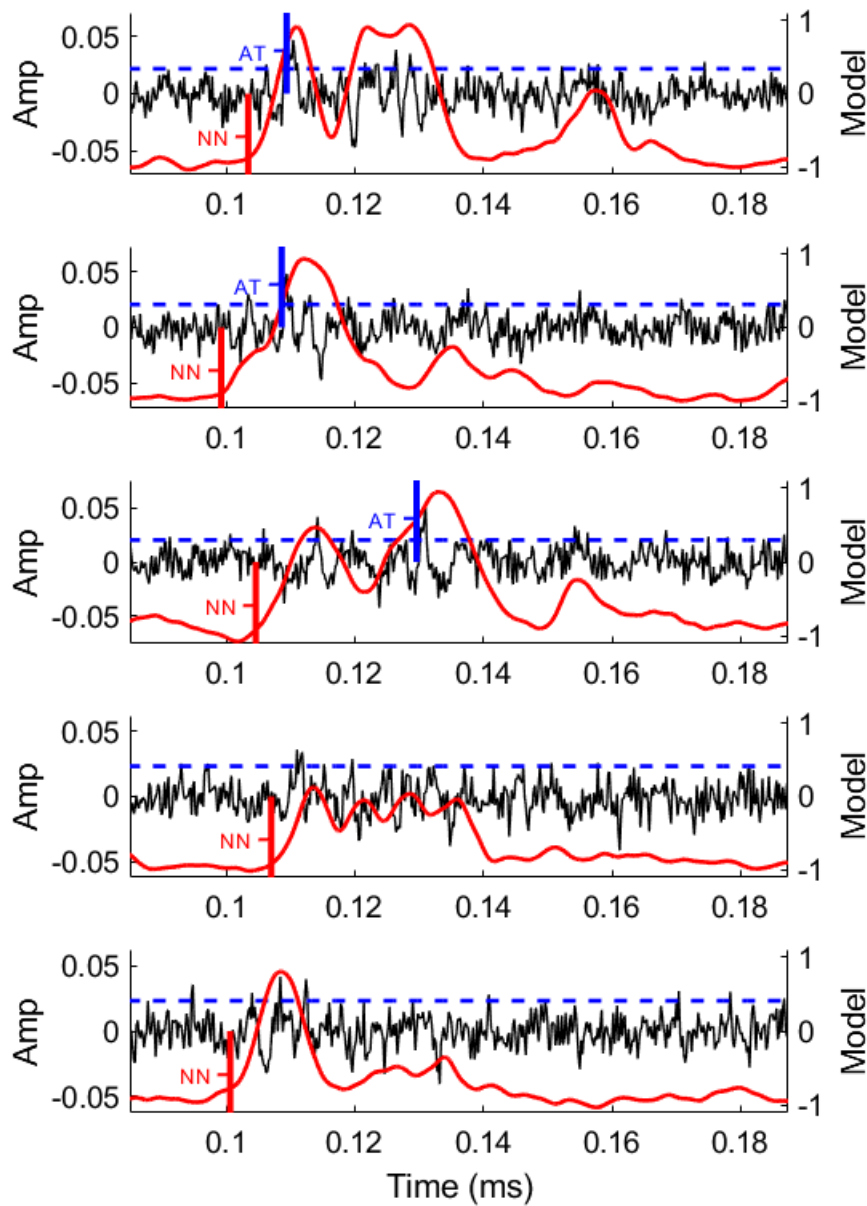
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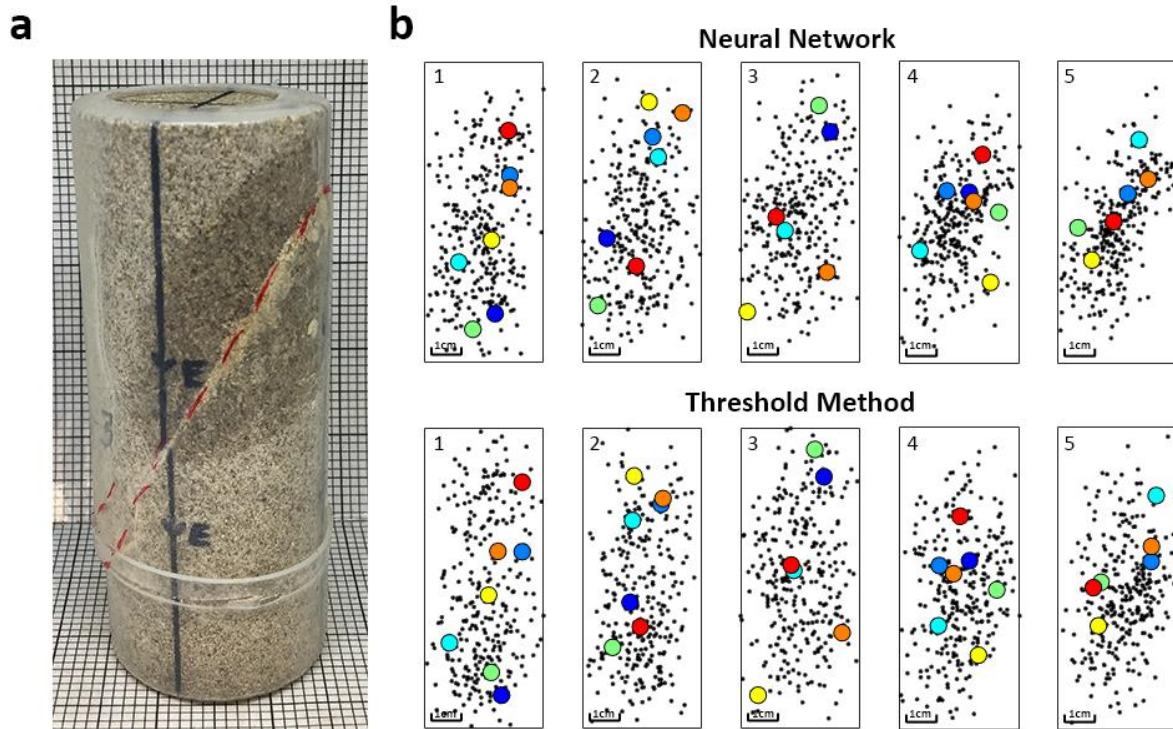
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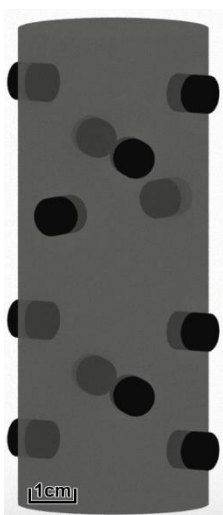
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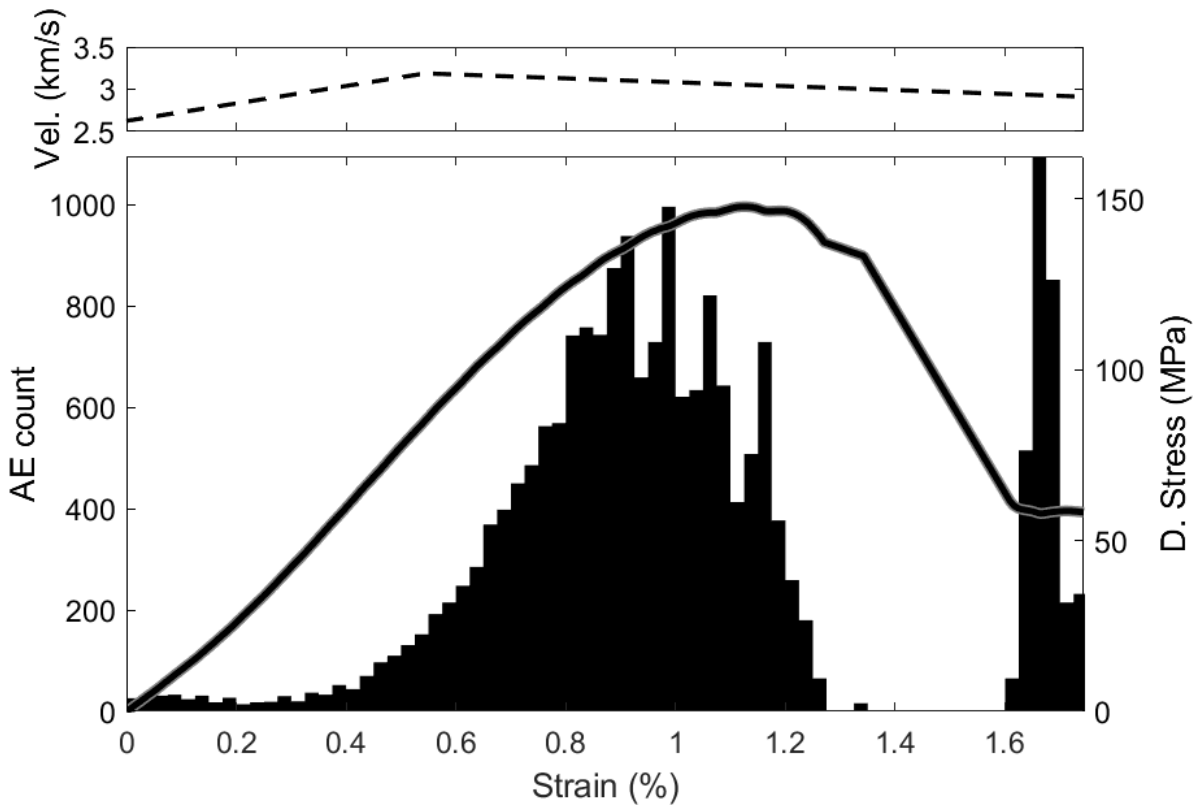
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557 Figure A1: Sensor array.

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560 Figure A2: Velocity data, stress-strain curve and raw AE count for Darley Dale Sandstone deformed at 20 MPa
 561 confining pressure. Brittle failure occurs at approximately 1.2% strain.

562